Positron Collection system

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Positron collection system



- 3 Adiabatic matching device
- 4 Conclusion & Questions





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- 2 Quarter wave transformer
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Figure: schematic and purpose of the positron collection system symbolized here by the matching device item.



• The phase space in a solenoid is described by :

$$\left(\frac{eB}{2}\right)^2 (x^2 + y^2) + (p_x^2 + p_y^2) = Cte$$
 (1)

Where :

- e is the particle charge.
- B is the magnetic field in the solenoid.
- x, y Spacial coordinates the solenoid.
- p_x , p_y are the particles momenta.
- This equation define the Hyper ellipsoid in the phase space.



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Quarter Wave Transformer





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Quarter Wave Transformer

• The position vector $\vec{q}(x, p_x, y, p_y)$

• The transfer matrix from the target z₀ to the exit of the solenoid according to the longitudinal position z_s can be written :

$$\vec{q}_B = M \left(z_s | z_0 \right) \, \vec{q}_0$$

$$M(z_{s}|z_{0}) = R_{2} M_{2} R_{1} M_{1}$$

Where:

- M_1 Transfer matrix in the first solenoid.
- M_2 Transfer matrix in the second solenoid.
- R_1 and R_2 are the rotation matrix.
- After decoupling:



$$\begin{pmatrix} X \\ P_X \end{pmatrix} = e^{-i(\xi_1 + \xi_2)} \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} \begin{pmatrix} X_0 \\ P_{X_0} \end{pmatrix}$$

(2)

Quarter Wave Transfor

where

$$M_{11} = \cos \chi_1 \cos \chi_2 - \frac{B_1}{B_2} \sin \chi_1 \sin \chi_2$$
(3)

$$M_{12} = \frac{2}{eB_1} \sin \chi_1 \cos \chi_2 + \frac{2}{eB_2} \cos \chi_1 \sin \chi_2$$
(4)

$$M_{21} = -\frac{eB_2}{2} \cos \chi_1 \sin \chi_2 - \frac{eB_1}{2} \sin \chi_1 \cos \chi_2$$
(5)

$$M_{22} = -\frac{B_2}{B_1} \sin \chi_1 \sin \chi_2 + \cos \chi_1 \cos \chi_2 .$$
(6)

• The rotation angle in the (x, y) plane

$$\chi = \int_0^l \frac{eB}{2p} dz$$



(7)

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$$XX^{*} + \left(\frac{2}{eB_{2}}\right)^{2} P_{X} P_{X}^{*} = \left[\cos^{2} \chi_{1} + \left(\frac{B_{1}}{B_{2}}\right)^{2} \sin^{2} \chi_{1}\right] x_{0} x_{0}^{*} \qquad (8)$$
$$+ \left[\left(\frac{2}{eB_{1}}\right)^{2} \sin^{2} \chi_{1} + \left(\frac{2}{eB_{2}}\right)^{2} \cos^{2} \chi_{1}\right] p_{x_{0}} p_{x_{0}}^{*}$$
$$+ \frac{2}{eB_{1}} \sin \chi_{1} \cos \chi_{1} \left[1 - \left(\frac{B_{1}}{B_{2}}\right)^{2}\right] (x_{0}^{*} p_{x_{0}} + x_{0} p_{x_{0}}^{*})$$

•
$$X = x + iy$$
 and $P_x = p_x + ip_y$
• $X^* = x - iy$ and $P_x^* = p_x - ip_y$
• $XX^* = x^2 + y^2$
• $P_x P_x^* = p_x^2 + p_y^2$
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After symplification we get:

$$\left(\frac{eB_2}{2}\right)^2 XX^* + P_X P_X^* = \left(\frac{eB_2}{2}\right)^2 (x^2 + y^2) + (p_x^2 + p_y^2) = \text{Cte} \qquad (9)$$

We can get :

$$XX^* + \left(\frac{2}{eB_2}\right)^2 P_X P_X^* = \operatorname{Cte}.$$

• The acceptance condition :

$$XX^* \leq a^2$$

a is the aperture at the end of the long solenoidWe can write :



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QWT Volume acceptance

The volume acceptance V of the QWT is then defined as:

$$V_{QWT} = \int_{XX^* \le a^2} dx \, dy \, dp_x \, dp_y$$

and can be written :

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$$V(\chi_1) = \frac{2\pi^2}{3} \left(\frac{eB_2 a^2}{2}\right)^2 \left[1 - \left(1 - \frac{1}{\sin^2 \chi_1 + \left(\frac{B_1}{B_2}\right)^2 \cos^2 \chi_1}\right)^{\frac{3}{2}}\right]$$

•

Volume acceptance

The momentum p_m at maximum transmission is obtained from the equation

$$\frac{dV(\chi_1)}{d\chi_1} = 0 \tag{10}$$

Will get:

$$\chi_1 = \pi/2 \tag{11}$$

Which leads to:

$$p_m = \frac{eB_1L_1}{\pi} \tag{12}$$

• Correlation between the length of the high field and the magnetic field intensity.

QWT characteristique

The accepted momentum dispersion in the QWT

• At the maximum volume acceptance:

$$\frac{\delta p}{p} = \frac{4}{\pi} \, \frac{B_2}{B_1}$$

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Transverse acceptance

• Replacing the rotation angle $\chi_1 = \pi/2$ which lead to maximize the volume acceptance:

$$XX^{*} + \left[\frac{2}{eB_{2}}\right]^{2} = \left[\frac{B_{1}}{B_{2}}\right]^{2} X_{0} X_{0}^{*} + \left[\frac{2}{eB_{1}}\right]^{2} P_{X_{0}} P_{X_{0}^{*}} = \text{Cte}$$
(13)

• Cylindrical coordinates:

$$\left[\frac{B_1}{B_2}\right]^2 r_0^2 + \left[\frac{2}{eB_1}\right]^2 \left[P_{r_0}^2 + \frac{P_{\phi_0}}{r^2}\right] = \text{Cte}$$
(14)

With

$$\operatorname{Cte} - \left(\frac{2}{eB_2}\right)^2 P_X P_X^* \le a^2.$$
(15)

• where $(r_0, P_{r_0}, P_{\phi_0})$ are the coordinates of the particles at target.

$$\begin{bmatrix} \frac{B_1}{B_2} \end{bmatrix}^2 r_0^2 + \begin{bmatrix} \frac{2}{eB_1} \end{bmatrix}^2 \begin{bmatrix} P_{r_0}^2 + \frac{P_{\phi_0}^2}{r^2} \end{bmatrix} = a^2 + \begin{bmatrix} \frac{2}{eB_2} \end{bmatrix}^2 \frac{P_{\phi_0}^2}{a^2} \tag{16}$$

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Transverse acceptance

• The maximum extension in the spacial demension for an ellipsoid equation is deteminded if:

$$P_{r_0}=P_{\phi_0}=0$$

• Providing the radial accepptance of the QWT:

$$r_0^{max} = \frac{B_2}{B_1} a.$$
 (18)

• Example: Considering the case (B₁, B₂, a)=(2.5 T,0.5 T,20 mm), we obtain

$$r_0^{max} = 4 \text{ mm}$$
 (19)
 $P_{r_0}^{max} = 5.7 \text{ MeV}/c.$ (20)

Transverse acceptance

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Adiabatic matching device

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AMD Volume acceptance

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• The accepted phase space under AMD magnetic field is :

$$XX^* + \left[\frac{2}{eB_2}\right]^2 P_X P_X^* = \frac{B_1}{B_2} X_0 X_0^* + \frac{4}{e^2 B_1 B_2} P_{X_0} P_{X_0}^* = \text{Cte}.$$
 (21)

• The corresponding volume acceptance for such a device :

$$V_{AMD} = \int_{XX^* \le a^2} dx \, dy \, dp_x \, dp_y = \frac{2\pi^2}{3} \left[\frac{eB_2 a^2}{2}\right]^2.$$
(22)

 Where B₂ is the end value of the magnetic field at the end of the AMD

$$B(z) = \frac{B_1}{1 + \mu z} \tag{23}$$

• where P_0 is a particular central scalar momentum. The value of the smallness parameter

$$\epsilon = \frac{P_0}{B_2} \frac{dB}{dz} \tag{25}$$

• is a characteristic constant of the adiabatic operation ($\epsilon \ll 1$). From

$$\int_{0}^{L_{1}} \frac{\epsilon}{P_{0}} dz = \int_{B_{1}}^{B_{2}} \frac{dB}{B_{2}}$$
(26)

• We obtain
$$B(z) = \frac{B_1 B_2 L_1}{B_2 L_1 + (B_1 - B_2) z} \,. \tag{27}$$

AMD Volume acceptance

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• Cylindrical coordinates:

$$[\frac{B_1}{B_2}](\frac{r_0}{a})^2 + (\frac{P_{r0}}{\frac{1}{2}e\sqrt{B_1B_2a}})^2 + (\frac{P_{\phi 0}}{\frac{1}{2}}eB_2a^2)^2[\frac{B_1}{B_2}\frac{1}{[\frac{r_0}{a}]^2} - 1] \le 1 \quad (28)$$

• Making $P_{r0} = 0$ and $P_{\phi 0} = 0$ in (28), we get the radial acceptance for the AMD:

$$r_0^{max} = \sqrt{\frac{B_2}{B_1}}a \tag{29}$$

Making $r_0 = 0$ and $P_{\phi 0} = 0$ in (28), we get the transverse momentum acceptance for the AMD::

$$p_{x0}^{max} = \frac{1}{2}e \ a\sqrt{B_2 B_1}$$
(30)

• For $B_1 = 2.5$ T and $B_2 = 0.5$ T with $\mu = 40$ m^{-1} and L = 10 cm the acceptance caracteristique are :

$$r_0^{max} = \sqrt{\frac{B_2}{B_1}} a = \sqrt{\frac{0.5}{2.5}} \times 20 mm \approx 8.9 mm$$
 (31)

$$p_{x0}^{max} = \frac{1}{2} ea \sqrt{B_2 B_1} = \frac{1}{2} e \times 20 mm \sqrt{2.5 \times 0.5} \le 3.34 \frac{MeV}{c}$$
 (32)

Transverse acceptance AMD

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- With a suitable length for the QWT we got a momentum acceptance which is much larger than the momentum acceptance in the AMD. Moreover, the wide momentum acceptance allow us to use a small radius aperture or large angle for the positron distribution at the target.
- The radial acceptance in AMD is larger than the radial acceptance in QWT.
- The volume of the phase space of the AMD is independent of the longitudinal momentum.

