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Analyzing electron beam using BPM from “Lecture Notes on Topics in Accelerator Physics, Section 4: Fast Ion Instability”  
by Alex Chao, November 2002 SLAC-PUB-9574

Assume we have a beam position monitor (BPM) located at  $s = 0$  along the accelerator. Let an electron bunch of length  $l$  circulate in a storage ring of circumference  $C = cT_0$ . Since the beam is bunched, the BPM will measure an alternating current signal via pick-up electrodes. Based on the amplitudes of the signals on the electrodes, the position of the beam can be determined. The time-dependent signal seen by the BPM is

$$\text{signal}(t) = \sum_{k=0}^{\infty} y_e(kC|ct - kC)|_{0 < ct - kC < l}$$

where  $y_e$  is the transverse distance of an electron from the beam centroid and  $k$  sums over multiple turns. We can take a Fourier transform of the BPM signal into frequency space:

$$\begin{aligned} \text{spectrum}(\Omega) &\propto \int_0^{\infty} dt e^{-i\Omega t} \text{signal}(t) \\ &= \sum_{k=0}^{\infty} \int_0^{l/c} dt' e^{-i\Omega(t'+kT_0)} y_e(kC|ct') \\ &= \sum_{k=0}^{\infty} e^{-i(\Omega+\omega_\beta)kT_0} \int_0^{l/c} dt' e^{-i(\Omega-\omega_I)t'} \tilde{y}_e(kC|ct') \\ &= y_0 \sum_{k=0}^{\infty} e^{-i(\Omega+\omega_\beta)kT_0} \int_0^{l/c} dt' e^{-i(\Omega-\omega_I)t'} \frac{e^{\eta'}}{\sqrt{2\pi\eta'}} \end{aligned}$$

where  $\eta' = t' \sqrt{K\omega_I kC/2\omega_\beta l}$ . From the second to third step, the form of  $y_e$  was used:

$$y_e(s|z) = \tilde{y}_e(s|z) e^{-i\omega_\beta s/c + i\omega_I z/c}$$

From the third to the fourth step, the form of  $\tilde{y}_e$  in the asymptotic regime ( $\eta \gg 1$ ) was used:

$$\tilde{y}_e(s|z) = y_0 I_0(\eta) \approx y_0 \frac{e^\eta}{\sqrt{2\pi\eta}}$$

The integral in the last step is of the form:

$$I = \int_0^{l/c} dt' \frac{e^{(B-iA)t'}}{\sqrt{2\pi B t'}} = \frac{1}{\sqrt{2\pi B}} (-B + iA)^{-\frac{1}{2}} \gamma\left(\frac{1}{2}, (-B + iA) \frac{l}{c}\right)$$

where  $A = \Omega - \omega_I$  and  $B = \sqrt{K\omega_I kC/2\omega_\beta l}$  and  $\gamma(\alpha, x)$  is the lower incomplete Gamma function:

$$\gamma(\alpha, x) = \int_0^x t^{\alpha-1} e^{-t} dt$$

For  $|x| \gg 1$ ,  $\gamma(\alpha, x) \approx -x^{\alpha-1} e^{-x}$ . Thus, with  $|A|l/c \gg Bl/c \gg 1$  (from the validity criterion  $\frac{\omega_I l}{c} \gg \eta \gg 1$  with

$$\eta = \frac{z}{c} \sqrt{\frac{K\omega_I s}{2\omega_\beta l}},$$

$$I \approx \sqrt{\frac{l/c}{2\pi B}} e^{Bl/c} \left( \frac{e^{-iAl/2c} \sin \frac{Al}{2c}}{Al/2c} \right)$$

We can then plug this into the equation for spectrum ( $\Omega$ ). If we measure the signal in a small window around a large  $k = \bar{k}$  (note that the signal obviously diverges for  $k \rightarrow \infty$ ), we have

$$|\text{spectrum}(\Omega)| \propto y_0 \sqrt{\frac{l/c}{2\pi\bar{B}}} e^{\bar{B}l/c} \left| \frac{\sin\left(\frac{\Omega - \omega_I}{2c}l\right)}{(\Omega - \omega_I)l/2c} \right| \sum_{p=-\infty}^{\infty} \delta(\Omega + \omega_\beta - p\omega_0)$$

where  $\omega_0 = 2\pi/T_0$  is the revolution angular frequency,  $\bar{B} = \sqrt{K\omega_I\bar{k}C/2\omega_\beta l}$  (similar to  $\eta$  above), and we have plugged in  $A = \Omega - \omega_I$ . We see that the electron beam spectrum contains  $\delta$ -function peaks at  $\Omega = p\omega_0 - \omega_\beta$  corresponding to the lower betatron sidebands of all revolution harmonics. It also contains a broad envelope  $\frac{\sin[(\Omega - \omega_I)l/2c]}{(\Omega - \omega_I)l/2c}$  around  $\Omega = \omega_I$ , the characteristic ion frequency, with width  $\Delta\Omega \pm \pi c/l$ . Thus, for longer bunches, the width of the envelope decreases and becomes more defined. The entire spectrum also grows with time according to the factor  $e^{\bar{B}l/c}/\sqrt{2\pi\bar{B}l/c}$ , as we would expect.

To see what this spectrum looks like, we can plot  $|\text{spectrum}(\Omega)|$  as a function of  $\Omega$  for several cases. Let's plug in some reasonable numbers:  $\omega_\beta = 5\text{MHz}$ ,  $K = \frac{4\Sigma nNc^2r_e}{\gamma a^2}$ ,  $6.5 \times 10^7\text{s}^{-2}$ ,  $C = cT_0 = \frac{2\pi c}{\omega_0} = 2000\text{m}$ ,  $T_0 = C/c = 6.7\mu\text{s}$ , and  $\omega_0 \approx 0.942\text{MHz}$  (the revolution angular frequency). We'll make plots of  $\frac{|\text{spectrum}(\Omega)|}{y_0}$  vs.  $\Omega$  for various values for  $\bar{k}$ ,  $l$ , and  $\omega_I$ . We'll consider the characteristic frequencies  $\omega_I = \sqrt{\frac{2Nr_p c^2}{la^2 A}}$  for  $\text{H}_2^+$ ,  $\text{CH}_4^+$ ,  $\text{N}_2^+$ , and  $\text{CO}_2^+$  with  $N = 10^{11}$  electrons,  $a = 1\text{mm}$ ,  $r_p = 1.54 \times 10^{-16}\text{cm}$ , and  $A = \frac{M}{m_p}$  ( $A_{\text{H}_2^+} = 2$ ,  $A_{\text{CH}_4^+} = 16$ ,  $A_{\text{N}_2^+} = 28$ ,  $A_{\text{CO}_2^+} = 44$ ). Below is a table of parameters/calculated values for each plot:

| Plot | $\bar{k}$ | $\omega_I(\text{Hz})$                  | $l(\text{m})$ | $\bar{B}(\text{Hz})$ |
|------|-----------|--|---------------|----------------------|
| 1    | $10^1$    | $1.18 \times 10^8$ ( $\text{H}_2^+$ )  | 1             | $3.91 \times 10^6$   |
| 2    | $10^1$    | $4.16 \times 10^7$ ( $\text{CH}_4^+$ ) | 1             | $2.32 \times 10^6$   |
| 3    | $10^1$    | $3.15 \times 10^7$ ( $\text{N}_2^+$ )  | 1             | $2.02 \times 10^6$   |
| 4    | $10^1$    | $2.51 \times 10^7$ ( $\text{CO}_2^+$ ) | 1             | $1.81 \times 10^6$   |
| 5    | $10^2$    | $1.18 \times 10^8$ ( $\text{H}_2^+$ )  | 1             | $1.24 \times 10^7$   |
| 6    | $10^4$    | $1.18 \times 10^8$ ( $\text{H}_2^+$ )  | 1             | $1.24 \times 10^8$   |
| 7    | $10^1$    | $3.72 \times 10^7$ ( $\text{H}_2^+$ )  | 10            | $6.96 \times 10^5$   |
| 8    | $10^1$    | $1.18 \times 10^7$ ( $\text{H}_2^+$ )  | 100           | $1.24 \times 10^5$   |
| 9    | $10^1$    | $4.16 \times 10^6$ ( $\text{CH}_4^+$ ) | 100           | $7.36 \times 10^4$   |
| 10   | $10^1$    | $3.15 \times 10^6$ ( $\text{N}_2^+$ )  | 100           | $6.40 \times 10^4$   |
| 11   | $10^1$    | $2.51 \times 10^6$ ( $\text{CO}_2^+$ ) | 100           | $5.71 \times 10^4$   |
| 12   | $10^1$    | All Four $\omega_I$                    | 100           | N/A                  |

Table 1: Calculated values for  $\bar{B}$  for various  $\bar{k}$

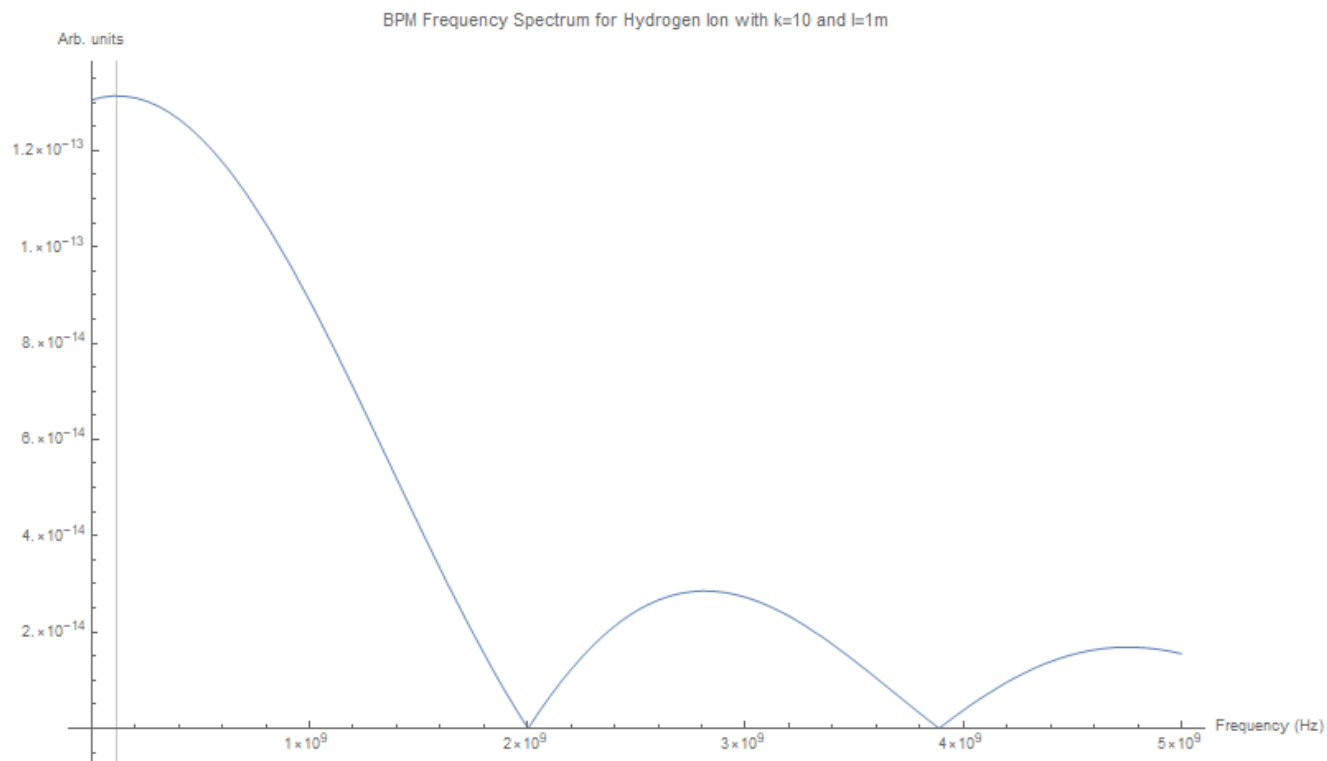


Figure 1: BPM Frequency Spectrum 1:  $H_2^+$  with  $k = 10$ , and  $l = 1m$

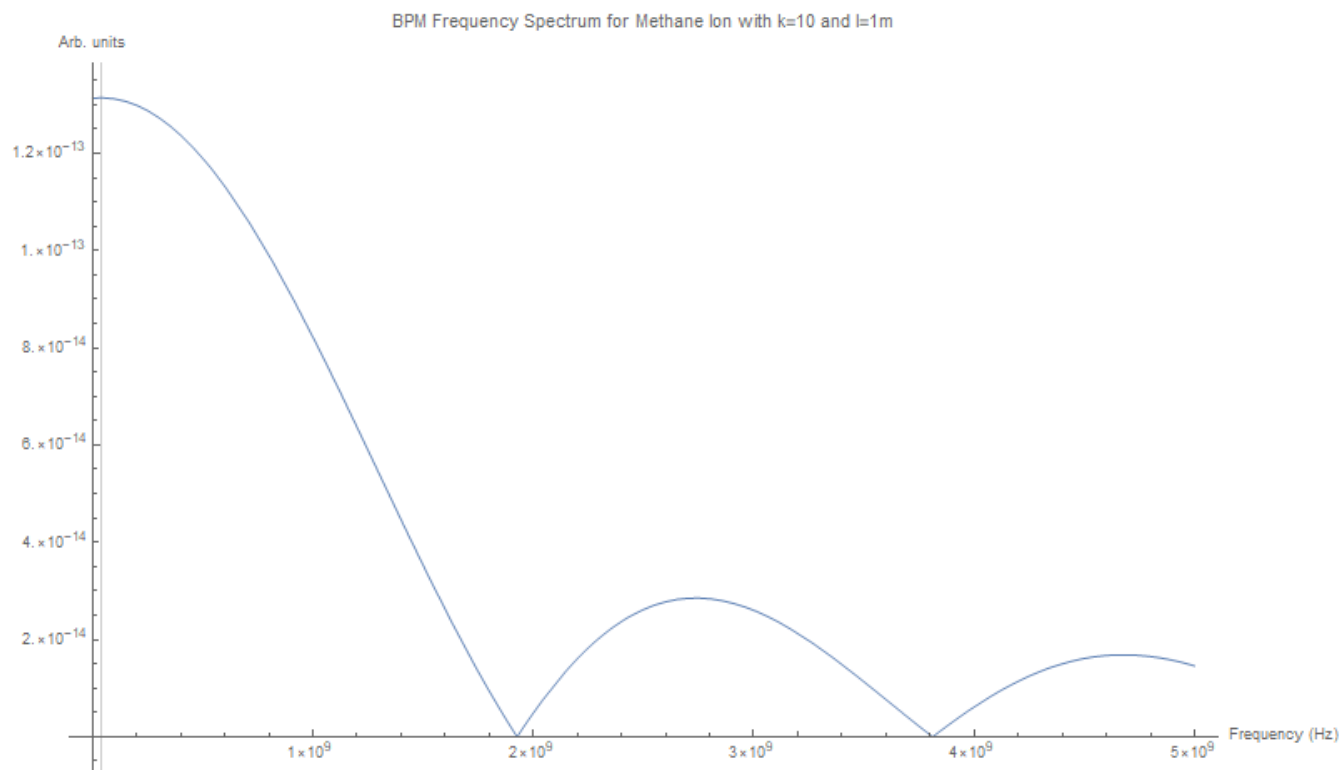


Figure 2: BPM Frequency Spectrum 2:  $CH_4^+$  with  $k = 10$  and  $l = 1m$

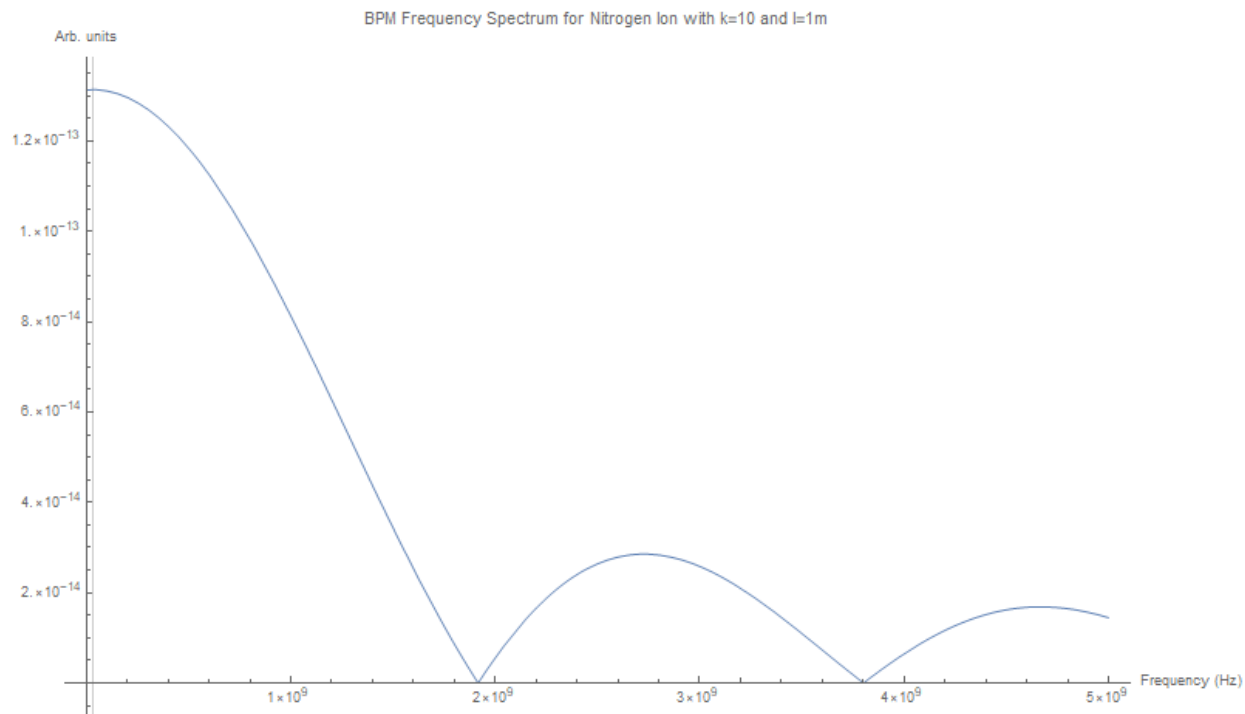


Figure 3: BPM Frequency Spectrum 3:  $N_2^+$  with  $k = 10$  and  $l = 1m$

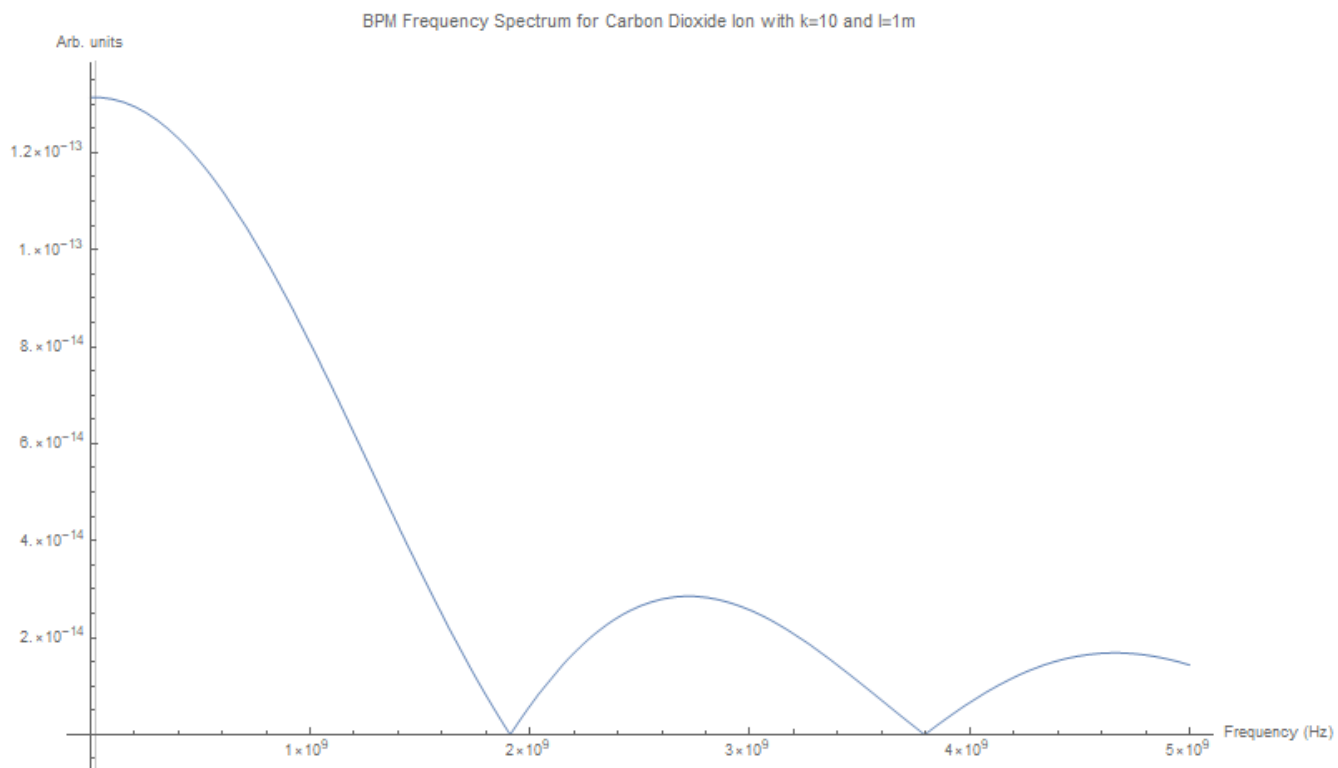


Figure 4: BPM Frequency Spectrum 4:  $CO_2^+$  with  $k = 10$  and  $l = 1m$

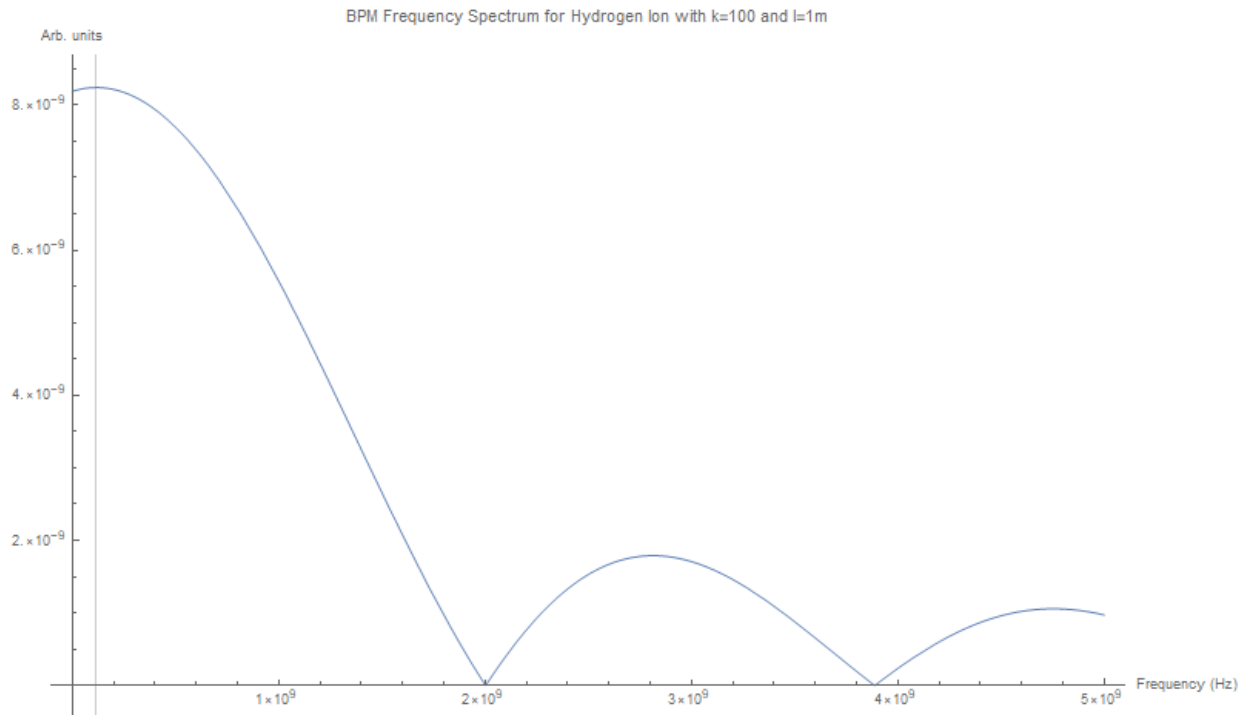


Figure 5: BPM Frequency Spectrum 5:  $H_2^+$  with  $k = 100$  and  $l = 1m$

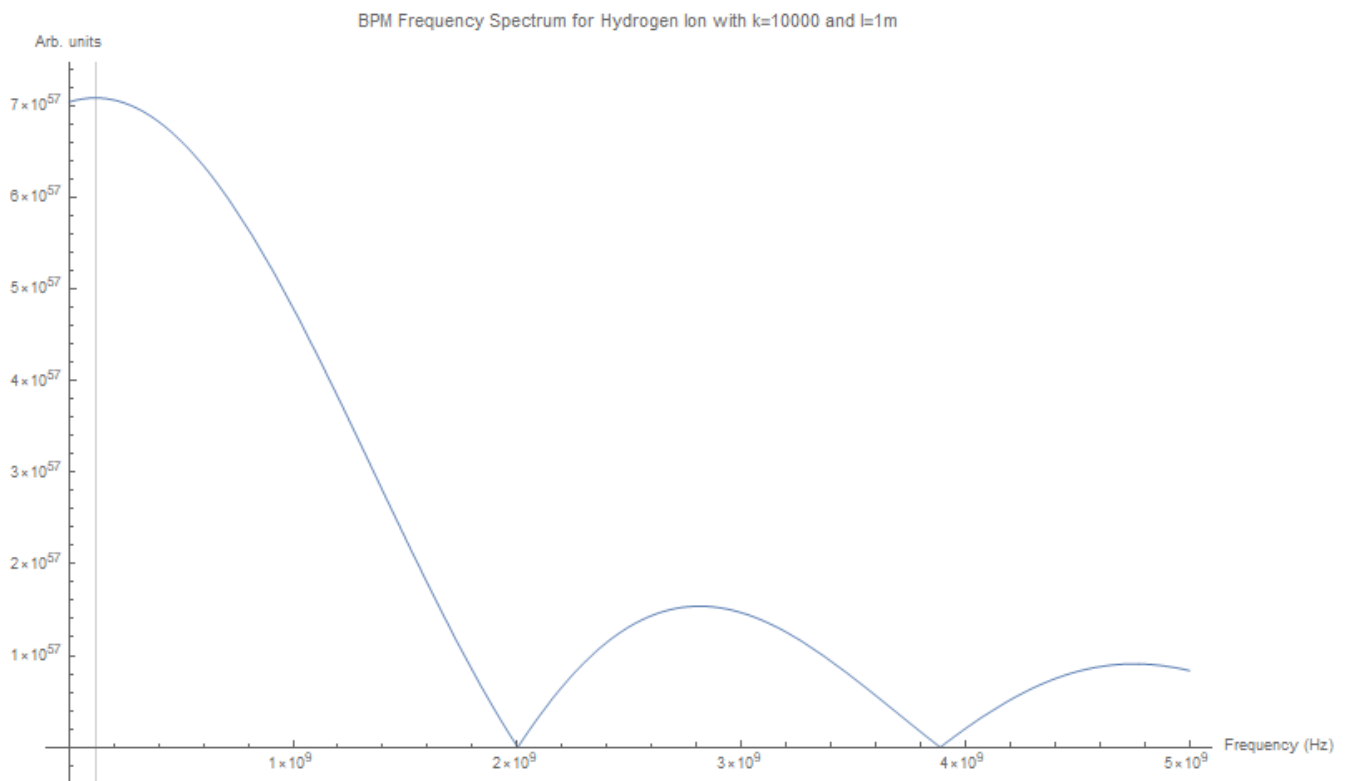


Figure 6: BPM Frequency Spectrum 6:  $H_2^+$  with  $k = 10000$  and  $l = 1m$

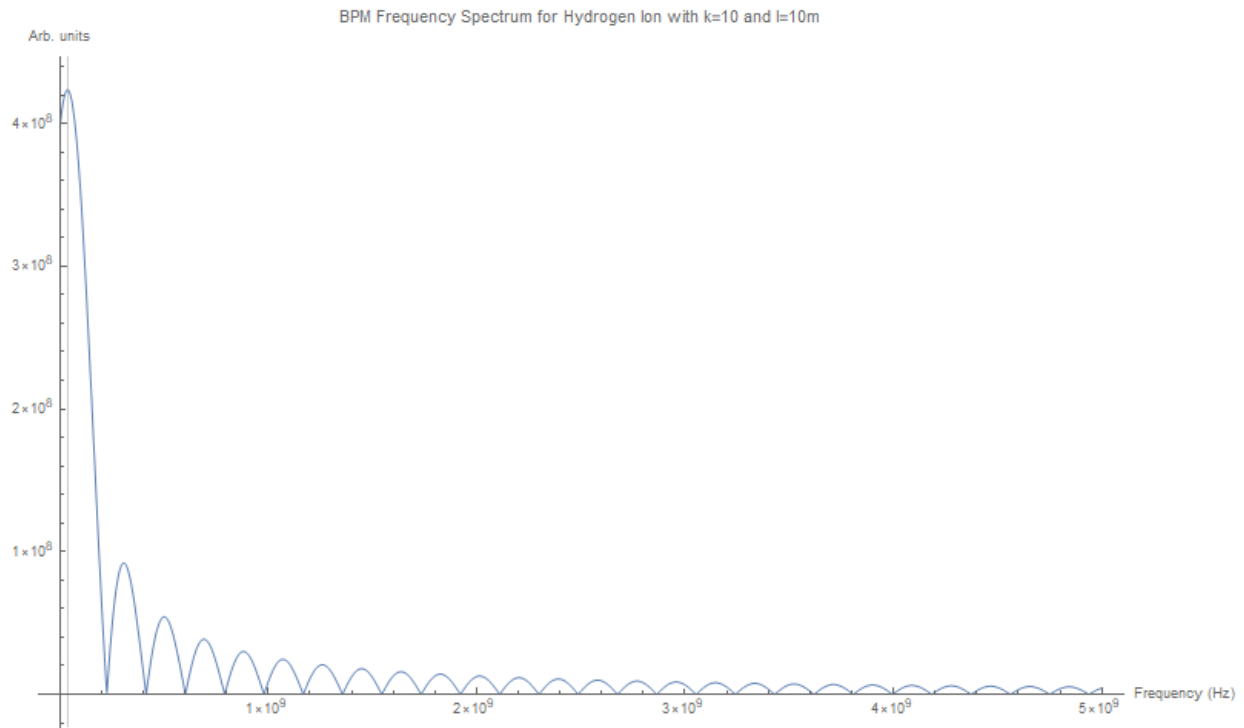


Figure 7: BPM Frequency Spectrum 7:  $H_2^+$  with  $k = 10$  and  $l = 10m$

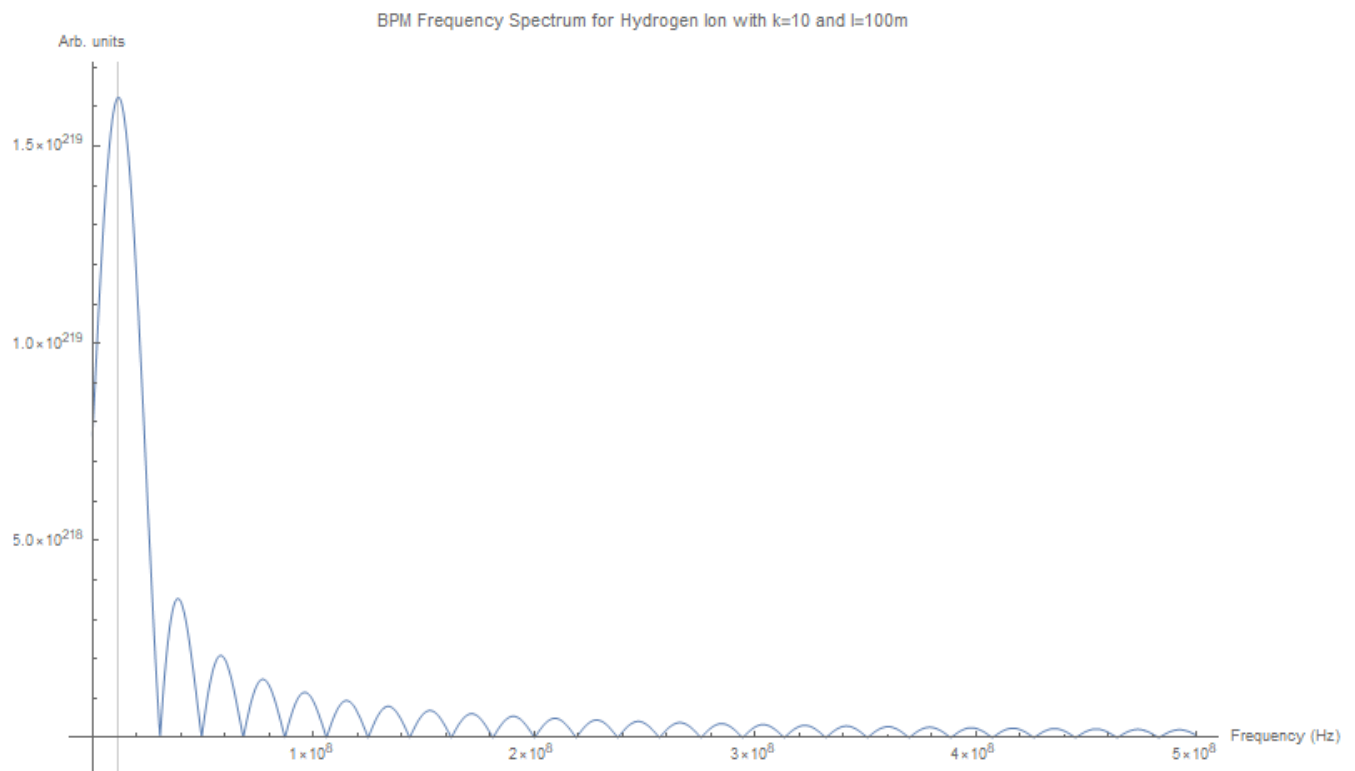


Figure 8: BPM Frequency Spectrum 8:  $H_2^+$  with  $k = 10$ ,  $l = 100m$

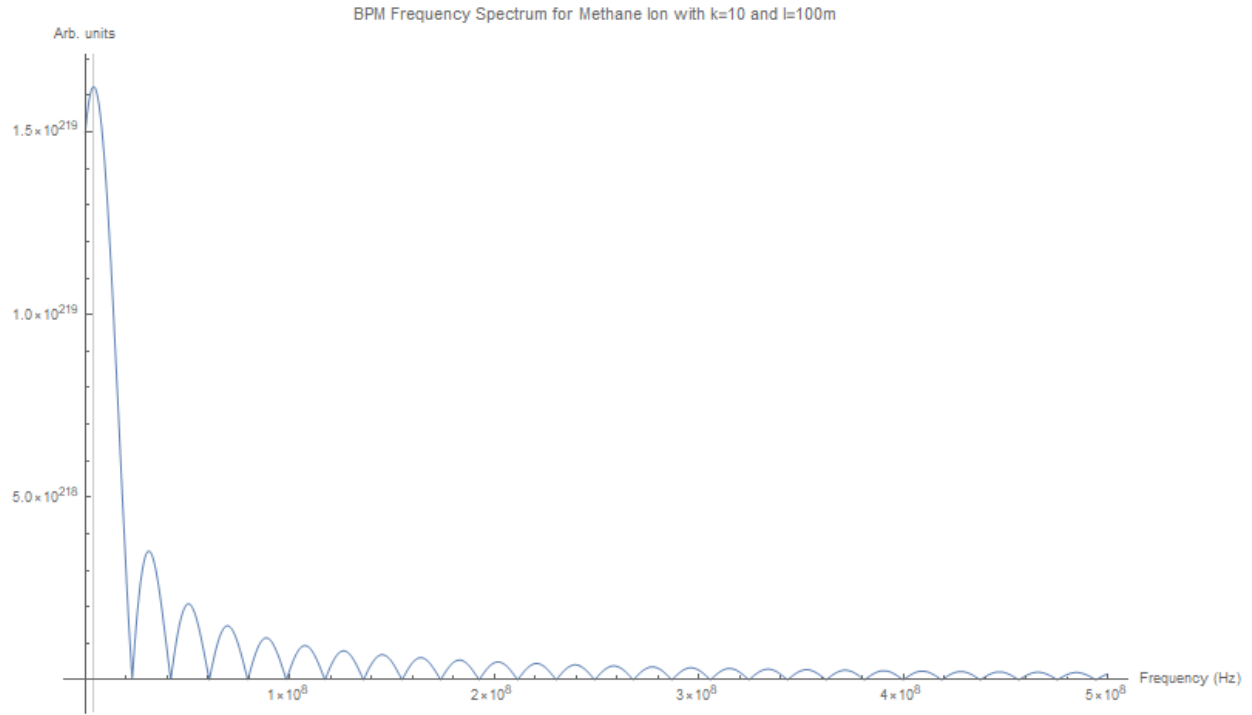


Figure 9: BPM Frequency Spectrum 9:  $\text{CH}_4^+$  with  $k = 10$  and  $l = 100m$

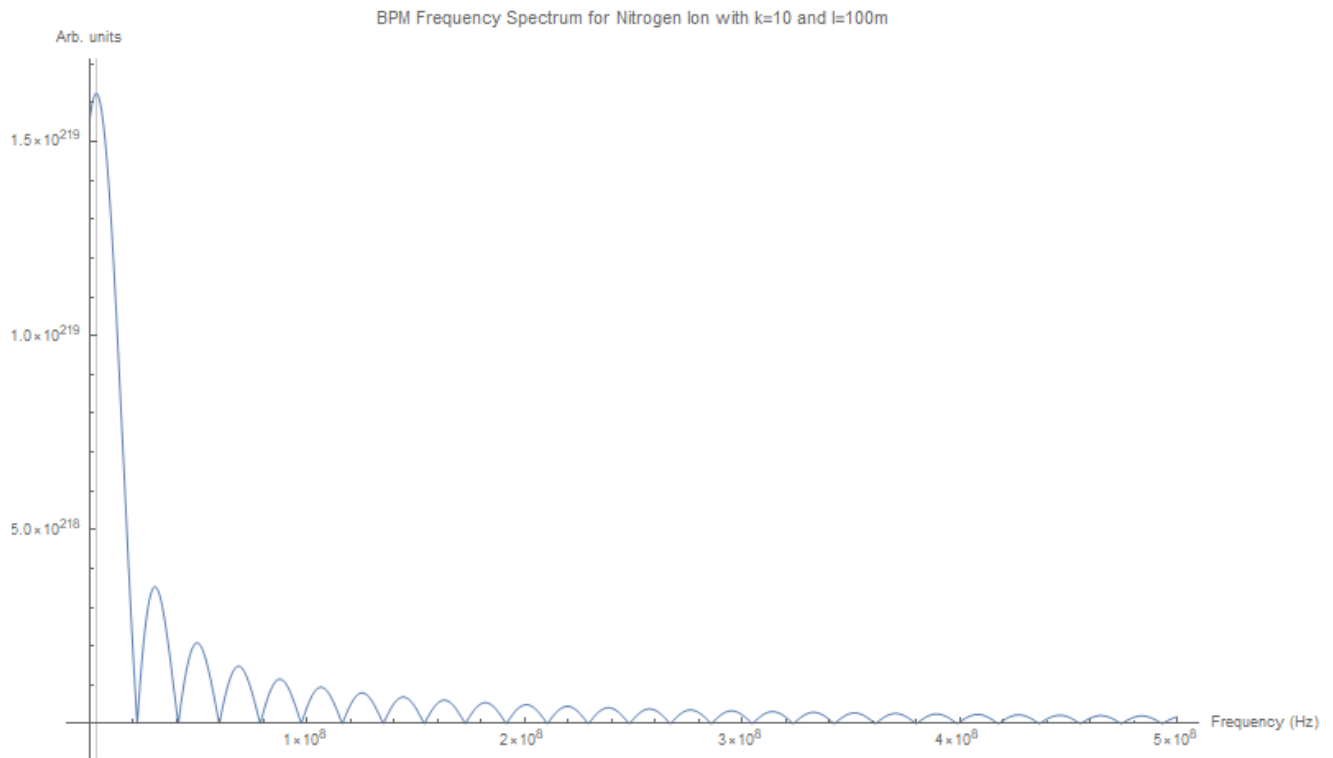


Figure 10: BPM Frequency Spectrum 10:  $\text{N}_2^+$  with  $k = 10$  and  $l = 100m$

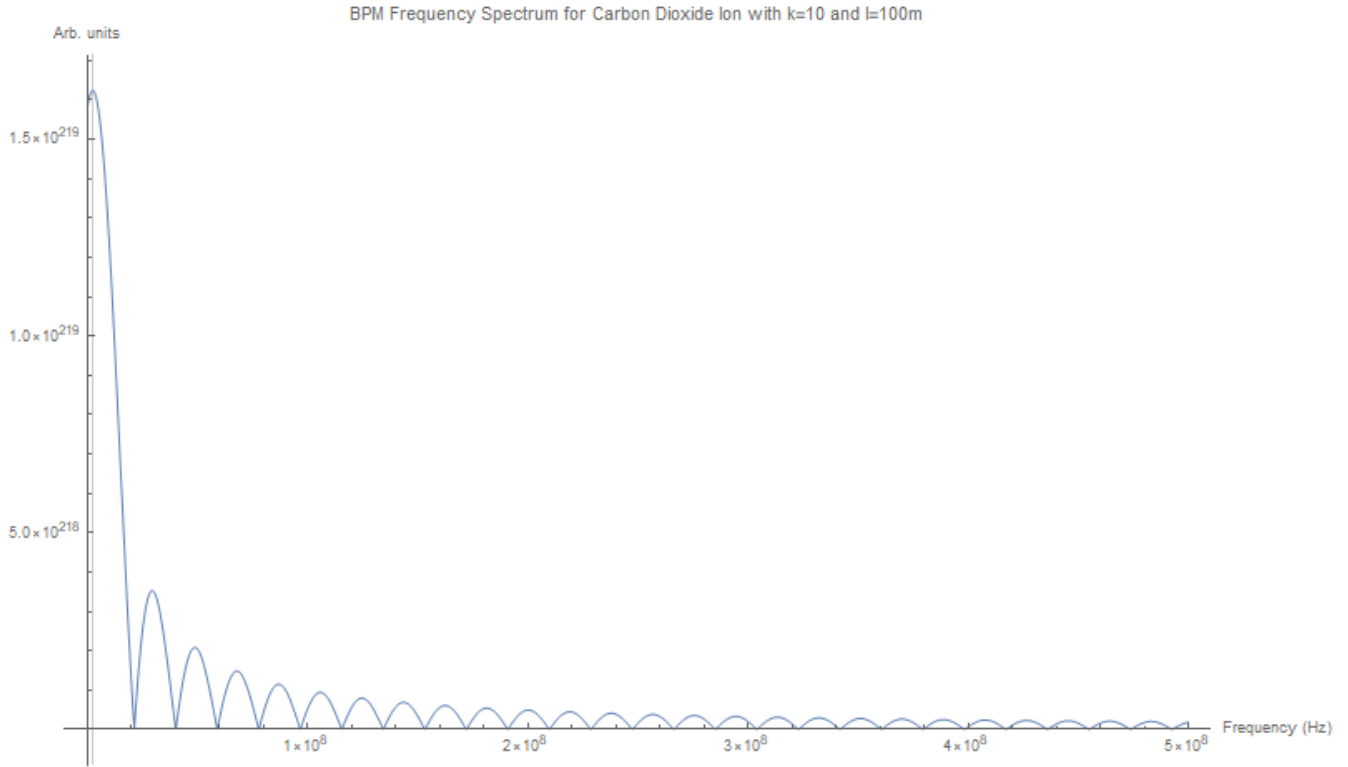


Figure 11: BPM Frequency Spectrum 11:  $\text{CO}_2^+$  with  $k = 10$  and  $l = 100m$

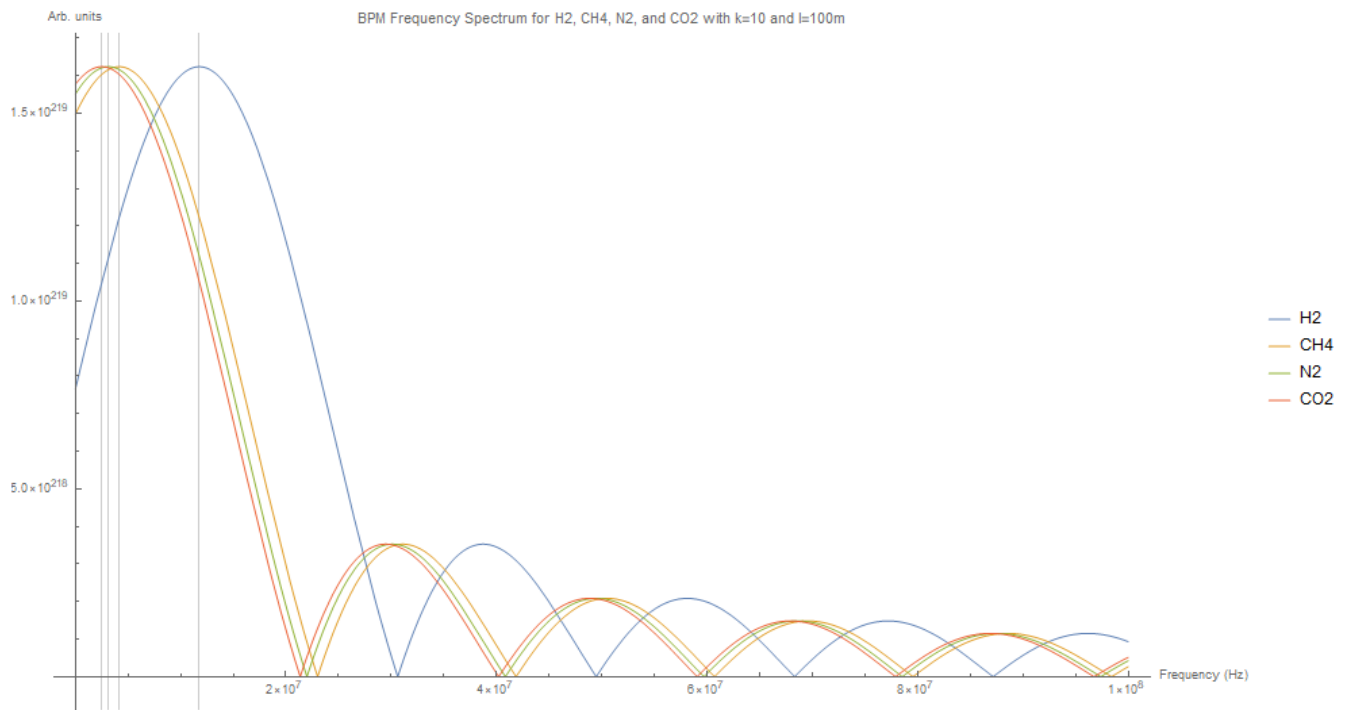


Figure 12: BPM Frequency Spectrum 12:  $\text{H}_2^+$ ,  $\text{CH}_4^+$ ,  $\text{N}_2^+$  and  $\text{CO}_2^+$  with  $k = 10$  and  $l = 100m$