

Direct measurement of the bubble-nucleation energy threshold in a CF₃I bubble chamber

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We have directly measured the energy threshold and efficiency for bubble nucleation from iodine recoils in a CF₃I bubble chamber in the energy range of interest for a dark matter search. These interactions cannot be probed by standard neutron calibration methods, so we develop a new technique by observing the elastic scattering of 12 GeV/*c* negative pions. The pions are tracked with a silicon pixel telescope and the reconstructed scattering angle provides a measure of the nuclear recoil kinetic energy. The bubble chamber was operated with a nominal threshold of (13.6 ± 0.6) keV. Interpretation of the results depends on the response to fluorine and carbon recoils, but in general we find agreement with the predictions of the classical bubble-nucleation theory. This measurement confirms the applicability of CF₃I as a target for spin-independent dark matter interactions and represents a novel technique for calibration of superheated fluid detectors.

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Recent years have seen a resurgence in the use of superheated liquids and bubble chambers as continuously sensitive nuclear recoil detectors searching for dark matter in the form of weakly interacting massive particles (WIMPs) [1–3]. At a low degree of superheat, bubble chambers are insensitive to minimum ionizing backgrounds that normally plague WIMP searches but retain sensitivity to the nuclear recoils that would be characteristic of WIMP scattering. In a superheated liquid the process of radiation-induced bubble nucleation is described by the classical “hot spike” model [4]. For the phase transition to occur, the energy deposited by the particle must create a critically sized bubble, requiring a minimum energy deposition in a volume smaller than the critical bubble. Under mildly superheated conditions, the latter requirement renders the bubble chamber insensitive to minimum ionizing particles.

The radius of the critical bubble is given by the condition that the bubble be in (unstable) equilibrium with the surrounding superheated fluid [5]. This demands the pressure balance

$$P_b - P_l = \frac{2\sigma}{r_c}, \quad (1)$$

where P_b is the pressure inside the bubble, P_l is the pressure in the liquid, σ is the bubble surface tension,

and r_c is the critical bubble radius. The pressure P_b is fixed by the condition that the chemical potential inside and outside the bubble be equal, giving

$$(P_b - P_l) \approx (P_{\text{sat}} - P_l) \frac{\rho_l - \rho_v}{\rho_l}, \quad (2)$$

where P_{sat} is the pressure in a saturated system at the given temperature and ρ_l and ρ_v are the liquid and vapor densities, respectively, in the saturated system [6].

In Seitz’s hot spike model for bubble nucleation, the entire energy necessary to create the critical bubble must come from the particle interaction that nucleates the bubble. This is in contrast to earlier models that required only the work (free energy) to come from the particle interaction, with the remaining bubble-formation energy supplied by heat flowing in from the surrounding superheated fluid [7]. As the name hot spike implies, the nucleation site in Seitz’s model begins as a high-temperature seed, so it cannot draw heat from the surrounding fluid.

Once the decision is made to consider the total bubble creation energy rather than just the free energy, the threshold energy calculation is completely described by Gibbs [5]. This energy is given by

$$E_T = 4\pi r_c^2 \left(\sigma - T \frac{\partial \sigma}{\partial T} \right) + \frac{4\pi}{3} r_c^3 \rho_b (h_b - h_l) - \frac{4\pi}{3} r_c^3 (P_b - P_l) + \mathcal{O}\left(\frac{\delta}{r_c}\right). \quad (3)$$

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Here, T is the temperature of the system, ρ_b is the bubble vapor density, and h_b and h_l are the specific enthalpies of the bubble vapor and superheated liquid, respectively. The surface tension σ and temperature derivative are taken along the usual saturation curve. The three terms give, from left to right, the heat necessary to create the bubble surface, the heat needed to vaporize the fluid to make the bubble interior, and a reversible work term done in expanding the bubble to the critical size that must be subtracted to avoid double-counting work present in both of the first two terms. To good approximation $h_b - h_l$ may be replaced by the normal heat of vaporization at temperature T .

The greatest uncertainty in determining the thermodynamic E_T is the relation between the surface tension at a flat liquid-vapor interface and the surface tension for a very small bubble. This relation is described by the ‘‘Tolman length’’ δ , which is unknown but is expected to be some fraction of the intermolecular distance [8]. This translates to an uncertainty on E_T of $\sim 3\%$. For the rest of this paper, we refer to the calculated threshold in Eq. (3) as the Seitz threshold.

The Seitz model assumes the efficiency for bubble nucleation is 100% for all interactions that deposit $E \geq E_T$ over a volume small compared to the critical bubble. The length scales for nuclear recoil cascades in the energy region between 5 and 20 keV relevant for a WIMP search are similar to the critical radius, so the Seitz model may or may not give a good description of bubble nucleation, and direct calibrations of bubble-nucleation efficiency are necessary.

The working fluid discussed in this paper is iodotri-fluoromethane or CF_3I , which contains two highly sensitive WIMP target nuclei: fluorine, for spin-dependent interactions, and iodine, for spin-independent interactions. Neutrons are typically used to mimic WIMPs in calibrating the nuclear recoil response of a WIMP detector, and neutron sources have been used to measure the nucleation threshold for carbon and fluorine recoils in CF_3I , CF_3Br [1,9] and C_4F_{10} [10] under various superheat conditions. However, iodine recoils contribute only a small fraction to the total neutron-nucleated bubble rate in CF_3I . Therefore, neutron sources are an ineffective calibration tool for iodine recoils in COUPP. We have used heavy daughter nuclei produced in alpha decays as a proxy [11], but these are high energy recoils of ~ 100 keV. This paper describes a measurement of bubble-nucleation efficiency for iodine recoils near our dark matter search thresholds.

I. EXPERIMENTAL DESIGN

Our bubble chambers are insensitive to minimum ionizing particles, allowing us to exploit a new calibration technique using charged pions as WIMP surrogates to produce nuclear recoils by strong elastic scattering. We measure the pion scattering angle using silicon pixel detectors. The nuclear recoil kinetic energy can be calculated

by $E = (p\theta)^2/2M$ on an event by event basis, where p is the beam momentum, θ the scattering angle, and M the nuclear mass of the target.

For a CF_3I target, a measured scattering angle corresponds to a different recoil energy depending on which nucleus is involved in the interaction; in this paper, we will refer to iodine equivalent recoil energy E_{Ie} as the energy given to an iodine nucleus for a given pion scattering angle. For a 12 GeV/ c pion beam, approximately 75% of the rate of pions scattering into angles corresponding to E_{Ie} between 5 and 20 keV is due to elastic scattering on iodine, with smaller contributions from carbon, fluorine, and inelastic scattering [12]. Therefore, the bubble-nucleation efficiency for iodine recoils in a bubble chamber with Seitz threshold between 5 and 20 keV can be inferred from a measurement of the fraction of pion-scattering events that nucleate bubbles in the chamber as a function of E_{Ie} .

The measurement was performed in the Fermilab Test Beam Facility [13] using a 12 GeV/ c mainly π^- beam with $\sigma_p/p = 4\%$ and an angular spread of < 1.1 mrad. The absolute momentum of the beam is known to 3%. The pions were tracked with a silicon pixel telescope [14] consisting of four upstream and four downstream silicon pixel plaquettes, with a spatial coverage of $14 \text{ mm} \times 14 \text{ mm}$. The total length of the telescope was 90 cm. The angular resolution was 0.6 mrad (σ) in the horizontal (x) direction and 0.7 mrad in the vertical (y) direction, with roughly equal contributions from multiple Coulomb scattering (MCS) in the target and the spatial resolution of the telescope. Plastic scintillators triggered the pixel telescope on each beam particle.

A small bubble chamber was designed for this measurement consisting of a quartz test tube with inner diameter 10 mm and 1-mm-thick wall, filled with 7 cm^3 of CF_3I . The small size is required to minimize MCS in the short radiation length of CF_3I (58 mm). The bubble chamber was operated at a pressure of 30.0 ± 0.1 psi (absolute) and a temperature of 34.2 ± 0.2 °C with a nominal Seitz threshold of (13.6 ± 0.6) keV. The iodine equivalent threshold scattering angle is 4.7 mrad. An acoustic transducer was attached to the top of the test tube to record the acoustic emission produced during bubble formation, providing the time of bubble nucleation with $\sim 10 \mu\text{s}$ resolution. Temperature control was provided by a water bath around the bubble chamber.

Bubble chamber data were taken between March 14 and March 28, 2012, with a beam flux of ~ 1000 particles per 4-second beam spill with one spill per minute. The size of the beam spot was wider than both the bubble chamber and the pixel telescope. The chamber was expanded to the superheated state 22 seconds before the arrival of the beam, allowing time for pressure and temperature transients to dissipate after expansion. The observation of bubbles by a 100-Hz video camera system created a bubble

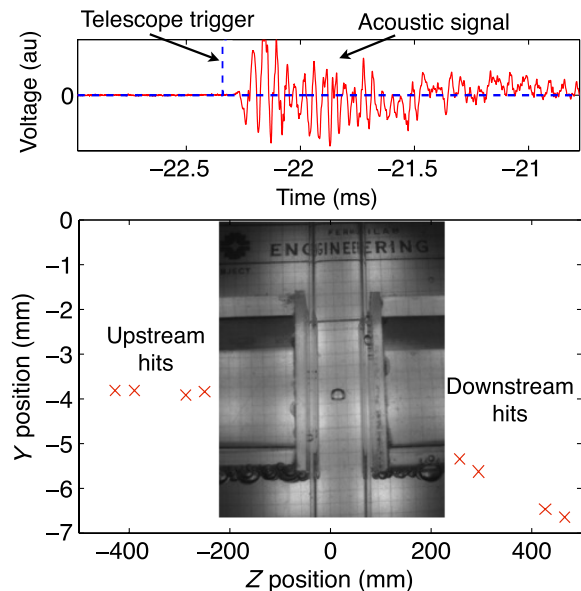


FIG. 1 (color online). An example event ($\theta = 6$ mrad), including the relative timing of the telescope trigger and acoustic signal, one camera image of the bubble, and the y and z positions of the telescope hits. The pion beam is in the $+z$ direction. The camera image is not to scale but the test tube has an inner diameter of 10 mm. Beam tubes in the water bath to minimize the material traversed by the pion beam are visible to either side of the bubble chamber.

chamber trigger, causing the video images and associated data to be recorded and the chamber to be recompressed. After recompression, the chamber was dead for the remainder of the beam spill, allowing us to collect at most one bubble event per minute. We collected about four good single-bubble events per hour, with the primary losses due to premature bubble chamber triggers, bubbles forming outside of the region covered by the telescope planes, multiple bubble events and large-angle scatters outside the acceptance of the downstream plaquettes. The last two categories are predominantly the result of inelastic interactions. Figure 1 shows an example scattering event.

At the end of the run the CF_3I was removed and a target empty data set was taken. In addition, data were taken in a test run in December 2011 with no target, as well as solid targets of quartz, graphite, Teflon or $(\text{C}_2\text{F}_4)_n$, and crystalline iodine.

II. ANALYSIS

The primary analysis output is the bubble-nucleation fraction as a function of E_{Ie} , given by the ratio

$$r_{\text{bub}} = \frac{N_{\text{bub}}}{(N_{\text{tot}} - N_{\text{multi}})f_{\text{CF3I}}}, \quad (4)$$

where N_{bub} is the observed number of pion tracks creating single bubbles, N_{tot} is the total number of pion tracks, N_{multi} is the number of tracks creating multiple bubbles,

and f_{CF3I} is the fraction of scatters that occur in the active CF_3I volume, determined by a comparison of the number of scatters in the target-full data set to the number in the target-empty data set normalized to the number of pion tracks (N_{emp}):

$$f_{\text{CF3I}} = (N_{\text{tot}} - N_{\text{emp}})/N_{\text{tot}}. \quad (5)$$

An angular smearing correction is made to N_{emp} to include the MCS from the absent CF_3I by convolution with the standard Gaussian approximation for MCS [15].

Each pion track is fitted for an upstream and downstream component, with an associated scattering angle and 3D point of closest approach of the two components. The upstream and downstream track segments are required to have exactly one hit cluster in at least three of the four pixel planes, good fits to straight lines ($\chi^2/\nu < 4$), and to meet in space to within 0.5 mm. To exclude pions that passed through little or no CF_3I , the upstream track is required to pass within 4 mm of the center of the 10-mm-diameter bubble chamber in the x direction. The y location of the track is limited by the vertical extent of the pixel planes. Because the uncertainty on the location of the point of closest approach in the beam direction (z) depends strongly on the scattering angle, we require the z location to be within $3\sigma_z$ of the bubble chamber, where σ_z is the uncertainty on z for each individual event. Events with more than one track are rejected. As these track cuts are applied without regard to the presence of nucleations in the bubble chamber, their efficiency applies equally to N_{bub} , N_{tot} , N_{multi} , and N_{emp} , and therefore cancels in the final ratio r_{bub} .

The next step is to associate a bubble with a unique track using both time and space correlations. The timing requirement for correlating a track with a bubble is chosen to be $20 < \Delta t < 120 \mu\text{s}$. The bubble locations are reconstructed using standard COUPP techniques [1], and the difference between reconstructed bubble position and point of closest approach of the track components is required to be less than 2.1 mm in the x direction and less than 0.9 mm in the y direction. The combined event acceptance of these timing and spatial cuts is 0.958 ± 0.011 . After these data selection and quality cuts, 350 good single-bubble events remain. The final bubble-nucleation fraction r_{bub} is shown as the points in Fig. 2.

To disentangle the iodine component from carbon, fluorine and inelastic scattering, we perform a full simulation using GEANT4.9.5 [16]. The simulation was validated by comparing the simulated scattering angular distributions to data for no target, target-empty, target-full, and solid targets. In all cases, in the MCS-dominated small scattering angle region there is good (few percent) agreement with no adjustable parameters, suggesting that the telescope geometry is accurately modeled in the simulation. In the larger scattering angle region dominated by strong elastic scattering, the simulation systematically overestimates the

E. BEHNKE *et al.*

observed scattering rate by $\sim 40\%$. As this ratio is measured to be the same for Teflon (1.45 ± 0.10) and iodine (1.41 ± 0.12), we assume that the relative contributions of iodine, fluorine and carbon are being accurately described by the MC.

A significant systematic uncertainty is introduced by our developing understanding of the carbon and fluorine recoil nucleation efficiency in this low energy regime. Ongoing studies with *ad hoc* neutron sources [17] will reduce this uncertainty in the future, but here we apply the exponential carbon and fluorine efficiency model from [11]:

$$\epsilon(E) = 1 - \exp(-\alpha[(E - E_T)/E_T]), \quad (6)$$

where E is the nuclear recoil energy and E_T is the threshold energy. In [11] we found $\alpha = 0.15$ to be consistent with AmBe neutron calibration data.

We test the hypothesis that the iodine recoil nucleation efficiency follows the nominal Seitz model of a step function ($\alpha_I \gg 1$) with 100% efficiency above the Seitz threshold by fitting a step function to the data in the region $9 < E_{Ie} < 42$ keV, allowing E_T to float. The fit returns $E_T = (16.8^{+0.8}_{-1.1})$ keV, where the error bars are statistical. This value is 2.1σ higher than the Seitz model threshold (13.6 ± 0.6) keV $\pm 6\%$ (sys), where the systematic error includes absolute energy scale uncertainties of 3% in the beam momentum and 1% in the scattering angle stemming from uncertainty in the z positions of the plaquettes. The fit is shown as the dashed blue line in Fig. 2.

Figure 3 shows the inferred iodine nucleation efficiency as a function of E_{Ie} with the iodine component isolated by subtracting the simulated contributions from carbon, fluorine and inelastic scatters. The dashed blue curve is the best

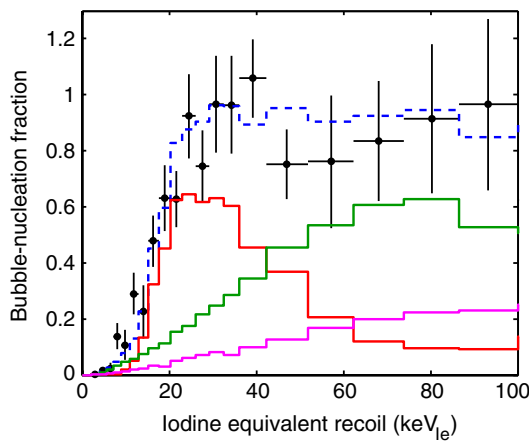


FIG. 2 (color online). The fraction of pion scattering events that produced bubbles as a function of iodine equivalent recoil energy. The solid curves show the simulated contribution from individual recoil species (from high to low at 20 keV, red for iodine, green for fluorine, and pink for carbon and inelastics), with the blue dashed curve showing the sum. The iodine curve shown takes a step function efficiency model for iodine recoils using the best fit threshold of $(16.8^{+0.8}_{-1.1})$ keV.

PHYSICAL REVIEW D **88**, 021101(R) (2013)

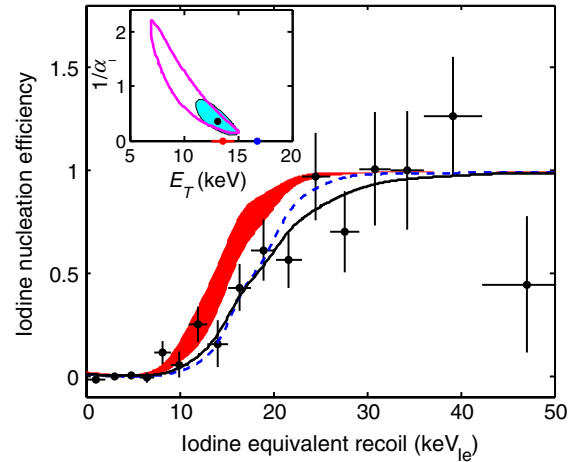


FIG. 3 (color online). The data points represent the measured iodine nucleation efficiency as a function of iodine equivalent recoil energy, where the contributions from carbon, fluorine and inelastic scatters have been subtracted. The gradual turn-on is predominantly due to the angular resolution of the experiment, as illustrated by both the red region, representing the step function model with the threshold varied within the uncertainty on the Seitz theory prediction, $E_T = (13.6 \pm 0.6)$ keV $\pm 6\%$, and the dashed blue curve, representing the best fit step function with $E_T = 16.8$ keV. The black curve shows the best fit exponential model with the threshold constrained by the theory as described in the text. The inset shows 2σ contours for a fit to the exponential model with the threshold allowed to float (pink) or constrained by the theory (solid cyan). The colored dots represent the corresponding curves in the main plot.

fit step function with $E_T = (16.8^{+0.8}_{-1.1})$ keV. For comparison, the red region represents a step function at the predicted Seitz threshold, where the range represents the 1σ band including the thermodynamic uncertainty and the scale uncertainties in the absolute energy scale of the experiment.

Given the energy resolution smearing induced by MCS in this experiment, the preference for a value of E_T higher than the prediction cannot be easily distinguished from an exponential model like Eq. (6) for iodine nucleation efficiency with a lower threshold energy and a finite value of α_I . Previous studies have shown that the Seitz model accurately predicts the threshold at which bubble

TABLE I. Summary of fits to α_I , including 90% lower limits on α_I . The three different C/F efficiency models described in the text are tested, and in all cases the predicted Seitz threshold is treated as a nuisance parameter. By maximizing the subtracted C/F contribution, the step function with $\eta_{C,F} = 100\%$ represents the worst case for iodine efficiency.

C/F efficiency model	α_I	90% L.L.	χ^2/ν
Flat model ($\eta_{C,F} = 49\%$)	$4.5^{+4.4}_{-1.3}$	>2.8	14.4/10
Exp. ($\alpha_{C,F} = 0.15$)	$2.8^{+1.6}_{-0.8}$	>1.8	8.2/10
Step function ($\eta_{C,F} = 100\%$)	$2.2^{+0.9}_{-0.7}$	>1.4	6.6/10

nucleation begins for heavy radon daughter nuclei in CF_3I [1,11]. We therefore perform a second fit applying the exponential model to iodine recoils, taking the Seitz threshold calculation as an external input to the analysis to explore the allowed range of α_I . The best fit is shown as the black curve in Fig. 3. The inset shows 2σ contours for fits to the exponential model with the threshold constrained by our prediction (shaded region) and free (unshaded region), along with the best fit points.

To assess the systematic errors associated with carbon and fluorine recoils, we refit the data with two alternative models for carbon and fluorine efficiency: the flat model from [11] with energy-independent nucleation efficiency, $\eta_{C,F} = 49\%$, above threshold, and a step function with $\eta_{C,F} = 100\%$. The latter case represents the worst possible scenario for the response of the bubble chamber to iodine recoils, as the response to carbon and fluorine is maximized. We use the exponential model for iodine recoils, allowing α_I to float and treating E_T as a nuisance parameter constrained by the prediction. The results of these fits are summarized in Table I. Extended fits over the energy interval $9 < E_{Ie} < 100$ keV have a negligible effect on the iodine fit parameters but disfavor the flat C/F efficiency model with $\eta_{C,F} = 49\%$.

In conclusion, we have directly measured the efficiency for iodine bubble nucleation in a CF_3I bubble chamber operated with a nominal threshold of (13.6 ± 0.6) keV. For some models of carbon and fluorine efficiency, the response to iodine recoils is consistent with a step function at the Seitz threshold, but in all cases there is a preference

for either a softer turn-on or a slightly higher threshold. Even in the worst case scenario for iodine, however, the response of the chamber to iodine recoils is much closer to the nominal Seitz model than it is for carbon and fluorine recoils. This was expected from the considerably larger stopping power of iodine, which facilitates the concentration of energy that leads to critical bubble formation. Systematic uncertainties from both the absolute beam momentum calibration and the carbon and fluorine response limit the present measurement. This measurement provides confirmation of the sensitivity of COUPP bubble chambers to spin-independent WIMP interactions with iodine nuclei, a confirmation that was not attainable using standard neutron source calibrations. The technique of employing hadron elastic scattering as a tool to measure bubble-nucleation thresholds is now established, enabling the measurement of bubble-nucleation energies on an event by event basis. We have begun studies of the feasibility to repeat this technique with different fluids.

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