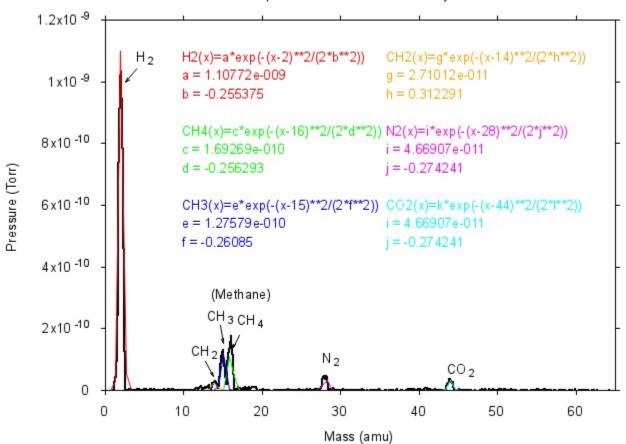
Josh Yoskowitz Ionization Cross Sections for "After 2 Days" RGA Spectrum

1 Purpose

To calculate electron-impact ionization cross sections for gas species found in the "After 2 Days" residual gas analyzer (RGA) spectrum taken on 5/21/18. The spectrum was analyzed using gnuplot and is shown below in Figure 1. Each substantial peak was identified and fit with a Gaussian function in order to determine the partial pressures of the various species of residual gas in the gun chamber. NOTE: The peak values must be divided by the correction factors listed here: https://www.mksinst.com/docs/ur/GaugeGasCorrection.aspx



RGA Spectrum for the "After Two Days" Data

Figure 1: Analysis of the RGA spectrum for the "After 2 Days" data (before correction factor)

2 Calculation of the Ionization Cross Section

The equation for the calculation of the ionization cross section σ_i of the i^{th} gas species can be found in Reiser [1] and was originally developed by Slinker et. al. [3]:

$$\sigma_i = \frac{8a_0^2\pi I_R A_1}{m_e c^2 \beta^2} f\left(\beta\right) \left(\ln \frac{2A_2 m_e c^2 \beta^2 \gamma^2}{I_R} - \beta^2\right) \tag{1}$$

Numerically, this can be rewritten as:

$$\sigma_{i[m^{2}]} = \frac{1.872 \times 10^{-24} A_{1}}{\beta^{2}} f\left(\beta\right) \left[\ln\left(7.515 \times 10^{4} A_{2} \beta^{2} \gamma^{2}\right) - \beta^{2}\right]$$
(2)

In these two equations, $a_0 = 5.29 \times 10^{-11}$ m is the Bohr radius, $I_R = 13.6$ eV is the Rydberg energy, $m_e c^2$ is the rest mass energy of the electron, and β and γ are relativistic factors, A_1 and A_2 are empirical constants that depend on the type of gas species, and $f(\beta)$ is a function used when fitting data at low energies, i.e. $T_e \approx I_i$ where T_e is the kinetic energy of the electron and I_i is the ionization energy for the i^{th} gas species. Expressions for A_1 , A_2 , and $f(\beta)$ are given below:

$$f(\beta) = \frac{I_i}{T_e} \left(\frac{T_e}{I_i} - 1 \right) = \frac{2I_i}{m_e c^2 \beta^2} \left(\frac{m_e c^2 \beta^2}{2I_i} - 1 \right)$$
(3)

$$A_1 = M^2 \tag{4}$$

$$A_2 = \frac{e^{\frac{1}{M^2}}}{7.515 \times 10^4} \tag{5}$$

where C and M^2 are parameters given by Rieke and Prepejchal [2]. For H₂, CH₄, CH₃, N₂, and CO₂ the values of C, $M^2 = A_1$, A_2 , and the ionization energy I_i from NIST (https://webbook.nist.gov/) are given in the table below:

Gas Species	$A_1 = M^2$	C	A_2	$I_i(eV)$
H_2	0.695	8.115	1.5668	15.4
CH_4	4.23	41.85	0.2635	12.6
N_2	3.74	34.84	0.1478	15.6
CO_2	5.75	55.92	0.2227	13.8

Table 1: Values for C, $M^2 = A_1$, and A_2 given by Rieke and Prepejchal and I_i given by NIST for gas species found in the RGA spectrum.

Since at high energies, $\beta_e \gg \beta_{ion}$, we will assume that in the above equations, $\beta \approx \beta_e$. As an example calculation, for a 200keV electron beam, $T_e = 200$ keV, $m_e c^2 = 511$ keV, the cross section for H₂ gas is:

$$\begin{split} T_e &= (\gamma_e - 1) \, m_e c^2 = 200 \text{keV} \\ m_e c^2 &= 511 \text{keV} \\ \gamma_e &= 1 + \frac{T_e}{m_e c^2} = 1.39 \\ \beta_e &= \sqrt{1 - \frac{1}{\gamma_e^2}} = 0.695 \left(= 2.08 \times 10^8 \frac{\text{m}}{\text{s}} \right) \\ f\left(\beta_e\right) &= \frac{I_i}{T_e} \left(\frac{T_e}{I_i} - 1\right) \approx 1 \\ \sigma_i &= \frac{1.872 \times 10^{-24} A_1}{\beta_e^2} f\left(\beta_e\right) \left[\ln\left(7.515 \times 10^4 A_2 \beta_e^2 \gamma_e^2\right) - \beta_e^2\right] \\ &\approx 2.994 \times 10^{-23} \text{m}^2 \end{split}$$

3 Ionization Rate

The change in density of the electron and gas molecules over time is given by Reiser[1]

$$\frac{dn}{dt} = n_b n_g \sigma_i v = n_b n_g \sigma_i \beta_e c \tag{6}$$

At standard temperature ($T_0 = 273.15$ K) and pressure ($p_0 = 760$ torr = 1atm) the density of an ideal gas in a given volume is given by Loschmidt's number:

$$n_0 = \frac{p_0}{k_B T_0} \approx 2.687 \times 10^{25} \mathrm{m}^{-3} \tag{7}$$

Thus, for a given gas, its density is

$$n_g \left[m^{-3} \right] = (3.54 \times 10^{22}) p (torr)$$

For H₂, p (torr) can be read off from Figure 1. In this case, $p_{H_2} = 1.11 \times 10^{-9}$ torr with a correction factor of 0.46 for H₂, so $n_{H_2} = 8.54 \times 10^{13} \text{m}^{-3}$. Knowing the electron density in the beam, n_b , one can calculate the ionization rate $\frac{dn}{dt}$. For a numerical example, assume we have a 200keV, 1mA uniform, cylindrical electron beam with an average transverse size of 1mm. In this case, $T_e = 200 \text{keV}$, I = 1 mA, $\sigma_{H_2} = 2.994 \times 10^{-23} \text{m}^2$, and $n_b \approx 1.98 \times 10^{21} \text{m}^{-3}$. Then for H₂,

$$\frac{dn_{H_2}}{dt} = n_{H_2} n_b \sigma_{H_2} \beta_e c$$

= $\left(8.54 \times 10^{13} \text{m}^{-3}\right) \left(1.98 \times 10^{21} \text{m}^{-3}\right) \left(2.994 \times 10^{-23} \text{m}^2\right) \left(0.695\right) \left(2.998 \times 10^8 \frac{\text{m}}{\text{s}}\right)$
 $\approx 1.05 \times 10^{21} \text{m}^{-3} \text{s}^{-1}$

4 Ionization Cross Section vs. T_e

Starting from equation (2),

$$\sigma_i \left[\mathbf{m}^2 \right] = \frac{1.872 \times 10^{-24} A_1}{\beta^2} \frac{I_i}{T_e} \left(\frac{T_e}{I_i} - 1 \right) \left[\ln \left(7.515 \times 10^4 A_2 \beta^2 \gamma^2 \right) - \beta^2 \right]$$

we can rewrite β in terms of the electron beam kinetic energy T_e , which is proportional to the beam voltage:

$$\begin{split} T_{e} &= (\gamma - 1) \, m_{e} c^{2} \\ \gamma &= 1 + \frac{T_{e}}{m_{e} c^{2}} \\ \frac{1}{\sqrt{1 - \beta^{2}}} &= 1 + \frac{T_{e}}{m_{e} c^{2}} \\ 1 - \beta^{2} &= \left(\frac{1}{1 + \frac{T_{e}}{m_{e} c^{2}}}\right)^{2} = \left(\frac{m_{e} c^{2}}{m_{e} c^{2} + T_{e}}\right)^{2} \\ \beta^{2} &= 1 - \left(\frac{m_{e} c^{2}}{m_{e} c^{2} + T_{e}}\right)^{2} \end{split}$$

Thus,

$$\sigma_{i} = \frac{1.872 \times 10^{-24} A_{1}}{1 - \left(\frac{m_{e}c^{2}}{m_{e}c^{2} + T_{e}}\right)^{2}} \frac{I_{i}}{T_{e}} \left(\frac{T_{e}}{I_{i}} - 1\right) \left[\ln \left(7.515 \times 10^{4} A_{2} \left(1 - \left(\frac{m_{e}c^{2}}{m_{e}c^{2} + T_{e}}\right)^{2}\right) \left(1 + \frac{T_{e}}{m_{e}c^{2}}\right) \right] - \left(1 - \left(\frac{m_{e}c^{2}}{m_{e}c^{2} + T_{e}}\right)^{2}\right) \left(1 + \frac{T_{e}}{m_{e}c^{2}}\right) = \frac{1}{2} \left(1 - \left(\frac{m_{e}c^{2}}{m_{e}c^{2} + T_{e}}\right)^{2}\right) \left(1 - \left(\frac{m_{e}c^{2}}{$$

Using values in Table 1, a plot of σ_i vs. T_e for each of the gas species was made using Mathematica:

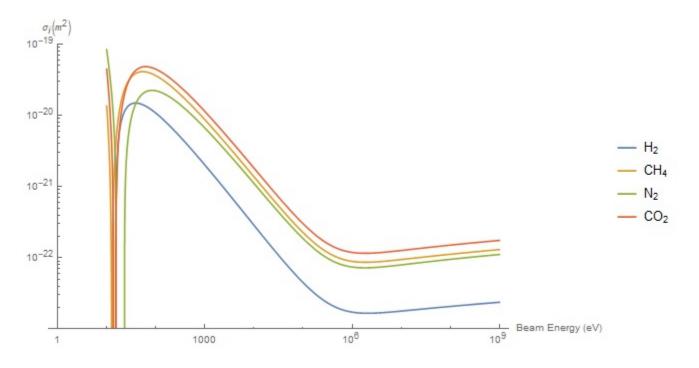


Figure 2: Plot of the ionization cross section σ_i vs. electron kinetic energy T_e

References

- [1] Martin Reiser. Theory and Design of Charged Particle Beams. Wiley VCH Verlag GmbH, 2008.
- [2] Foster F. Rieke and William Prepejchal. Ionization cross sections of gaseous atoms and molecules for high-energy electrons and positrons. *Physical Review A*, 6(4):1507–1519, oct 1972.
- [3] S. P. Slinker, R. D. Taylor, and A. W. Ali. Electron energy deposition in atomic oxygen. Journal of Applied Physics, 63(1):1-10, jan 1988.