

# Measurement of $^{16}\text{O}(\gamma,\alpha)^{12}\text{C}$ with a Bubble Chamber and a Bremsstrahlung Beam at Jefferson Lab Injector

Riad Suleiman

January 9, 2014



B. DiGiovine  
D. Henderson  
R. J. Holt  
K. E. Rehm



A. Robinson  
C. Ugalde



J. Benesch  
P. Degtiarenko  
A. Freyberger  
J. Grames  
G. Kharashvili  
D. Meekins  
M. Poelker  
Y. Roblin  
R. Suleiman  
C. Tennant  
V. Vylet



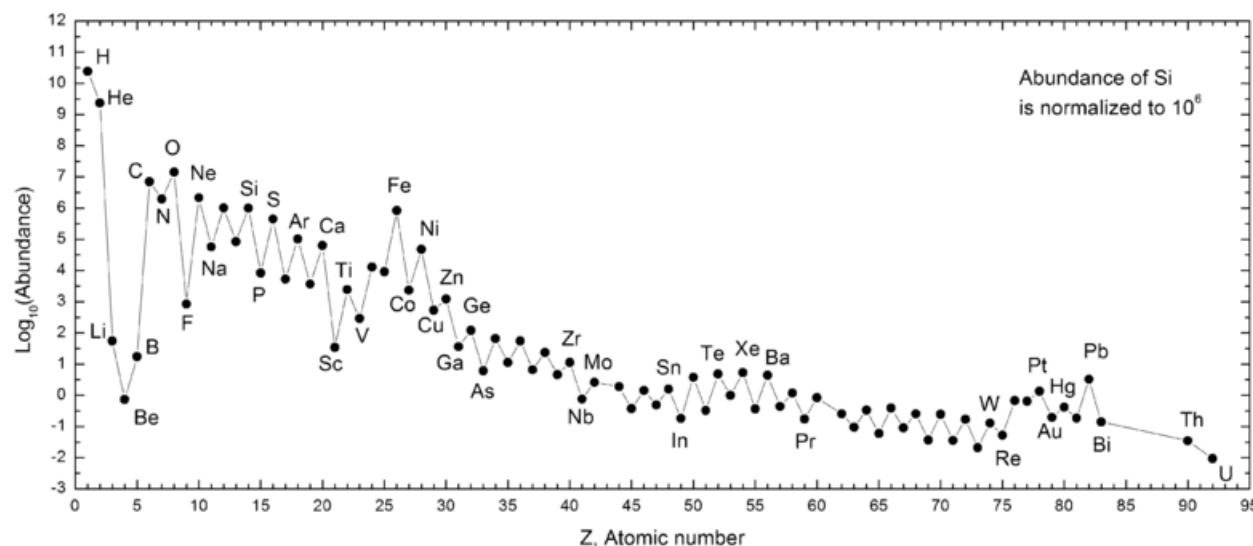
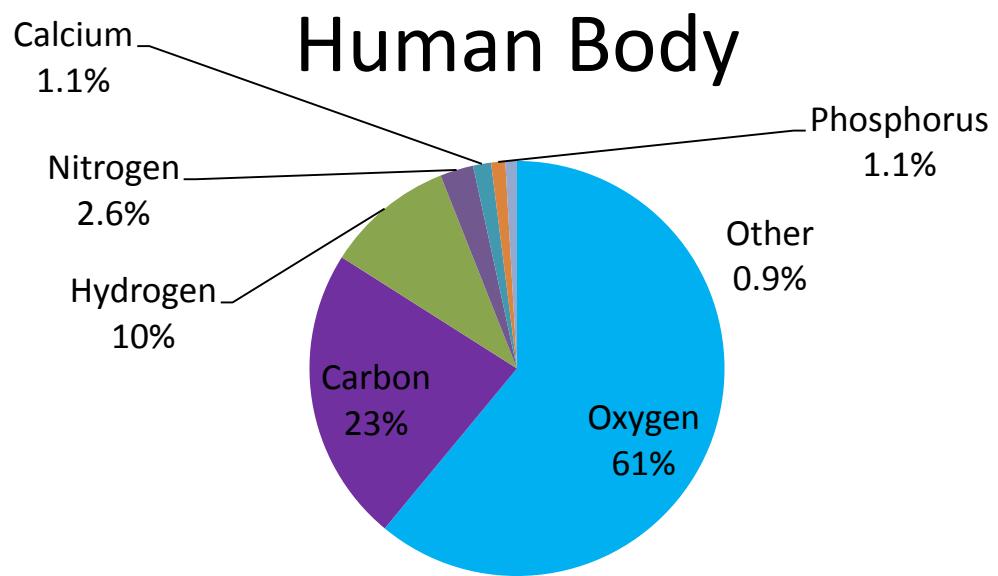
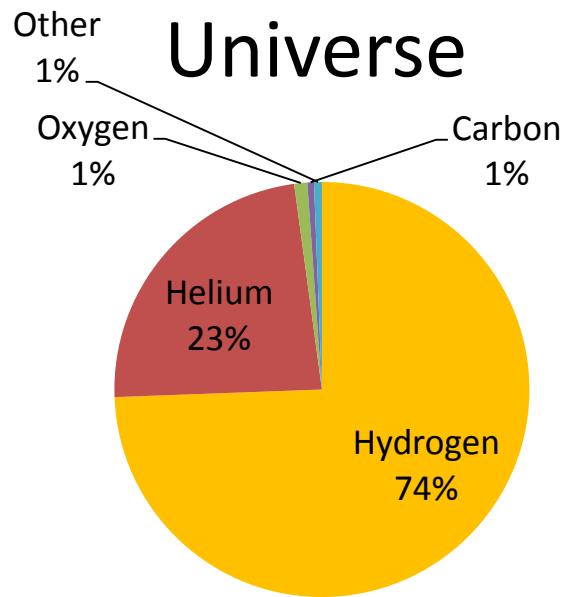
A. Sonnenschein

[https://wiki.jlab.org/ciswiki/index.php/Bubble\\_Chamber](https://wiki.jlab.org/ciswiki/index.php/Bubble_Chamber)

# OUTLINE

- Nucleosynthesis and the  $^{12}\text{C}(\alpha,\gamma)^{16}\text{O}$  Reaction
- Time Reversal Reaction:  $^{16}\text{O}(\gamma,\alpha)^{12}\text{C}$
- The Bubble Chamber
- Work at HIGS
- Experimental Setup at Jefferson Lab Injector
- Bremsstrahlung Beam and Penfold-Leiss Unfolding
- Statistical and Systematic Errors
- Backgrounds and Ion Energy Distributions
- Safety
- Summary and Outlook

# RELATIVE ABUNDANCE OF ELEMENTS BY WEIGHT

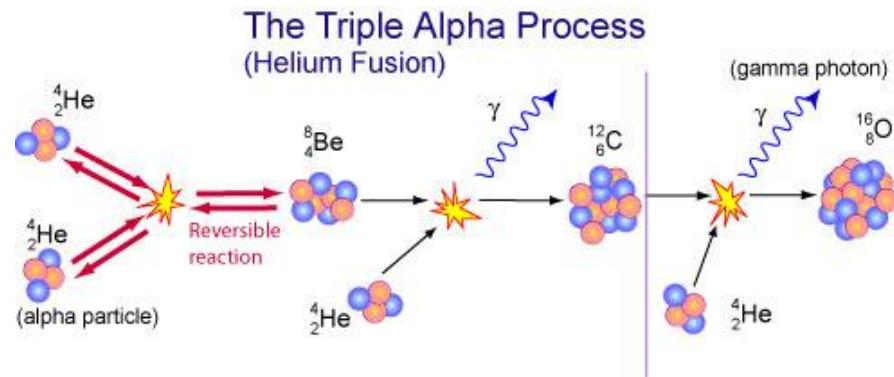


# STELLAR HELIUM BURNING

- Helium Reactions:



(slow, otherwise no  ${}^{12}\text{C}$  remains)



- $\alpha + {}^{12}\text{C}$  burning at very small cross section  $\sigma \approx 10^{-17}$  barn

➡ Currently, reaction rate error is large ( $\pm 35\%$ )

Goal  $< \pm 10\%$

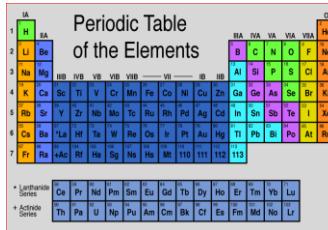
- Thermonuclear reaction rate involving two nuclei is:

Only narrow energy range is important (Gamow Peak)

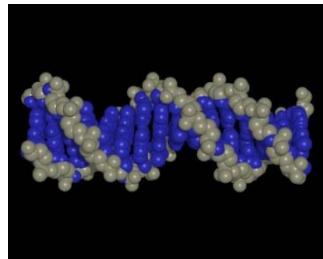
$$R = \sqrt{\frac{8}{\pi m (k_B T)^3}} \int_0^\infty E \sigma_{tot}(E) e^{-\frac{E}{k_B T}} dE$$

# THE $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$ Reaction

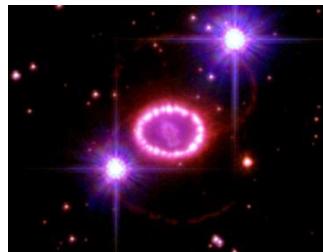
- The *holy grail* of nuclear astrophysics



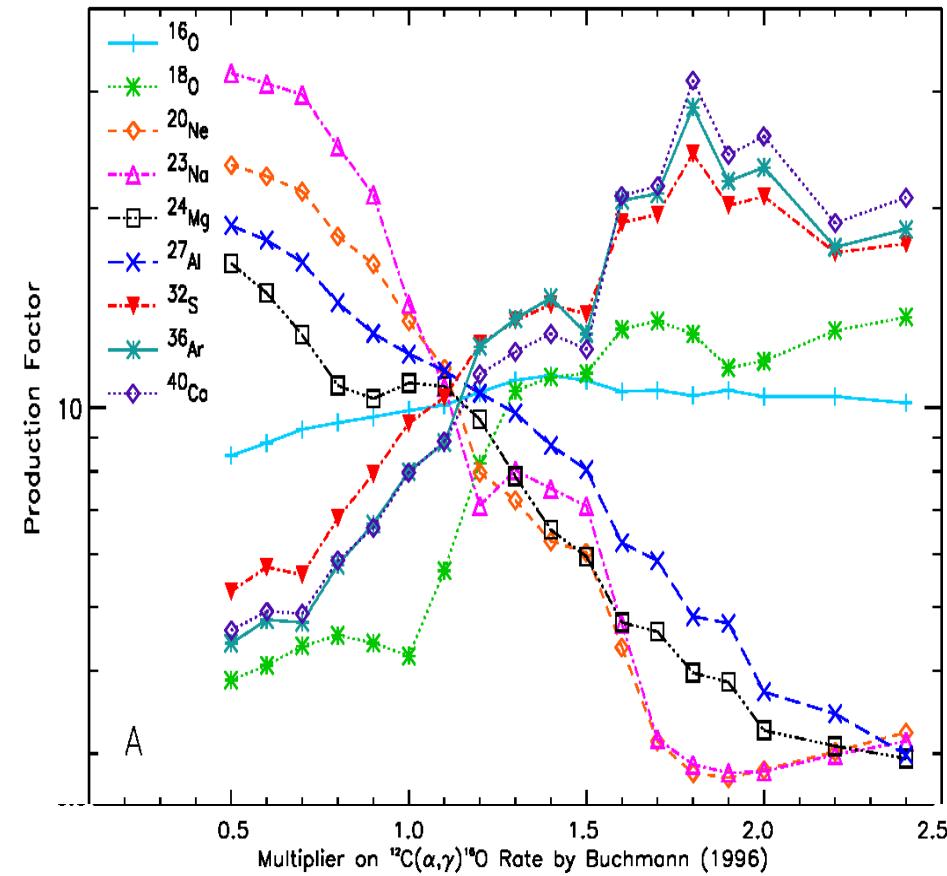
Affects the synthesis of most of the elements of the periodic table



Sets the  $\text{N}^{(12)\text{C}}/\text{N}^{(16)\text{O}}$  ( $\approx 0.4$ ) ratio in the universe



Determines the minimum mass a star requires to become a supernova



# THE GAMOW PEAK

- Narrow energy range where thermonuclear reactions is most likely to occur in stellar plasma is a product of two distributions:
  - Maxwell-Boltzmann energy distribution with  $e^{-E/k_B T}$
  - Penetration through Coulomb barrier with  $e^{-b/E^{1/2}}$

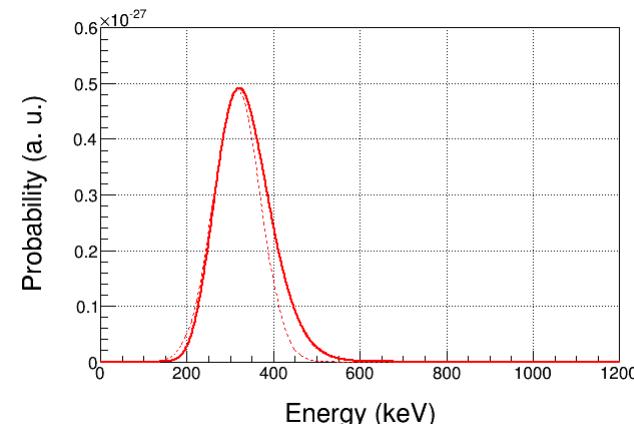
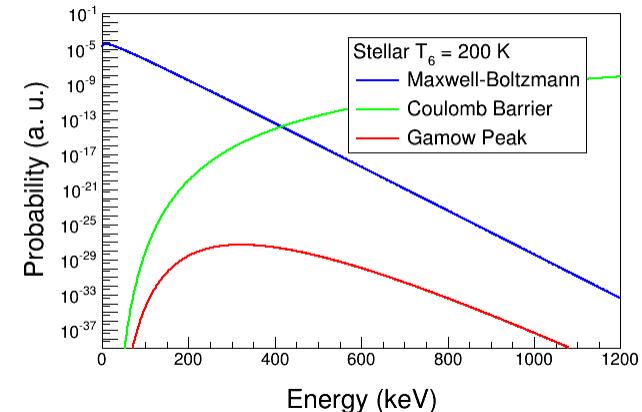
$$E_0 = 1.220 \left( Z_1^2 Z_2^2 A T_6^2 \right)^{1/3} \text{ keV}$$

$$W = 0.2368 \left( Z_1^2 Z_2^2 A T_6^5 \right)^{1/6} \text{ keV}$$

- For  $\alpha + {}^{12}\text{C}$  ( $Z_1=2$ ,  $Z_2=6$ ,  $A=3$ ),

and stellar  $T=200 \times 10^6 \text{ K}$ :

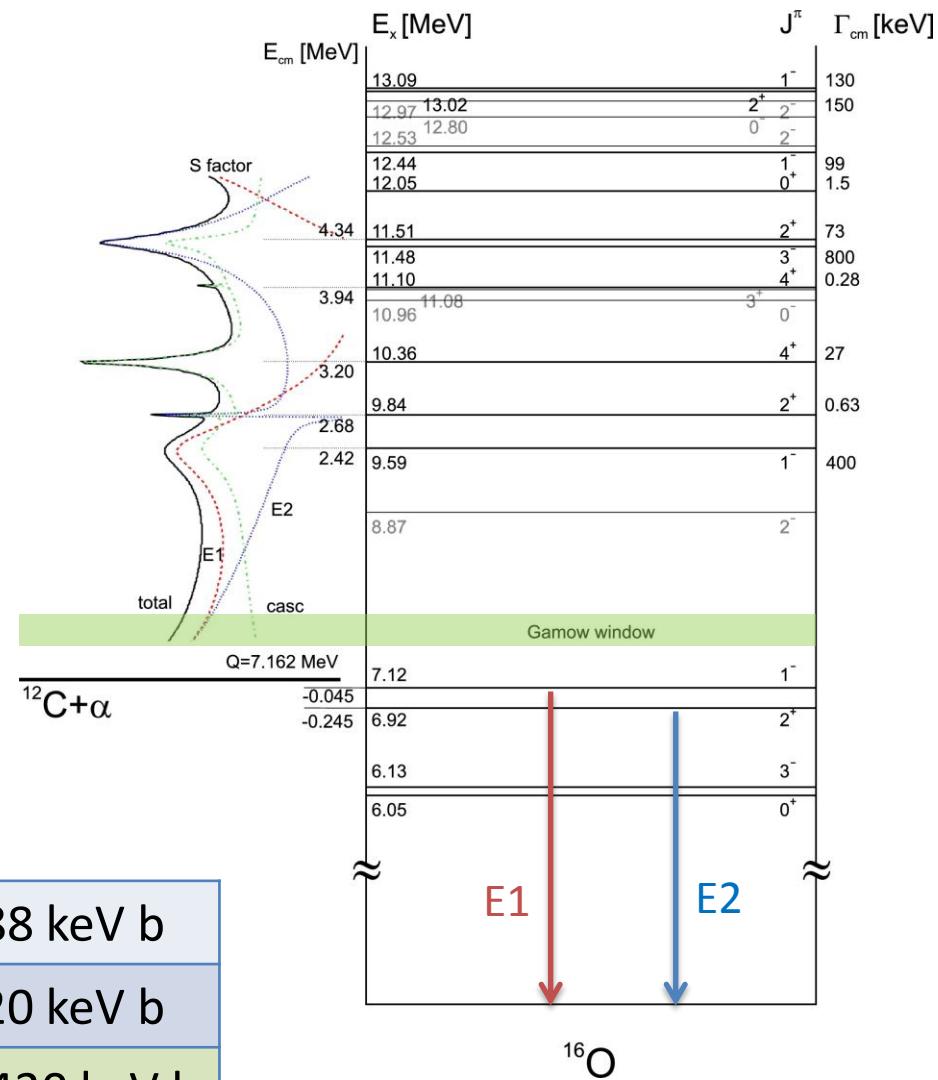
- **Gamow Peak,  $E_0 \approx 300 \text{ keV}$ ,  $W \approx 50 \text{ keV}$**
- Maximum of Maxwell–Boltzmann energy distribution,  $k_B T = 17 \text{ keV}$



# $\alpha + {}^{12}\text{C}$ RADIATIVE CAPTURE

- $\sigma(E_0)$  is dominated by  $p$ -wave (E1) and  $d$ -wave (E2) radiative capture to ( $J^\pi=0^+$ )  ${}^{16}\text{O}$  ground state
- Two bound states, at 6.92 MeV ( $J^\pi=2^+$ ) and 7.12 MeV ( $J^\pi=1^-$ ), with sub-threshold resonances at  $E_R=-0.245$  and -0.045 MeV, provide most of  $\sigma(E_0)$  through their finite widths
- Distinguish E1 and E2 by measuring  $\gamma$  angular distributions

|                                 |  |
|---------------------------------|--|
| Transition $\rightarrow 0$ (E1) | $S_{\text{E}1}(300) = 1\text{--}288 \text{ keV b}$   |
| Transition $\rightarrow 0$ (E2) | $S_{\text{E}2}(300) = 7\text{--}120 \text{ keV b}$   |
| Total                           | $S_{\text{tot}}(300) = 40\text{--}430 \text{ keV b}$ |



# Heroic efforts in search of $^{12}\text{C}(\alpha,\gamma)^{16}\text{O}$

## ➤ Previous Experiments:

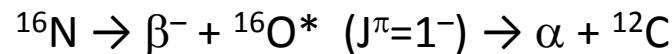
### A. Direct Measurements:

- I. Helium ions on carbon target:  $^{12}\text{C}(\alpha,\gamma)^{16}\text{O}$
- II. Carbon ions on helium gas:  $^4\text{He}(^{12}\text{C}, \gamma)^{16}\text{O}$  or  $^4\text{He}(^{12}\text{C}, ^{16}\text{O})\gamma$  (Schürmann)

| Experiment | Beam Current (mA) | Target (nuclei/cm <sup>2</sup> )    | Time (h) |
|------------|-------------------|-------------------------------------|----------|
| Redder     | 0.7               | $^{12}\text{C}$ , $3 \cdot 10^{18}$ | 900      |
| Ouellet    | 0.03              | $^{12}\text{C}$ , $5 \cdot 10^{18}$ | 1950     |
| Roters     | 0.02              | $^4\text{He}$ , $1 \cdot 10^{19}$   | 5000     |
| Kunz       | 0.5               | $^{12}\text{C}$ , $3 \cdot 10^{18}$ | 700      |
| EUROGAM    | 0.34              | $^{12}\text{C}$ , $1 \cdot 10^{19}$ | 2100     |
| GANDI      | 0.6 (?)           | $^{12}\text{C}$ , $2 \cdot 10^{18}$ | ?        |
| Schürmann  | 0.01              | $^4\text{He}$ , $4 \cdot 10^{17}$   | ?        |
| Plag       | 0.005             | $^{12}\text{C}$ , $6 \cdot 10^{18}$ | 278      |

### B. Indirect Measurements:

- I.  $\beta$ -delayed  $\alpha$  decay of  $^{16}\text{N}$  ( $J^\pi=2^-$ ,  $T_{1/2}=7.13$  s, BR=0.12%)



- II. Elastic  $\alpha - ^{12}\text{C}$  scattering

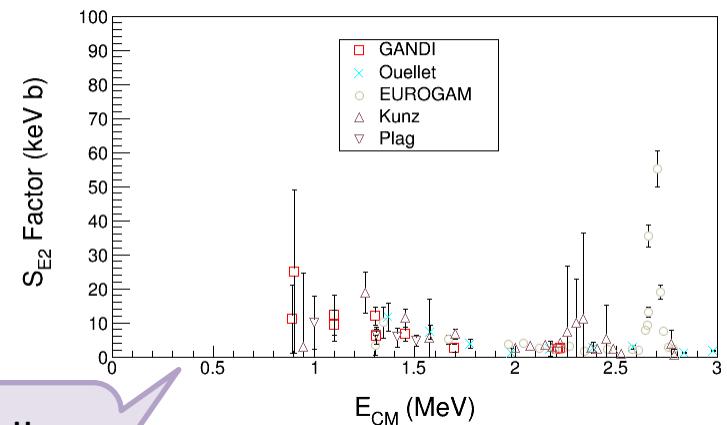
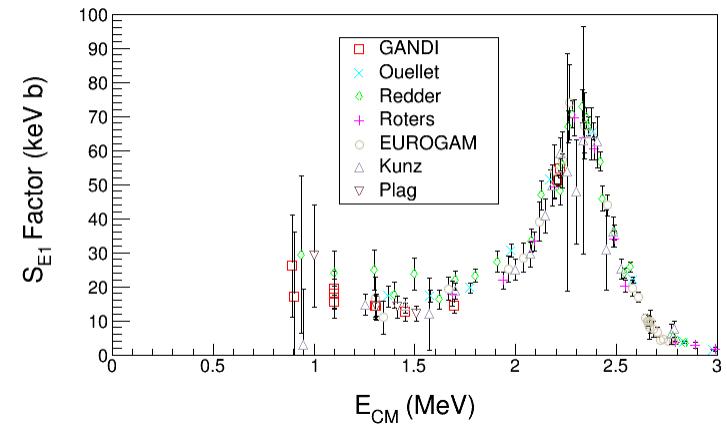
# ASTROPHYSICAL S-FACTOR $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$

- Define *S-Factor* to remove both  $1/E$  dependence of nuclear cross sections and Coulomb barrier transmission probability:

$$S \equiv E_{CM} \sigma(\alpha, \gamma) e^{2\pi\eta}$$

$$\eta = \frac{1}{137} Z_\alpha Z_{^{12}\text{C}} \sqrt{\frac{m_{^{12}\text{Ca}}}{2E_{CM}}}$$

| Author           | $S(300)$ (keV b)       |
|------------------|------------------------|
| Schürmann (2012) | $161 \pm 19^{+8}_{-2}$ |
| Hammer (2005)    | $162 \pm 39$           |
| Kunz (2001)      | $165 \pm 50$           |



R-matrix Extrapolation to stellar helium burning at  $E = 300$  keV

# RECIPROCITY RELATION: $(\gamma, \alpha)$ and $(\alpha, \gamma)$

➤ A( $\alpha, \gamma$ )B:

$$\sigma_{B\gamma}^{j \rightarrow i}(E_\gamma) = \frac{(2J_i + 1)(2J_\alpha + 1)}{2J_j + 1} \frac{m_{A\alpha} c^2 E_{A\alpha}}{E_\gamma^2} \sigma_{A\alpha}^{i \rightarrow j}(E_{A\alpha})$$

$$m_{A\alpha} c^2 = \frac{M(^{12}C) \cdot M(\alpha)}{M(^{12}C) + M(\alpha)} = 2796 \text{ MeV} \quad J_i = 0, J_j = 0, J_\alpha = 0$$

$$E_{A\alpha} = E_{CM}$$

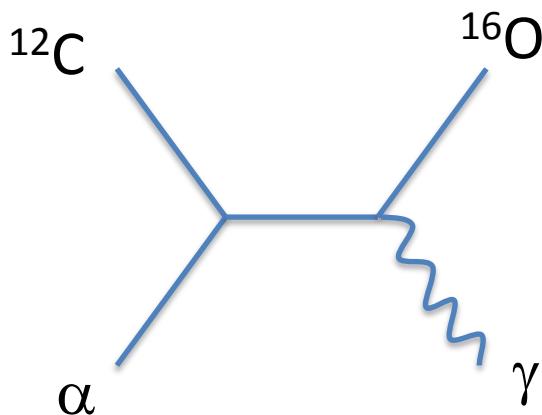
$$Q = m_A + m_\alpha - m_B = 7.162 \text{ MeV}$$

$$E_{CM} = \sqrt{m_B^2 + 2E_\gamma m_B} - m_B - Q \\ \cong E_\gamma - Q$$

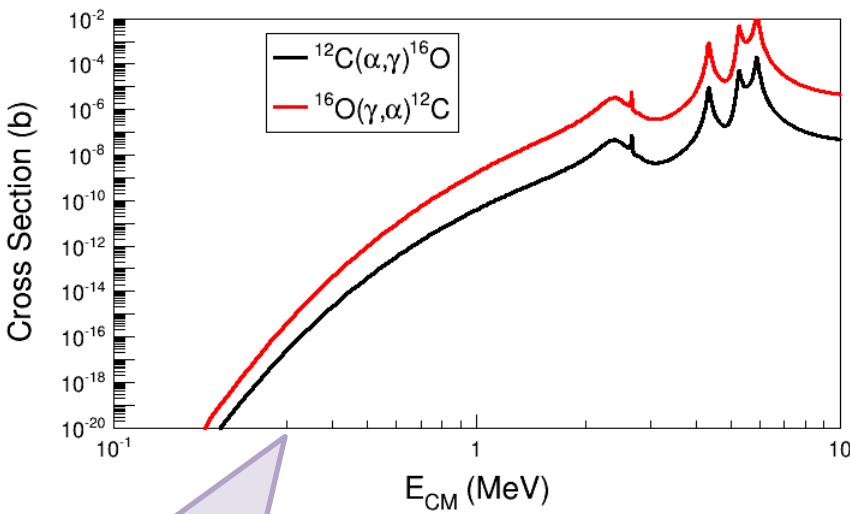
$$\sigma_{(\gamma, \alpha)}(E_\gamma) = \frac{m_{A\alpha} c^2 E_{CM}}{E_\gamma^2} \sigma_{(\alpha, \gamma)}(E_{CM})$$

➤  $\sigma(\gamma, \alpha)$  is over two orders of magnitude larger than  $\sigma(\alpha, \gamma)$

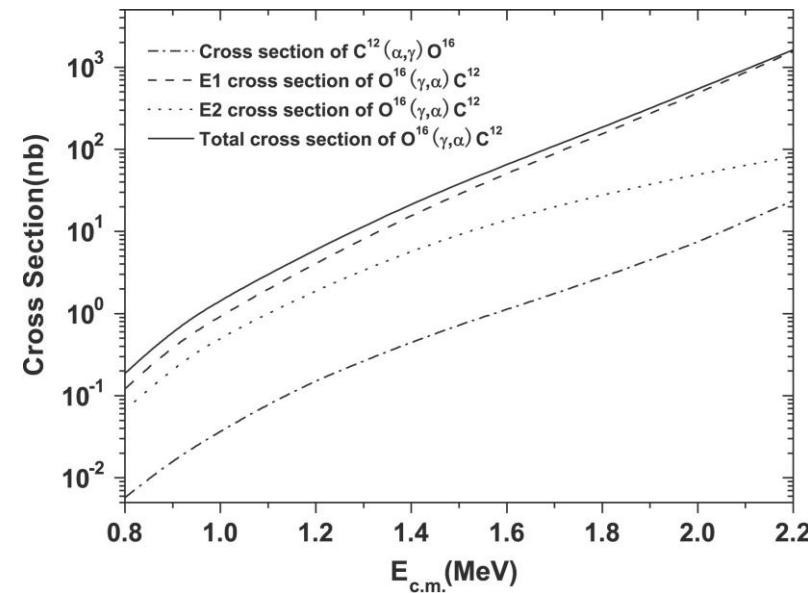
# TIME REVERSAL REACTION



- Bubble Chamber experiment measures total cross section, E1 + E2.
- We can separate E1 and E2 if we use linearly polarized  $\gamma$  but cannot measure  $\alpha$  and  $^{12}\text{C}$  angular distribution

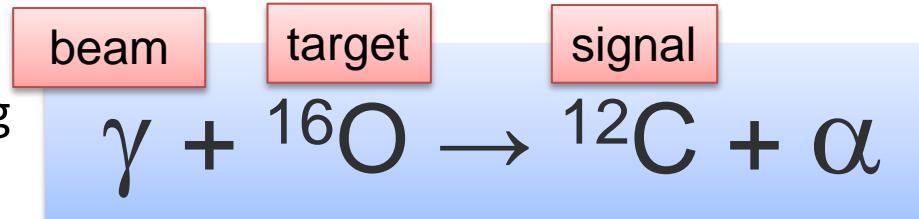


Stellar helium burning  
at  $E = 300$  keV



# NEW APPROACH: REVERSAL REACTION + BUBBLE CHAMBER

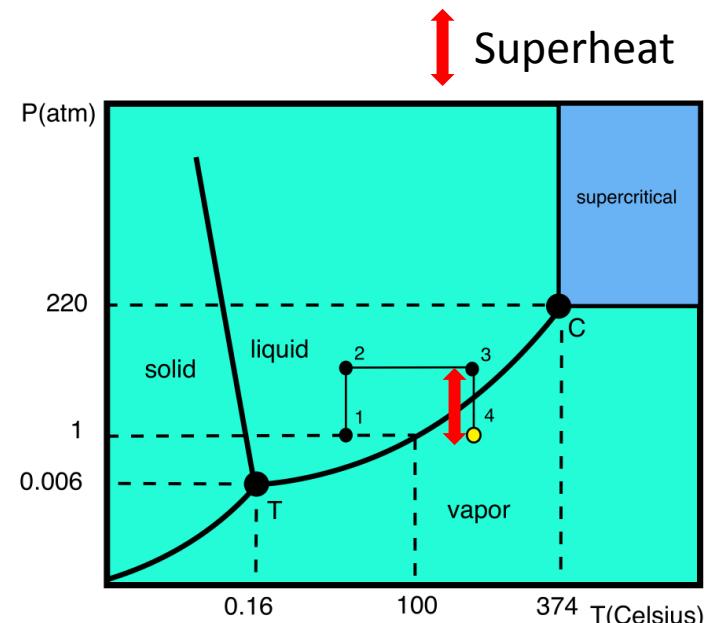
- Extra gain (factor of 100) by measuring time reversal reaction
- Target density up to  $10^4$  higher than conventional targets. Number of  $^{16}\text{O}$  nuclei =  $3.5 \cdot 10^{22} / \text{cm}^2$  (3.0 cm cell)
- Solid Angle and Detector Efficiency = 100%
- Superheated liquid will nucleate from  $\alpha$  and  $^{12}\text{C}$  recoils
- Electromagnetic debris (electrons and gammas, or positrons) do NOT trigger nucleation (detector is insensitive to  $\gamma$ -rays by at least 1 part in  $10^{11}$ ).



- Monochromatic  $\gamma$  beam at HIGS  $\approx 10^{7-8} \gamma/\text{s}$
- Bremsstrahlung at JLab  $\approx 10^9 \gamma/\text{s}$  (top 250 keV)

# THE BUBBLE CHAMBER

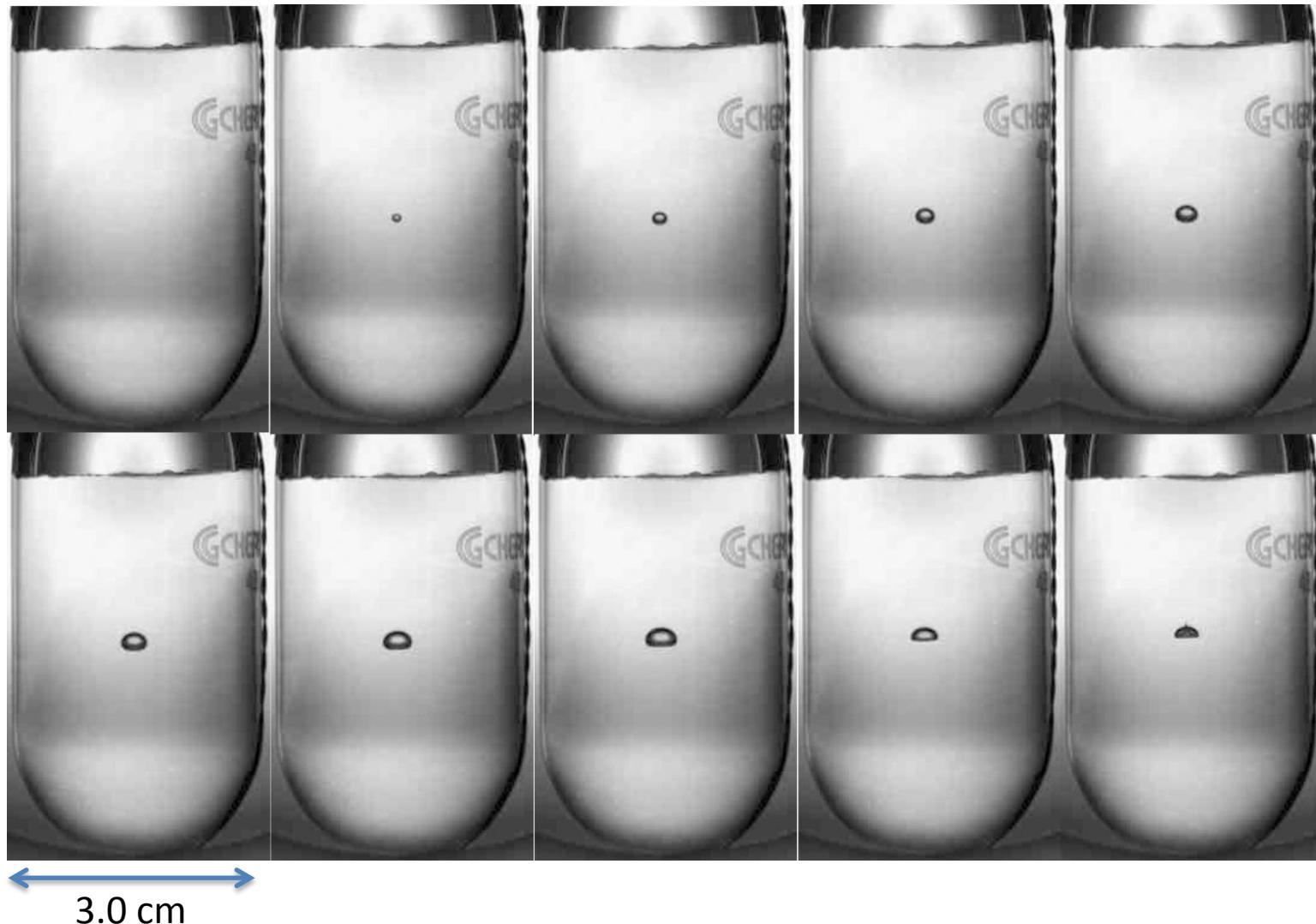
- Donald Glaser won Nobel Prize for inventing chamber to detect particles (1960)
- Now being used in Dark Matter Search Experiments: COUPP, PICASSO, SIMPLE
- Superheat Preparation:
  - Liquid is pressurized at ambient temperature (1 to 2)
  - Then pressure is kept constant while temperature is increased to above boiling point (2 to 3)
  - Finally pressure is slowly released while keeping temperature constant (3 to 4)
  - At this point (4), water is still liquid but now superheated
- Bubble Formation:
  - Particle energy loss will induce vaporization
  - Resultant vapor bubble is observable either **visibly or audibly**
  - Bubble growth is captured by a camera
  - Pressure is increased (4 to 3) to quench bubble. It takes about few seconds for liquid to return to a stable state
  - Superheat is restored by releasing pressure again (3 to 4), and cycle is repeated for each bubble event



# BUBBLE GROWTH AND QUENCHING

$^{19}\text{F}(\gamma,\alpha)^{15}\text{N}$  event in  $\text{C}_4\text{F}_{10}$

Fast Digital Camera:  $\Delta t = 10 \text{ ms}$



# BUBBLE CHAMBER PRINCIPLE

- I. Only bubbles with  $r > R_c$  grow to be macroscopic

$$R_c = 2s / (P_v - P_l)$$

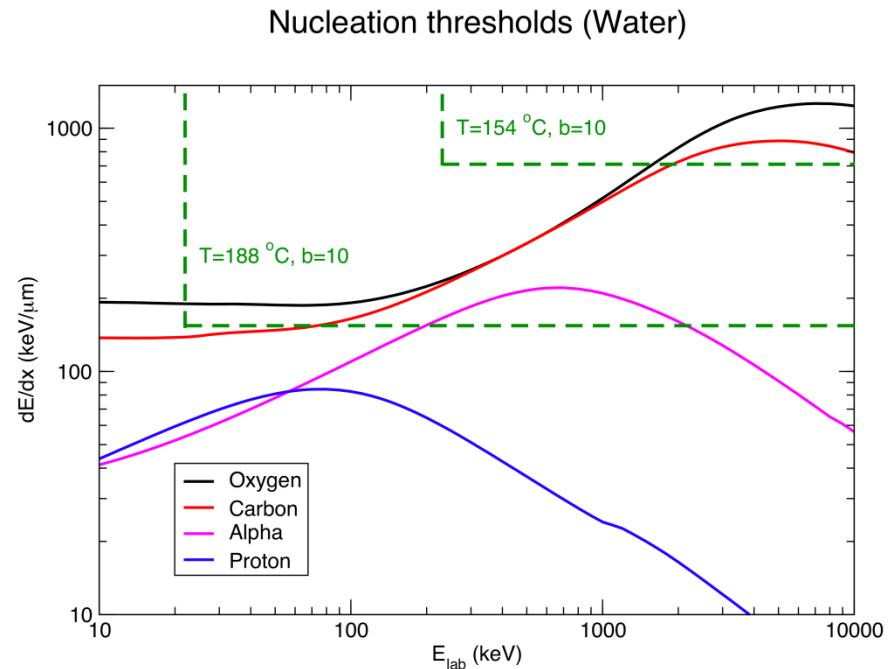
s: Surface tension

- II. Bubble requires minimum deposited energy ( $E_c$ ) within minimum distance  $L_c$  ( $=aR_c$ , 10s of nm to a few  $\mu\text{m}$ )

$$\frac{dE}{dx} > \left( \frac{dE}{dx} \right)_c = \frac{E_c}{aR_c}$$

a: free parameter (to determined experimentally)

- III. Particle must be over thresholds in both  $dE/dx$  and E



$$E \geq E_c = \frac{4}{3} \pi R_c^3 (\rho h + P_l) + 4\pi R_c^2 \left( s - T \frac{\partial s}{\partial T} \right)$$

# ACOUSTIC SIGNAL DISCRIMINATION

$\alpha$  should  
be louder

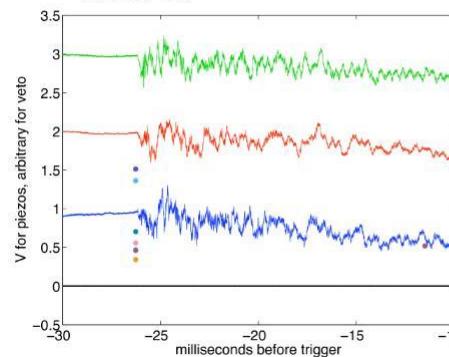
## I. Neutron Events:

- I.  $^{17}\text{O}(\gamma, n)^{16}\text{O}$
- II. Neutron–nucleus elastic scattering:  
 $^{16}\text{O}(n, n)$ ,  $^{14}\text{N}(n, n)$

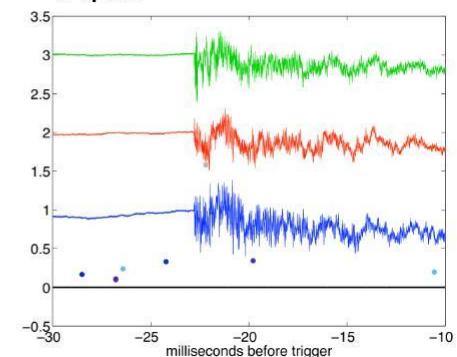
Ions  $^{16}\text{O}$  or  $^{14}\text{N}$  will generate a single bubble



Neutron



Alpha



## II. Alpha Events:

- I.  $^{16}\text{O}(\gamma, \alpha)^{12}\text{C}$
- II.  $^{17}\text{O}(\gamma, \alpha)^{13}\text{C}$
- III.  $^{18}\text{O}(\gamma, \alpha)^{14}\text{C}$

Ions  $^{12}\text{C}+\alpha$  or  $^{13}\text{C}+\alpha$  or  $^{14}\text{C}+\alpha$  will generate a combined multi-bubble

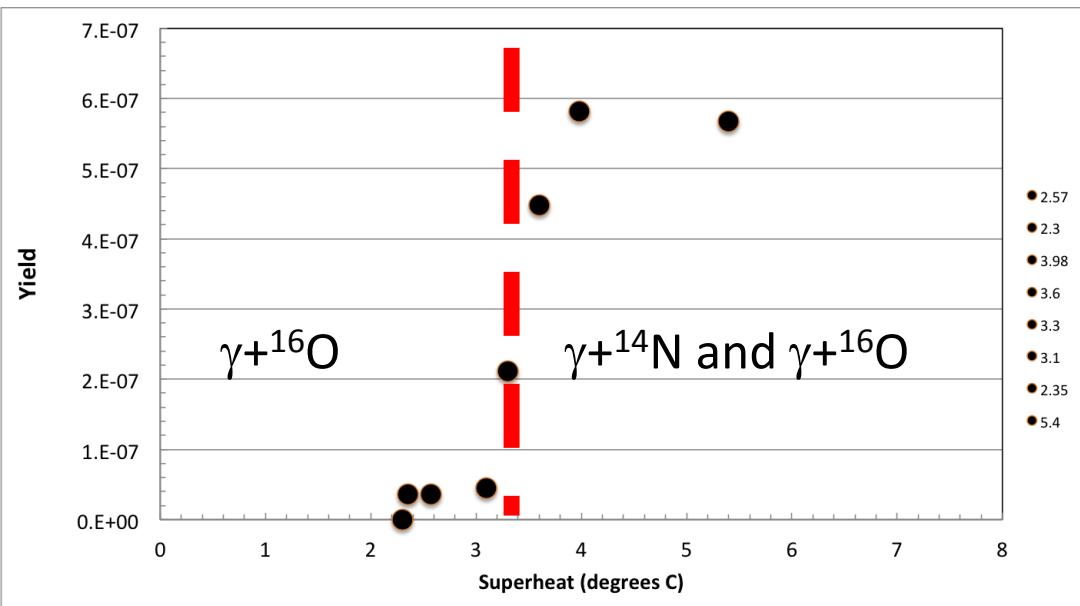
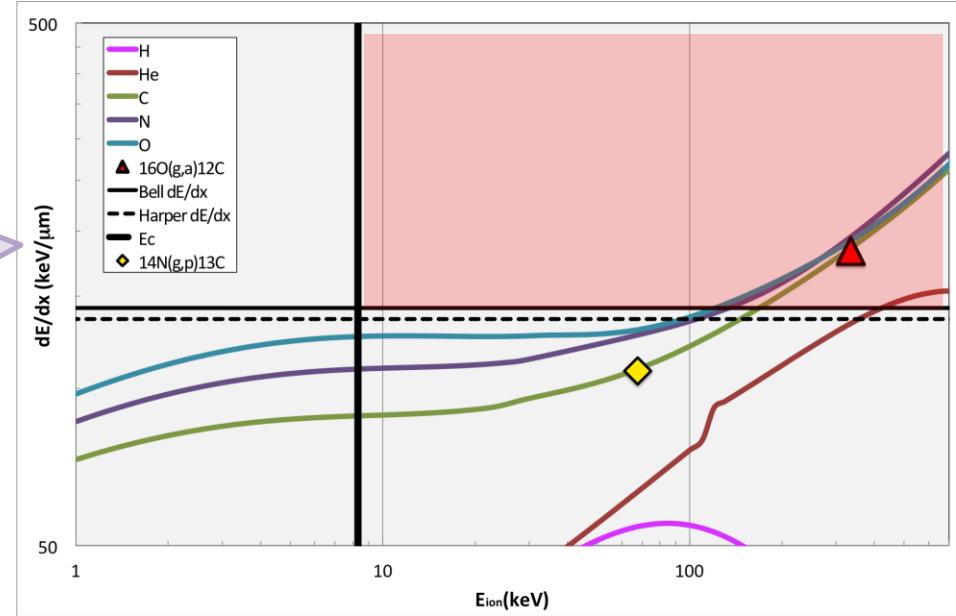
COUPP, FNAL, courtesy of A. Sonnenschein

Suppress neutron events by 100 from acoustic signal

## III. Bubble growth produces an audible click which is recorded by piezo-electric transducers

# EFFICIENCY CURVE

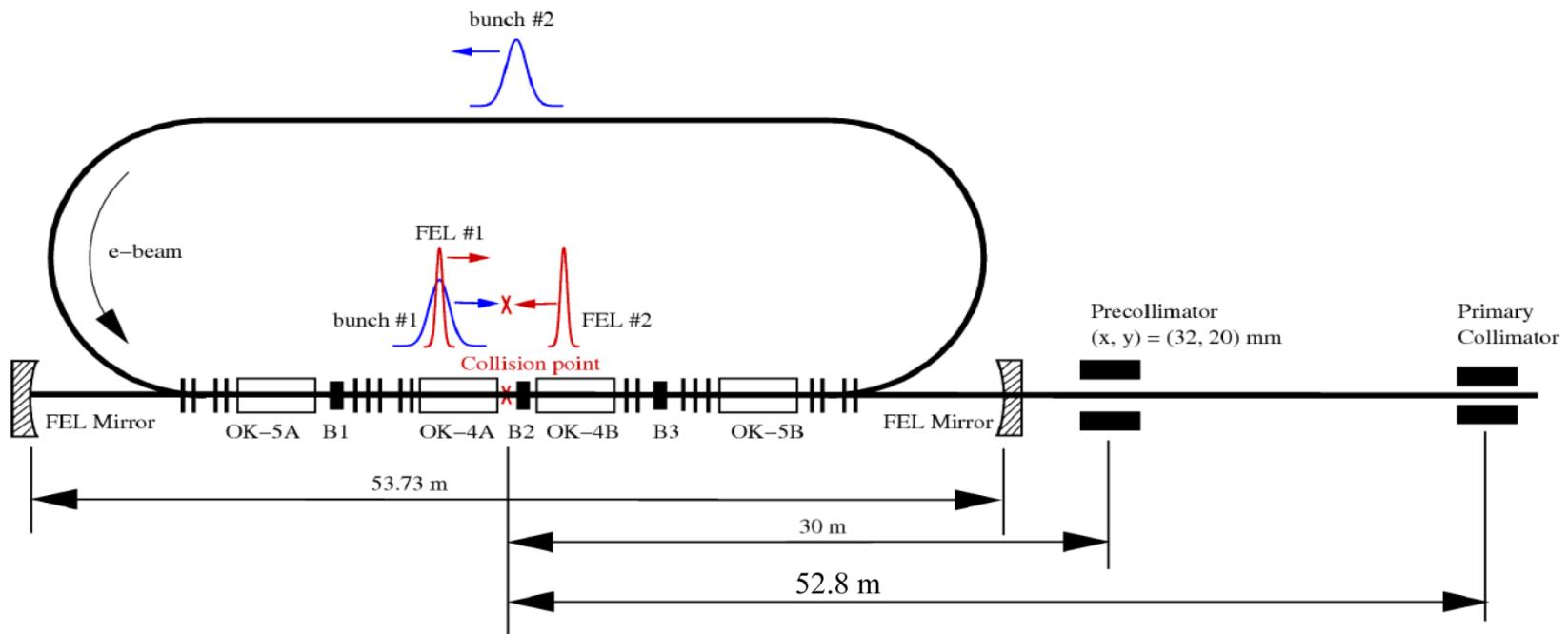
$\text{N}_2\text{O}$  thresholds,  
Superheat =  $3.3^\circ\text{C}$ ,  
 $E_\gamma=8.5 \text{ MeV}$



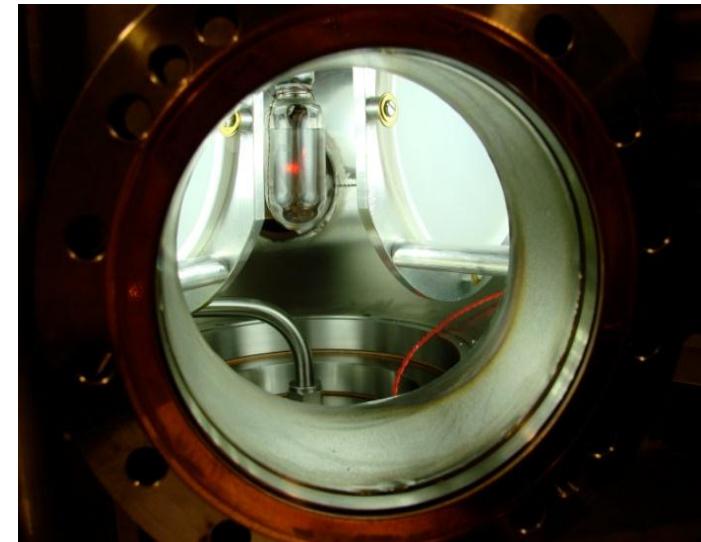
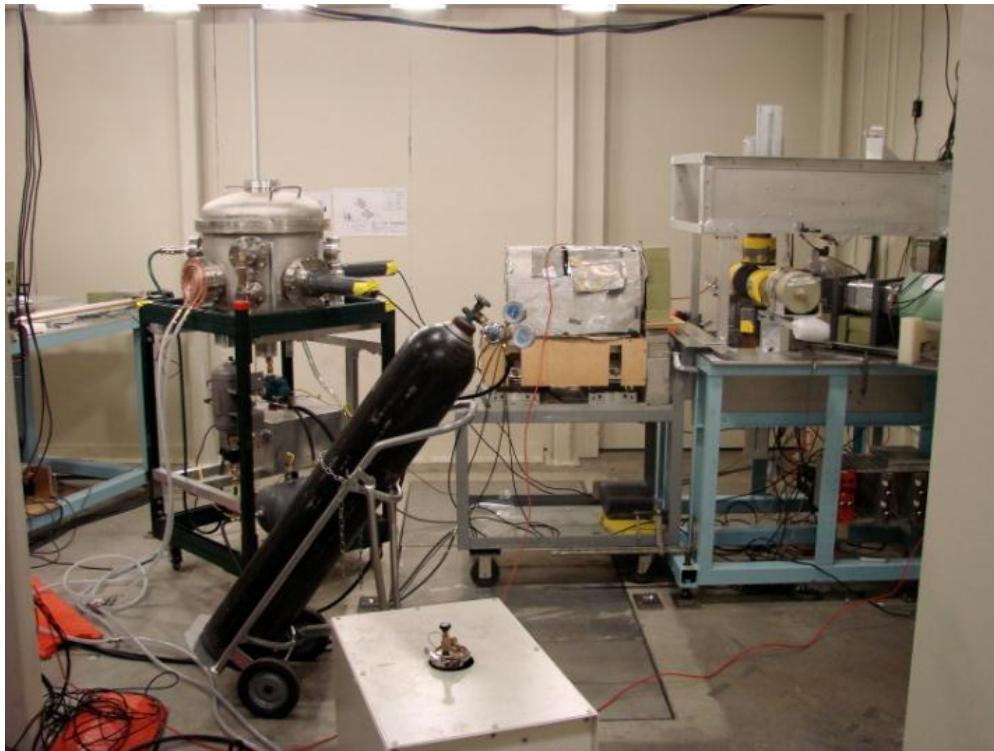
$\text{N}_2\text{O}$  efficiency curve,  
HIGS April 2013,  
 $E_\gamma = 9.7 \text{ MeV}$

# BUBBLE CHAMBER AT HIGS

- I. High Intensity Gamma Source (HIGS) at Duke University
- II.  $\gamma$ -rays generated by Compton backscattering of free-electron-laser (FEL) light from high-energy electron beam bunches



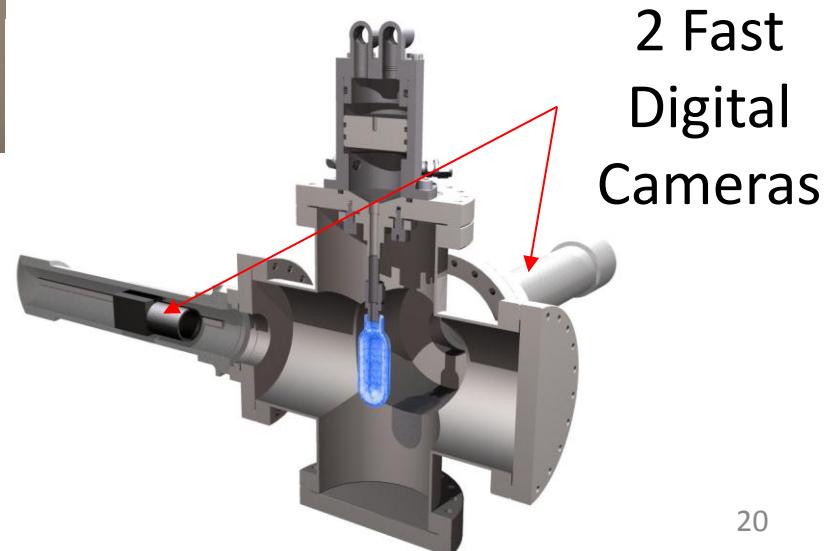
# MEASURING $^{19}\text{F}(\gamma,\alpha)^{15}\text{N}$ AT HIGS



$\text{C}_4\text{F}_{10}$  Bubble Chamber

T =  $30^\circ\text{C}$

P = 3 atm





## First determination of an astrophysical cross section with a bubble chamber: The $^{15}\text{N}(\alpha, \gamma)^{19}\text{F}$ reaction

C. Ugalde <sup>a,\*</sup>, B. DiGiovine <sup>b</sup>, D. Henderson <sup>b</sup>, R.J. Holt <sup>b</sup>, K.E. Rehm <sup>b</sup>, A. Sonnenschein <sup>c</sup>, A. Robinson <sup>d</sup>, R. Raut <sup>e,f,1</sup>, G. Rusev <sup>e,f,2</sup>, A.P. Tonchev <sup>e,f,3</sup>

<sup>a</sup> Department of Astronomy and Astrophysics, University of Chicago, Chicago, IL 60637, USA

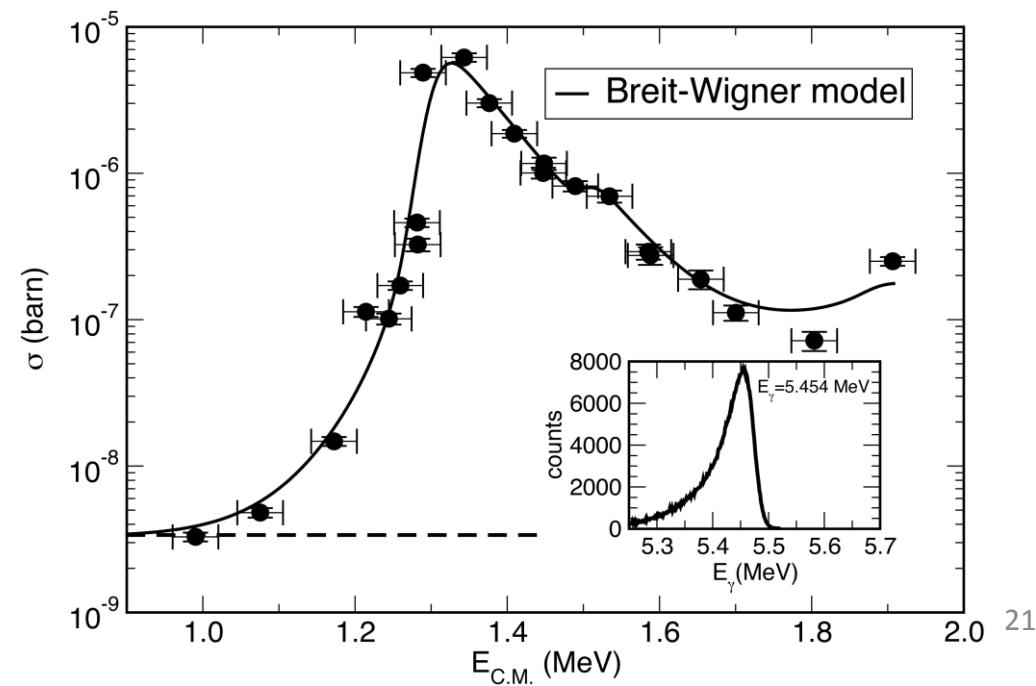
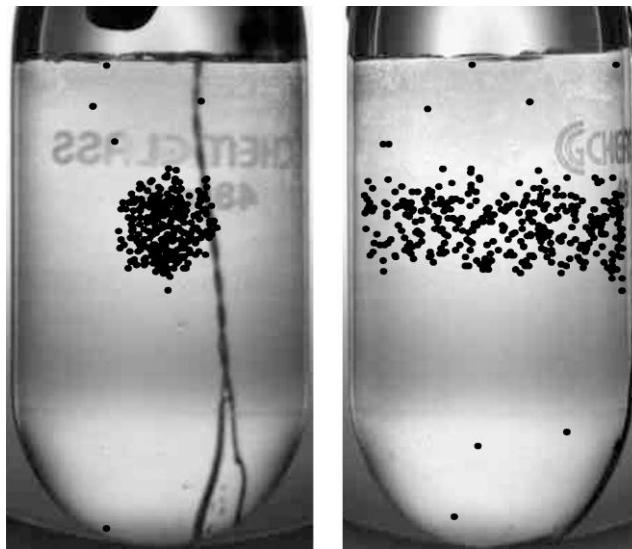
<sup>b</sup> Physics Division, Argonne National Laboratory, Argonne, IL 60439, USA

<sup>c</sup> Fermi National Accelerator Laboratory, Batavia, IL 60510, USA

<sup>d</sup> Department of Physics, University of Chicago, Chicago, IL 60637, USA

<sup>e</sup> Department of Physics, Duke University, Durham, NC 27708, USA

<sup>f</sup> Triangle Universities Nuclear Laboratory, Durham, NC 27708, USA



# BREMSSTRAHLUNG BACKGROUND AT HIGS

Vacuum:  $2 \times 10^{-10}$  Torr

Residual Gas:  $Z = 10$

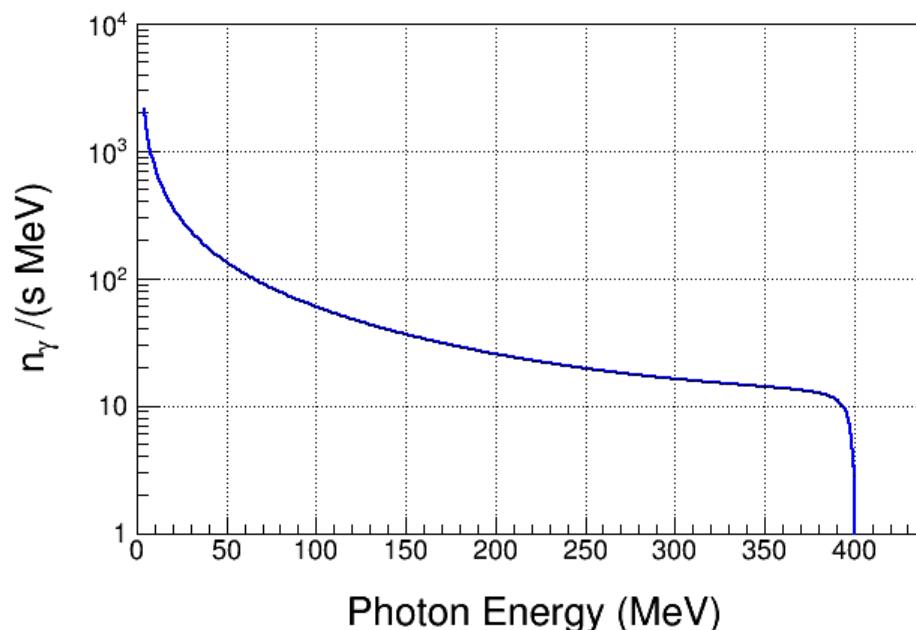
Electron Beam Energy: 400 MeV

Electron Beam Current: 41 mA

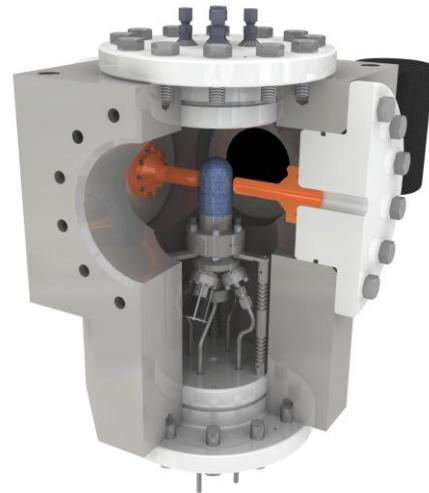
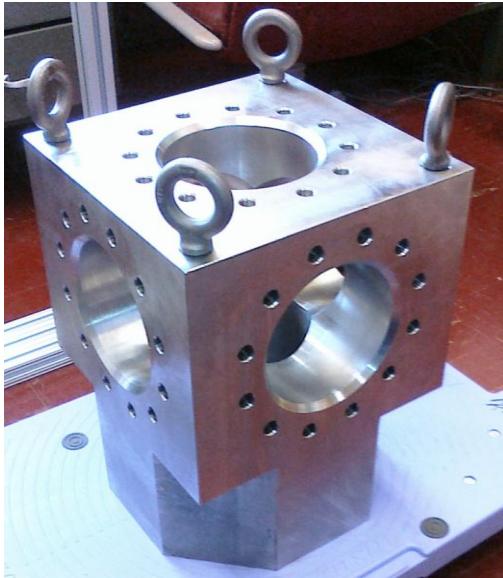
Interaction Length: 35 m



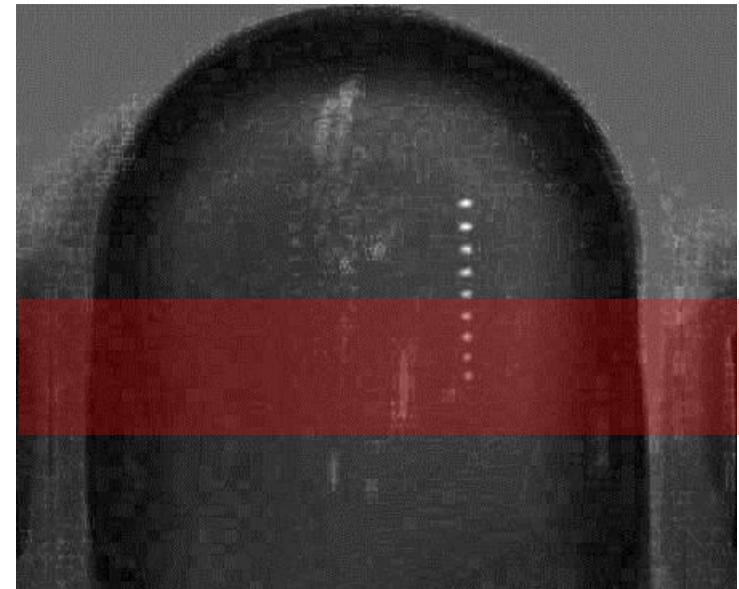
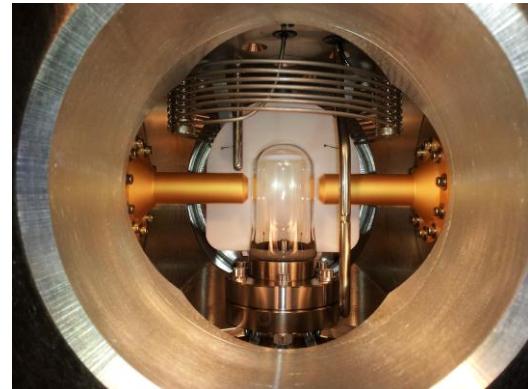
Strong Bremsstrahlung  
Background  
(when coupled with large  
cross sections at high energies)



# RECENT WORK



N<sub>2</sub>O Bubble Chamber:  
first  $\gamma + O \rightarrow \alpha + C$  bubble  
April 2013



# SUPERHEATED TARGETS

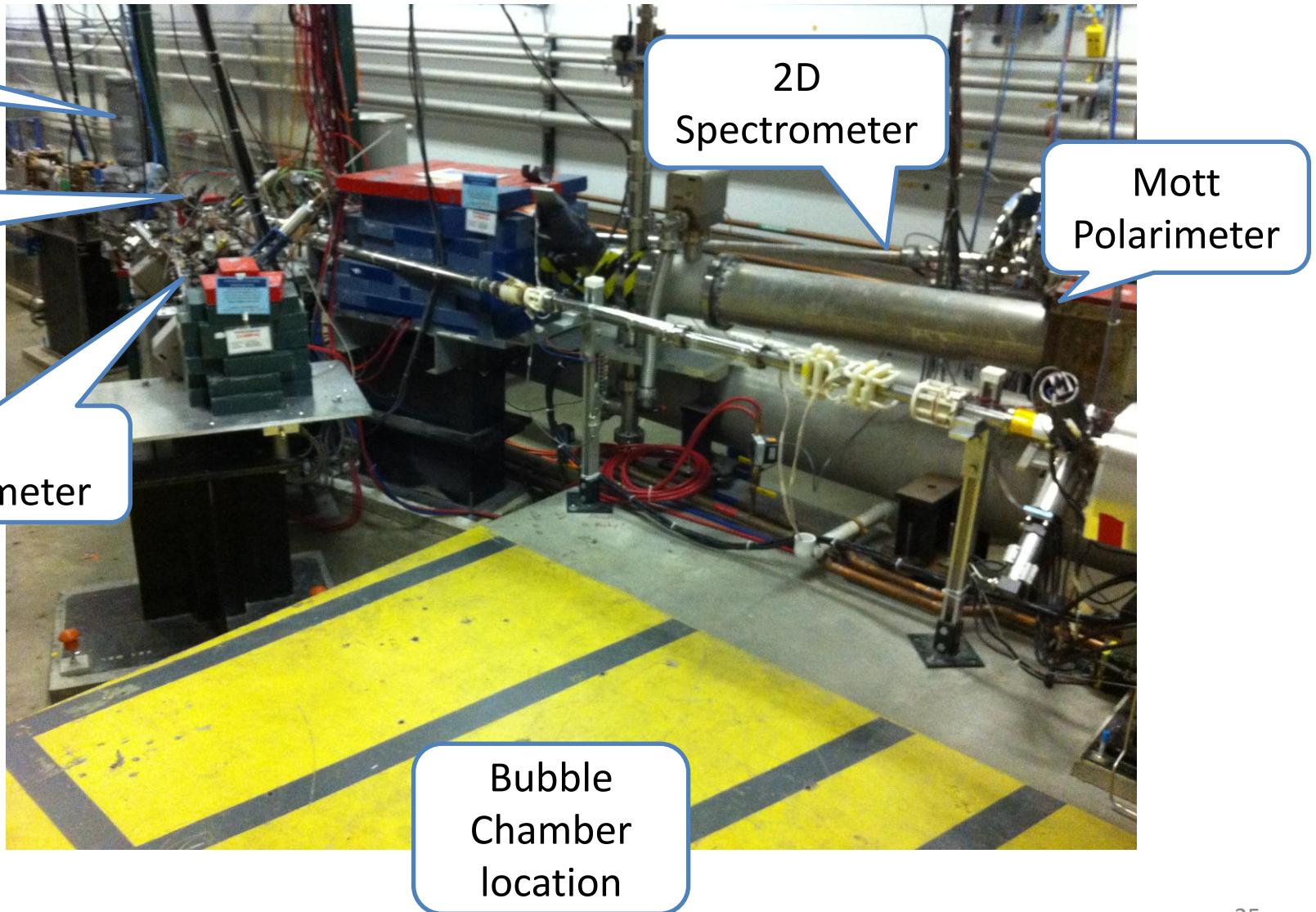
## I. List of superheated liquids to be used in experiment:

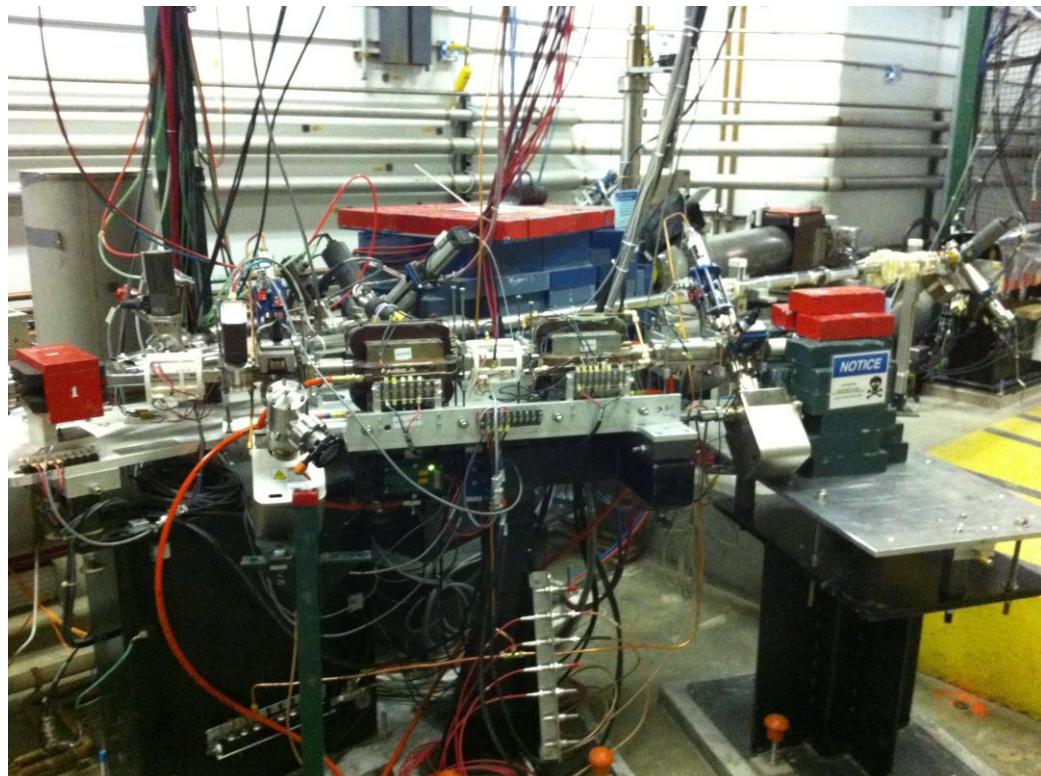
| N <sub>2</sub> O Targets | <sup>16</sup> O | <sup>17</sup> O  | <sup>18</sup> O  |
|--------------------------|-----------------|------------------|------------------|
| Natural Target           | 99.757%         | 0.038%           | 0.205%           |
| <sup>16</sup> O Target   |                 | Depleted > 5,000 | Depleted > 5,000 |
| <sup>17</sup> O Target   |                 | Enriched > 80%   | <1.0%            |
| <sup>18</sup> O Target   |                 | <1.0%            | Enriched > 80%   |

## II. Readout:

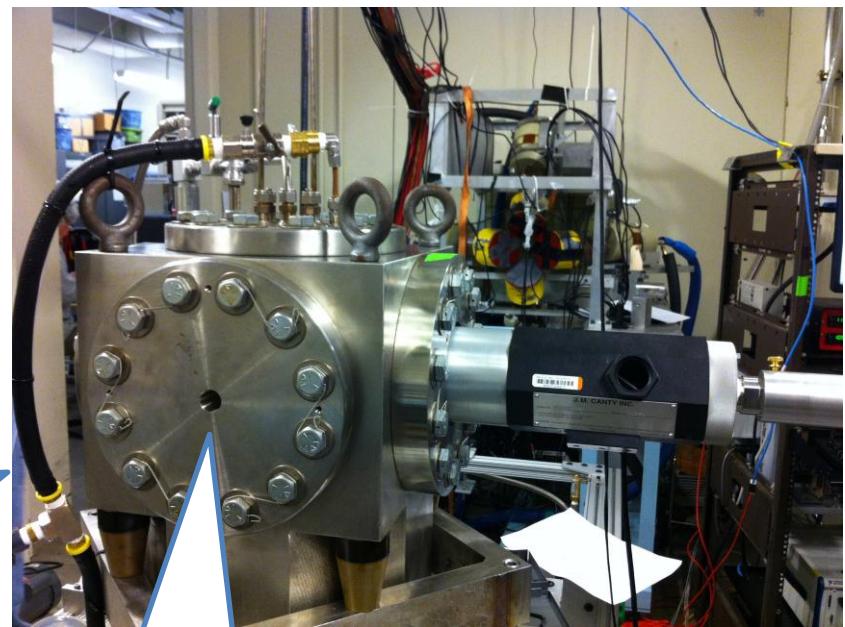
- I. Fast Digital Camera
- II. Acoustic Signal to discriminate between ( $\gamma, \alpha$ ) and ( $\gamma, n$ ) or ( $n, n$ ) events

# EXPERIMENTAL SETUP AT JLAB INJECTOR





5D  
Spectrometer



Bubble  
Chamber at  
HIGS

Photon Beam  
Entrance

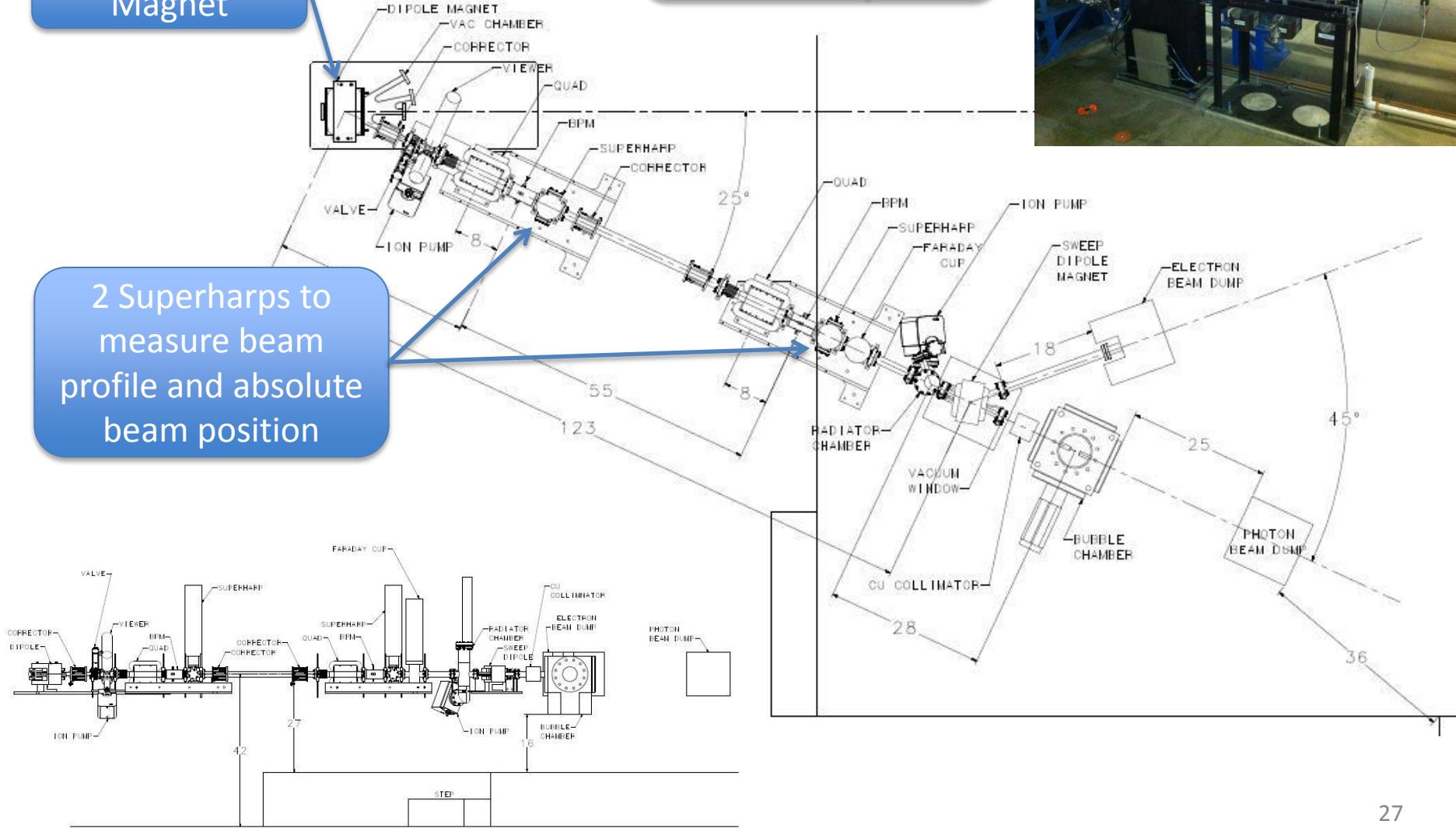
# BEAMLINE

Replace Dipole Magnet

New Fast Valve to protect from vacuum failure in front of  $\frac{1}{4}$  Cryo-unit

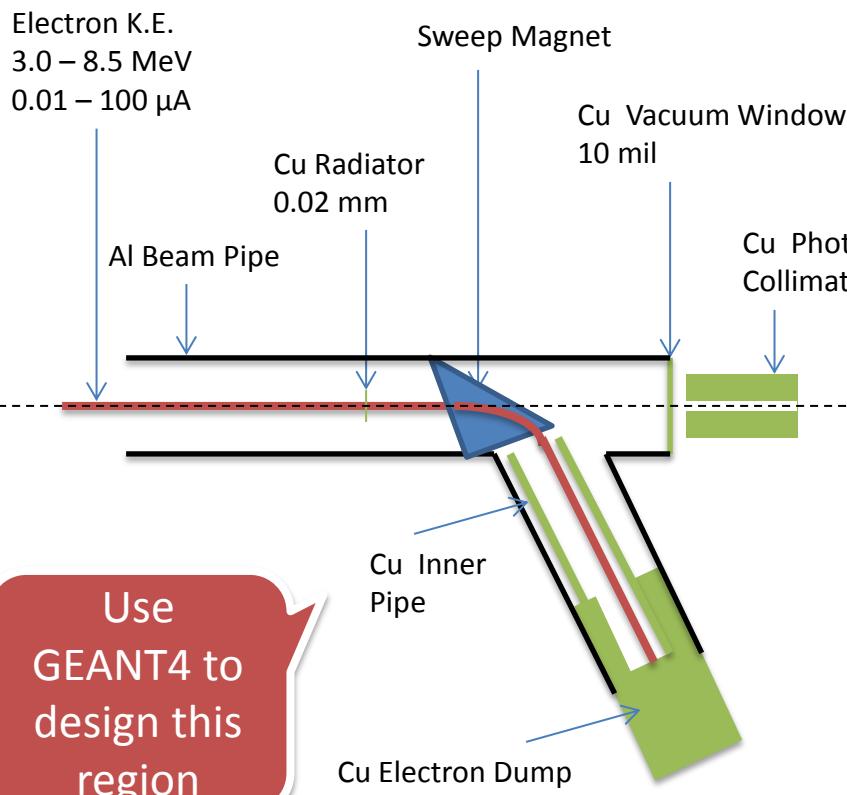


2 Superharps to measure beam profile and absolute beam position

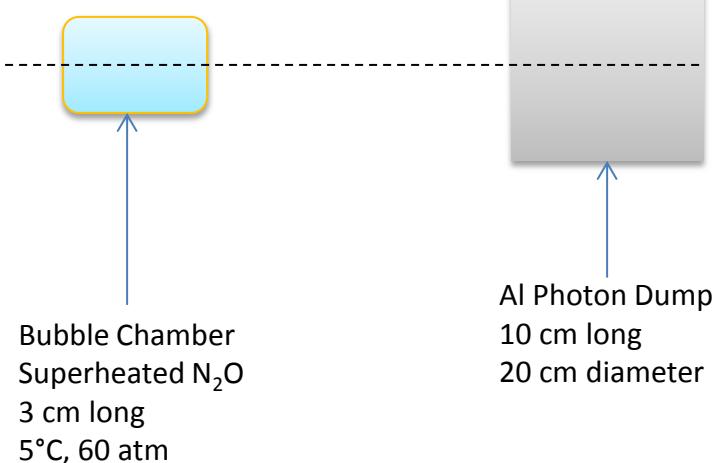


# SCHEMATICS

- Power deposited in radiator (100  $\mu$ A and 8.5 MeV) :
  - I. 0.02 mm: Energy loss = 21 keV, P = 2.1 W
  - II. 0.10 mm: Energy loss = 112 keV, P = 11 W
- Pure Copper and Aluminum (high neutron threshold):
  - I.  $^{63}\text{C}(\gamma, n)$  threshold = 10.86 MeV
  - II.  $^{27}\text{Al}(\gamma, n)$  threshold = 13.06 MeV



- I. Radiator motion and Sweep Dipole current must be in FSD
- II. BCM0L02 and Electron Dump in Beam Loss Accounting (BLA)



# BEAM REQUIREMENTS

## I. Beam Properties at Radiator:

|  |          |
|--|----------|
| Beam Kinetic Energy, (MeV)               | 7.9–8.5  |
| Beam Current ( $\mu\text{A}$ )           | 0.01–100 |
| Absolute Beam Energy                     | <0.1%    |
| Relative Beam Energy                     | <0.02%   |
| Energy Resolution (Spread), $\sigma_T/T$ | 0.06%    |
| Beam Size, $\sigma_{x,y}$ (mm)           | 1–2      |

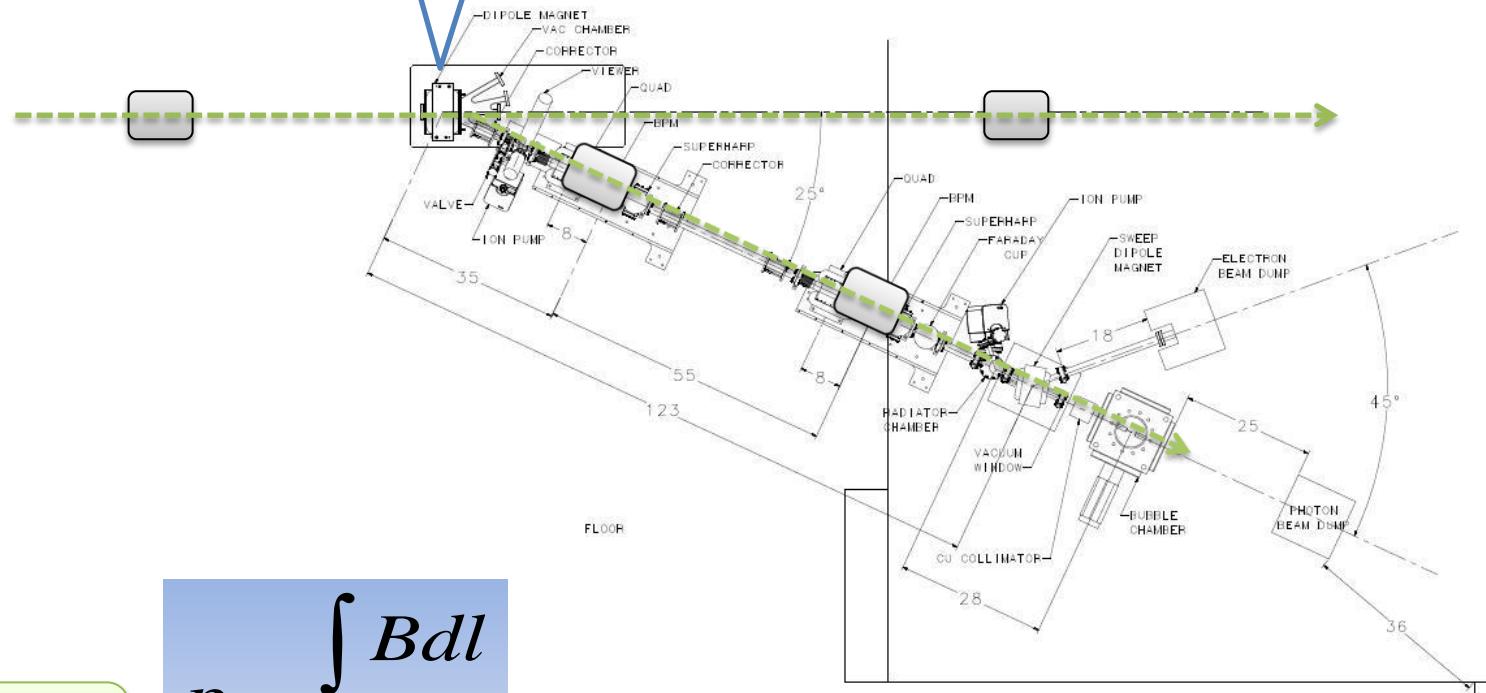
- II. PEPPo achieved  $p=8.25 \text{ MeV}/c$  or  $K.E.=7.75 \text{ MeV}$ . Maximum stable  $\frac{1}{4}$ -cryounit cavity gradients achieved: 8.4 MV/m and 6.1 MV/m (7.25 MV/m average). Vacuum in the beam line indicates that field emission and desorbed gas are the most problematic, but improve with processing.
- III. We may need to helium process the  $\frac{1}{4}$ -cryounit

# ABSOLUTE BEAM ENERGY



BPM

5 MeV  
Dipole



Electron Beam  
Momentum

$$p = \frac{\int B dl}{\theta}$$

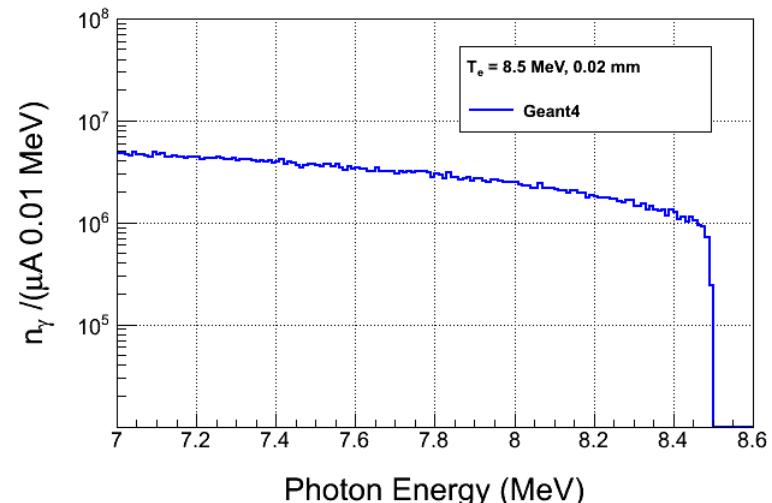
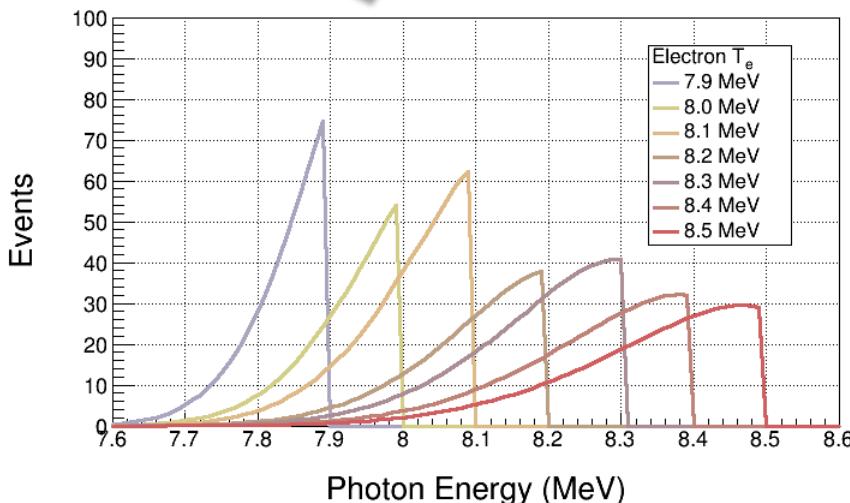
| Parameter                  | Term                           | Now          | Goal             |
|----------------------------|--------------------------------|--------------|------------------|
| Dipole – linearity         | $\delta B/B$                   | 0.25%        | 0.02%            |
| Dipole – spatial           | $\delta BL/BL$                 | 0.10%        | 0.02%            |
| Dipole – reproduce         | $\delta B/B$                   | 0.10%        | 0.02%            |
| Dipole – power supply      | $\delta I/I$                   | 0.20%        | 0.02%            |
| Position – surveys         | $\delta\theta/\theta$          | 0.01%        | 0.01%            |
| Position – BPM calibration | $\delta\theta/\theta$          | 0.05%        | 0.05%            |
| Stray magnetic field       | $\delta\theta/\theta$          | 0.05%        | 0.05%            |
| <b>Total</b>               | <b><math>\delta P/P</math></b> | <b>0.36%</b> | <b>&lt;0.10%</b> |

- I. Jay Benesch designed and now fabricating higher quality dipole (more uniformity, higher field)
- II. New Hall Probe: 0.01% accuracy, resolution to 2 ppm, and a temperature stability of 10 ppm/°C
- III. Relative beam energy error: <0.02%

# BREMSSTRAHLUNG BEAM

- Use both GEANT4 and FLUKA to calculate Bremsstrahlung spectra
- Monte Carlo simulation of Bremsstrahlung at radiotherapy energies is well studied, accuracy:  $\pm 5\%$

Bremsstrahlung Peaks

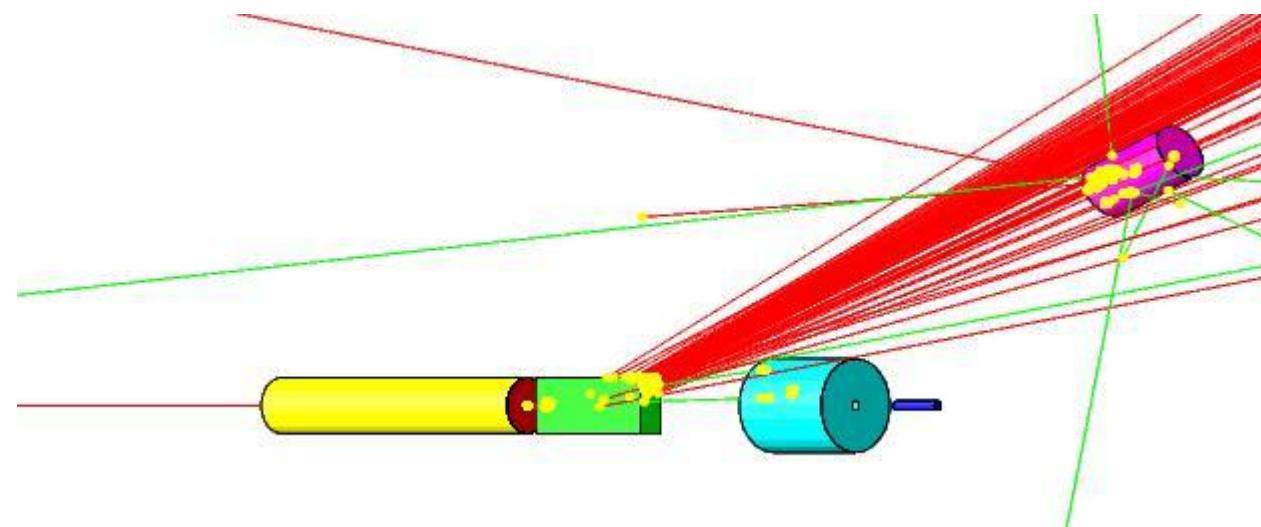


$^{16}\text{O}(\gamma, \alpha)^{12}\text{C}$  is ideal case for Bremsstrahlung beam and Penfold–Leiss Unfolding :

- I. Very steep; only photons near endpoint contribute to yield
- II. No-structure (resonances)

# GEANT4 SIMULATION

- Both GEANT4 and FLUKA use models that calculate wrong photo–nuclear cross sections. Both do not allow for user's cross sections.
  - I. Use GEANT4 and FLUKA to produce the photon spectrum impinging on the superheated liquid.
  - II. Fold the above photon spectrum with our cross sections in stand-alone codes.
- Use GEANT4 to design radiator, collimator, and dumps
- Geometry in GEANT4:



# PENFOLD-LEISS CROSS SECTION UNFOLDING

- Measure Yields at:  $E = E_1, E_2, \dots, E_n$  where,

$$E_i - E_{i-1} = \Delta, i = 2, n$$

$$Y(E_i) = \int_{th}^{E_i} n_\gamma(E_i, k) \sigma(k) dk \approx \sum_{j=1}^i N_\gamma(E_i, \Delta, E_j) \sigma(E_j)$$

- The solution can be written in two forms:

$$\sigma_i = \frac{1}{N_{ii}} \left[ y_i - \sum_{j=1}^{i-1} (N_{ij} \sigma_j) \right]$$

- Or, Matrix form:

$$\begin{bmatrix} y_1 \\ y_2 \\ \vdots \\ y_n \end{bmatrix} = \begin{bmatrix} N_{\gamma,11} & 0 & \cdots & 0 \\ N_{\gamma,21} & N_{\gamma,22} & \cdots & 0 \\ \vdots & \ddots & \ddots & 0 \\ N_{\gamma,n1} & N_{\gamma,n2} & \cdots & N_{\gamma,nn} \end{bmatrix} \begin{bmatrix} \sigma_1 \\ \sigma_2 \\ \vdots \\ \sigma_n \end{bmatrix}$$

$$[Y] = [N] \bullet [\sigma]$$

$$[\sigma] = [N]^{-1} \bullet [Y]$$

# STATISTICAL ERROR PROPAGATION

- Note:  $\frac{dy_i}{y_i} = \frac{1}{\sqrt{y_i}}$        $\frac{dN_{ij}}{N_{ij}} = \frac{1}{\sqrt{N_{ij}}} \approx 0$

$$dy_i = \sqrt{y_i}$$

$$dy_i = \sqrt{y_i + 2y_i^{bg}}$$

In case of  
background  
Subtraction

- With:

$$[B] = [N]^{-1}$$

$$[\sigma] = [B] \bullet [Y]$$

- Then:

$$[d\sigma^2] = [B] \bullet [dY^2] \bullet [B]^T$$

- Where:

$$[dY^2] = \begin{bmatrix} y_1 & 0 & \cdots & 0 \\ 0 & y_2 & \cdots & 0 \\ \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & \cdots & y_n \end{bmatrix}$$

$\text{var}(y_i, y_i) = y_i$   
 $\text{cov}(y_i, y_j) = 0$

$$[d\sigma^2] = \begin{bmatrix} d\sigma_1^2 & \text{cov}(\sigma_1, \sigma_2) & \cdots & \text{cov}(\sigma_1, \sigma_n) \\ \text{cov}(\sigma_2, \sigma_1) & d\sigma_2^2 & \cdots & \text{cov}(\sigma_2, \sigma_n) \\ \vdots & \ddots & \ddots & \vdots \\ \text{cov}(\sigma_n, \sigma_1) & \text{cov}(\sigma_n, \sigma_2) & \cdots & d\sigma_n^2 \end{bmatrix}$$

Although,  
 $\text{cov}(y_i, y_j) = 0$ ,  
 $\text{cov}(\sigma_i, \sigma_j) \neq 0$

$$(d\sigma_i)^2 = \frac{1}{N_{ii}^2} \left[ dy_i^2 + \sum_{j=1}^{i-1} (N_{ij} d\sigma_j)^2 + \sum_{k=1}^{i-1} \sum_{l=1}^{i-1} N_{ik} \text{cov}(\sigma_k, \sigma_l) N_{il} \right]$$

For mono-chromatic beam

$$\left( \frac{d\sigma_i}{\sigma_i} \right)^2 = \left( \frac{dy_i}{y_i} \right)^2 = \frac{1}{y_i}$$

# RESULTS

- I. Radiator Thickness = 0.02 mm
- II. Bubble Chamber Thickness = 3.0 cm. Number of  $^{16}\text{O}$  nuclei =  $3.474 \cdot 10^{22} / \text{cm}^2$
- III. Background subtraction of  $^{18}\text{O}(\gamma, \alpha)^{14}\text{C}$

$$[N] = \begin{bmatrix} 3.267e14 & 0 & 0 & 0 & 0 & 0 & 0 \\ 9.782e13 & 6.439e13 & 0 & 0 & 0 & 0 & 0 \\ 5.013e13 & 3.858e13 & 2.539e13 & 0 & 0 & 0 & 0 \\ 1.494e13 & 1.236e13 & 9.514e12 & 6.258e12 & 0 & 0 & 0 \\ 8.540e12 & 7.369e12 & 6.097e12 & 4.692e12 & 3.086e12 & 0 & 0 \\ 3.801e12 & 3.370e12 & 2.908e12 & 2.406e12 & 1.852e12 & 1.217e12 & 0 \\ 2.075e12 & 1.875e12 & 1.663e12 & 1.435e12 & 1.187e12 & 9.137e11 & 6.004e11 \end{bmatrix}$$

| Electron Beam K. E. | Beam Current ( $\mu\text{A}$ ) | Time (hour) | $y_i$ | $dy_i$ (no bg) | $dy_i/y_i$ (no bg, %) | $dy_i$ (with bg) | $dy_i/y_i$ (with bg, %) |
|---------------------|--------------------------------|-------------|-------|----------------|-----------------------|------------------|-------------------------|
| 7.9                 | 100                            | 100         | 545   | 23             | 4.2                   | 134              | 24.6                    |
| 8.0                 | 100                            | 20          | 581   | 24             | 4.1                   | 77               | 13.3                    |
| 8.1                 | 80                             | 10          | 852   | 29             | 3.4                   | 60               | 7.0                     |
| 8.2                 | 20                             | 10          | 634   | 25             | 3.9                   | 40               | 6.3                     |
| 8.3                 | 10                             | 10          | 812   | 28             | 3.4                   | 39               | 4.8                     |
| 8.4                 | 4                              | 10          | 746   | 27             | 3.6                   | 36               | 4.8                     |
| 8.5                 | 2                              | 10          | 763   | 28             | 3.7                   | 32               | 4.2                     |

# SYSTEMATIC ERROR PROPAGATION

- For absolute beam energy uncertainty of  $\delta E$  ( $= 0.1\%$ ) and zero relative beam energy uncertainty:

$$\frac{dy_i}{y_i} = \frac{y_i(E_i + \delta E) - y_i(E_i)}{y_i(E_i)}$$

$$\frac{dN_{ij}}{N_{ij}} = \frac{N_{ij}(E_i + \delta E) - N_{ij}(E_i)}{N_{ij}(E_i)}$$

$$E_0 = 7.8 + \delta E$$

$$E_i = E_0 + i\Delta$$

| $E_i$ (MeV) | $dy_i/y_i$ (%) | $d\sigma_i/\sigma_i$ (%) |
|-------------|----------------|--------------------------|
| 7.9         | 12.5           | 12.6                     |
| 8.0         | 10.8           | 10.5                     |
| 8.1         | 9.3            | 9.1                      |
| 8.2         | 8.0            | 7.1                      |
| 8.3         | 7.0            | 6.3                      |
| 8.4         | 6.3            | 5.8                      |
| 8.5         | 5.6            | 5.2                      |

This is the cross section dependence on energy

- Accounted for  $dN_{ij}$  due to energy error when calculating  $dy_i$

$$\approx \frac{\delta E}{i\Delta}$$

$$[dN_{ij} / N_{ij}] = \begin{bmatrix} 0.100 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0.058 & 0.050 & 0 & 0 & 0 & 0 & 0 \\ 0.041 & 0.039 & 0.033 & 0 & 0 & 0 & 0 \\ 0.031 & 0.031 & 0.029 & 0.025 & 0 & 0 & 0 \\ 0.025 & 0.025 & 0.025 & 0.023 & 0.020 & 0 & 0 \\ 0.021 & 0.021 & 0.021 & 0.021 & 0.020 & 0.017 & 0 \\ 0.018 & 0.018 & 0.018 & 0.018 & 0.018 & 0.017 & 0.022 \end{bmatrix}$$

- With:

$$[B] = [N]^{-1}$$

$$[\sigma] = [B] \bullet [Y]$$

- Then:

$$[d\sigma^2] = [B] \bullet ([dY^2] + [dN^2] \bullet [\sigma^2]) \bullet [B]^T$$

- Where:

Note: Correlation Coefficient =1

$$\text{var}(y_i, y_i) = (dy_i)^2$$

$$\text{cov}(y_i, y_j) = \rho_{ij} dy_i dy_j$$

$$[dY^2] = \begin{bmatrix} (dy_1)^2 & dy_1 dy_2 & \cdots & dy_1 dy_n \\ dy_2 dy_1 & (dy_2)^2 & \cdots & dy_n dy_n \\ \vdots & \ddots & \ddots & \vdots \\ dy_n dy_1 & dy_n dy_2 & \cdots & (dy_n)^2 \end{bmatrix}$$

No point-to-point systematic

$$[d\sigma^2] = \begin{bmatrix} d\sigma_1^2 & \text{cov}(\sigma_1, \sigma_2) & \cdots & \text{cov}(\sigma_1, \sigma_n) \\ \text{cov}(\sigma_2, \sigma_1) & d\sigma_2^2 & \cdots & \text{cov}(\sigma_2, \sigma_n) \\ \vdots & \ddots & \ddots & \vdots \\ \text{cov}(\sigma_n, \sigma_1) & \text{cov}(\sigma_n, \sigma_2) & \cdots & d\sigma_n^2 \end{bmatrix}$$

$$[dN^2] = \begin{bmatrix} (dN_{11})^2 & 0 & \cdots & 0 \\ (dN_{21})^2 & (dN_{22})^2 & \cdots & 0 \\ \vdots & \ddots & \ddots & \vdots \\ (dN_{n1})^2 & (dN_{n2})^2 & \cdots & (dN_{nn})^2 \end{bmatrix}$$

$$[\sigma^2] = \begin{bmatrix} \sigma_1^2 & 0 & \cdots & 0 \\ 0 & \sigma_2^2 & \cdots & 0 \\ \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & \cdots & \sigma_n^2 \end{bmatrix}$$

# SYSTEMATIC ERROR PROPAGATION

$$\begin{aligned}(d\sigma_i)^2 \simeq & \frac{1}{N_{ii}^2} \left[ dy_i^2 - 2dy_i \sum_{j=1}^{i-1} N_{ij} d\sigma_j \right. \\ & + \sum_{j=1}^{i-1} (N_{ij} d\sigma_j)^2 + \sum_{k=1}^{i-1} \sum_{l=1}^{i-1} N_{ik} \text{cov}(\sigma_k, \sigma_l) N_{il} \\ & \left. + \sum_{j=1}^{i-1} (dN_{ij} \sigma_j)^2 + (dN_{ii} \sigma_i)^2 \right]\end{aligned}$$

No point-to-point systematic

$\text{cov}(y_i, y_j) \neq 0,$   
 $\text{cov}(\sigma_i, \sigma_j) \neq 0$

# OTHER SYSTEMATIC ERRORS

|  |    |
|--|----|
| Beam Current, $\delta I/I$               | 3% |
| Photon Flux, $\delta\varphi/\varphi$     | 5% |
| Radiator Thickness, $\delta R/R$         | 3% |
| Bubble Chamber Thickness, $\delta T/T$   | 3% |
| Bubble Chamber Efficiency, $\varepsilon$ | 5% |

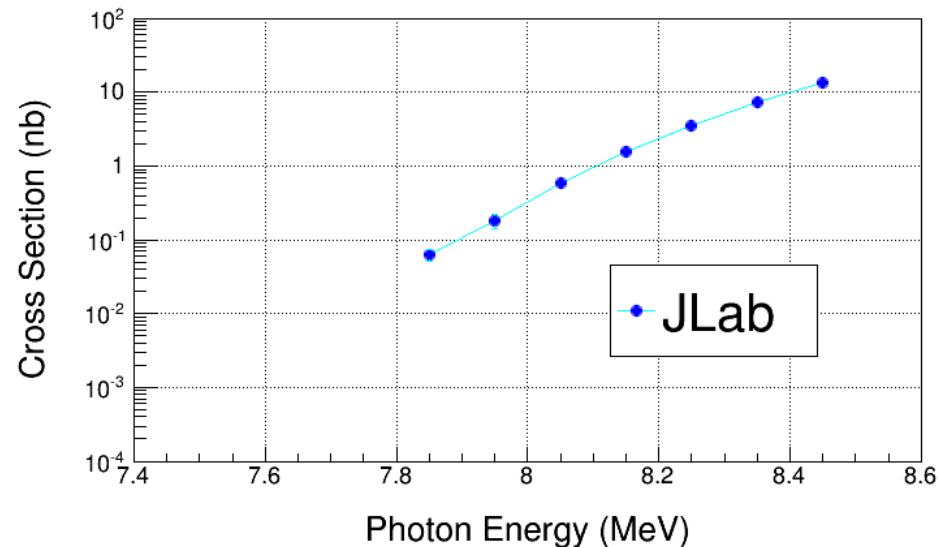
Simulation

- Then:

$$(dy_i)^2 = (dy_i(\delta E))^2 + \left[ \left( \frac{\delta I}{I} \right)^2 + \left( \frac{\delta R}{R} \right)^2 + \left( \frac{\delta T}{T} \right)^2 + \varepsilon^2 \right] y_i^2$$

$$(dN_{ij})^2 = \left( \frac{\delta\phi}{\phi} \right)^2 N_{ij}^2$$

| Electron Beam K. E. | Cross Section (nb) | Stat Error (no bg, %) | Stat Error (with bg, %) |
|---------------------|--------------------|-----------------------|-------------------------|
| 7.9                 | 0.046              | 4.4                   | 24.5                    |
| 8.0                 | 0.185              | 6.0                   | 20.7                    |
| 8.1                 | 0.58               | 6.3                   | 14.7                    |
| 8.2                 | 1.53               | 8.2                   | 13.8                    |
| 8.3                 | 3.49               | 9.1                   | 13.3                    |
| 8.4                 | 7.2                | 10.6                  | 13.8                    |
| 8.5                 | 13.6               | 12.2                  | 14.8                    |



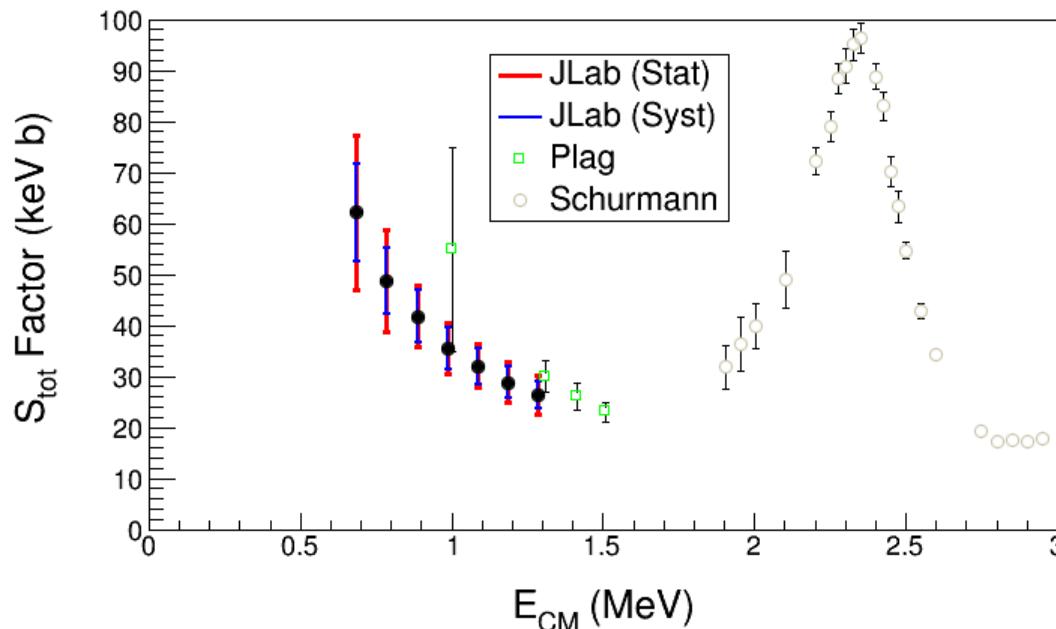
| Electron Beam K. E. | Cross Section (nb) | Sys Error (Energy, %) | Sys Error (Total, %) |
|---------------------|--------------------|-----------------------|----------------------|
| 7.9                 | 0.046              | 12.5                  | 15.3                 |
| 8.0                 | 0.185              | 10.2                  | 13.5                 |
| 8.1                 | 0.58               | 8.3                   | 12.2                 |
| 8.2                 | 1.53               | 7.0                   | 11.4                 |
| 8.3                 | 3.49               | 6.0                   | 10.7                 |
| 8.4                 | 7.2                | 5.3                   | 10.5                 |
| 8.5                 | 13.6               | 4.7                   | 10.1                 |

Note: Relative systematic errors do not get amplified in PL Unfolding

# JLAB PROJECTED $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$ S-Factor

- Statistical Error: dominated by background subtraction from  $^{18}\text{O}(\gamma, \alpha)^{14}\text{C}$  (depletion = 5,000)

| Electron Beam K. E. | Gamma Energy (MeV) | $E_{CM}$ (MeV) | Cross Section (nb) | $S_{tot}$ Factor (keV b) | Stat Error (%) | Sys Error (Total, %) |
|---------------------|--------------------|----------------|--------------------|--------------------------|----------------|----------------------|
| 7.9                 | 7.85               | 0.69           | 0.046              | 62.2                     | 24.5           | 15.3                 |
| 8.0                 | 7.95               | 0.79           | 0.185              | 48.7                     | 20.7           | 13.5                 |
| 8.1                 | 8.05               | 0.89           | 0.58               | 41.8                     | 14.7           | 12.2                 |
| 8.2                 | 8.15               | 0.99           | 1.53               | 35.5                     | 13.8           | 11.4                 |
| 8.3                 | 8.25               | 1.09           | 3.49               | 32.0                     | 13.3           | 10.7                 |
| 8.4                 | 8.35               | 1.19           | 7.2                | 28.8                     | 13.8           | 10.5                 |
| 8.5                 | 8.45               | 1.29           | 13.6               | 26.3                     | 14.8           | 10.1                 |



Bubble Chamber experiment measures total S-Factor,  $S_{E1} + S_{E2}$

# BACKGROUNDS

## I. Background from oxygen isotopes and nitrogen in N<sub>2</sub>O:

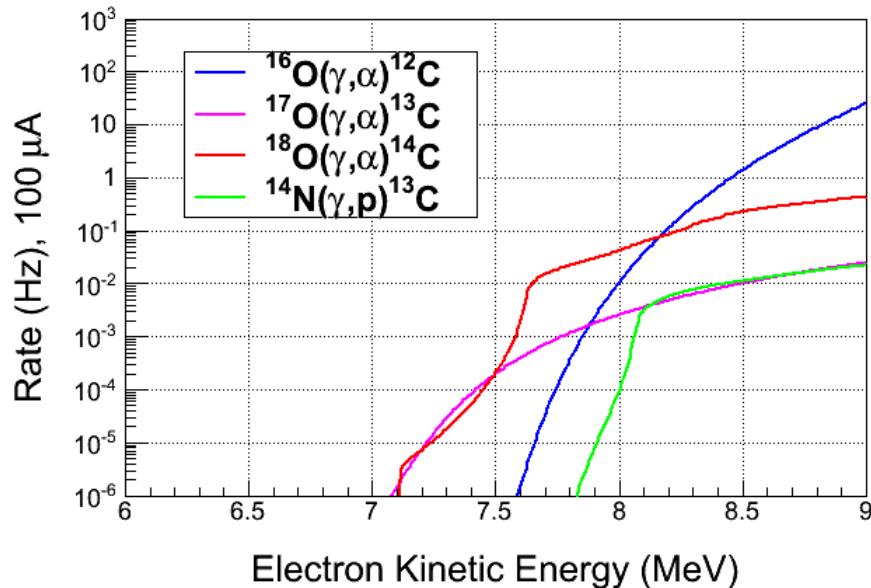
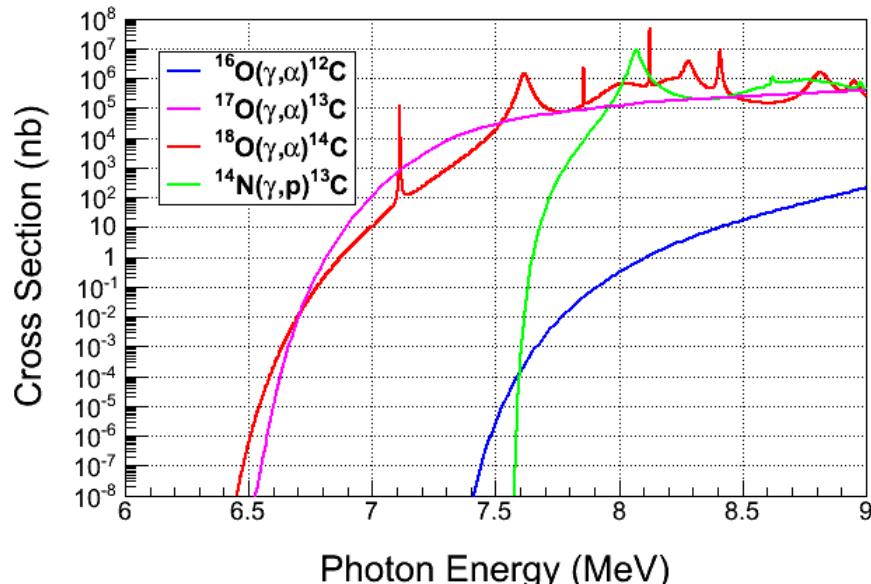
- $^{18}\text{O}(\gamma,\alpha)^{14}\text{C}$
- $^{17}\text{O}(\gamma,\alpha)^{13}\text{C}$
- $^{14}\text{N}(\gamma,p)^{13}\text{C}$

### ➤ Natural Abundance:

- I.  $^{17}\text{O}$ : 0.038%
- II.  $^{18}\text{O}$ : 0.205%

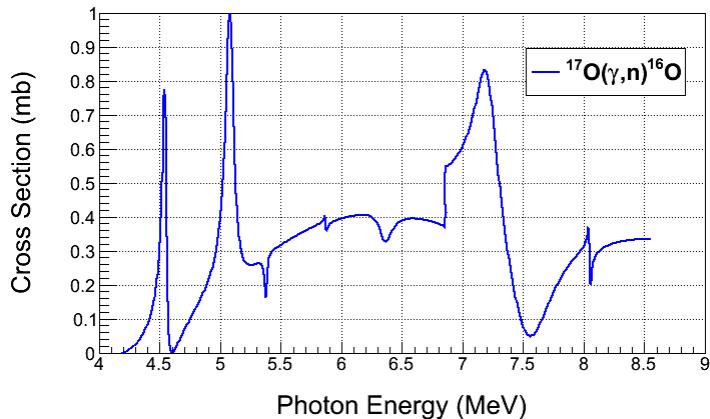
### ➤ Expected Rates:

- I.  $^{17}\text{O}(\gamma,\alpha)^{13}\text{C}$ , depletion=5,000
- II.  $^{18}\text{O}(\gamma,\alpha)^{14}\text{C}$ , depletion=5,000
- III.  $^{14}\text{N}(\gamma,p)^{13}\text{C}$ , Chamber eff.=  $10^{-8}$



## II. Background from:

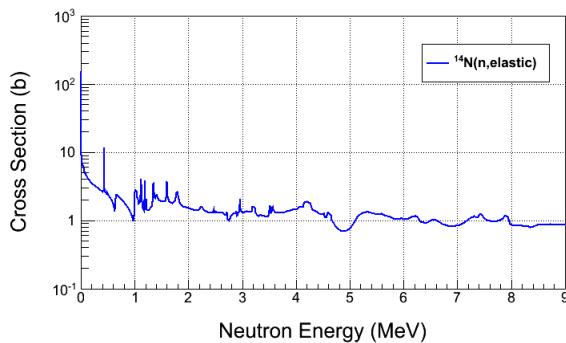
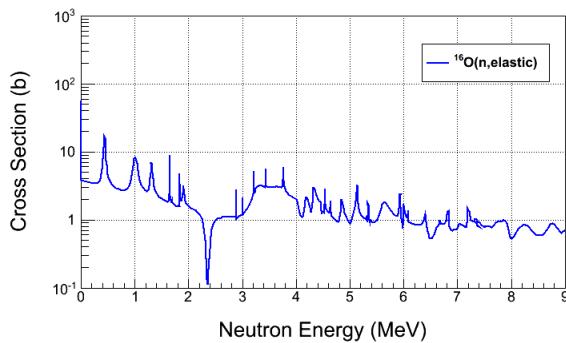
- $^{17}\text{O}(\gamma, \text{n})^{16}\text{O}$  and secondary ( $\text{n}, \text{n}$ ) neutron–nucleus elastic scattering



## III. Cosmic-ray background:

- $\mu^\pm$ –nuclear
- neutron–nuclear elastic scattering

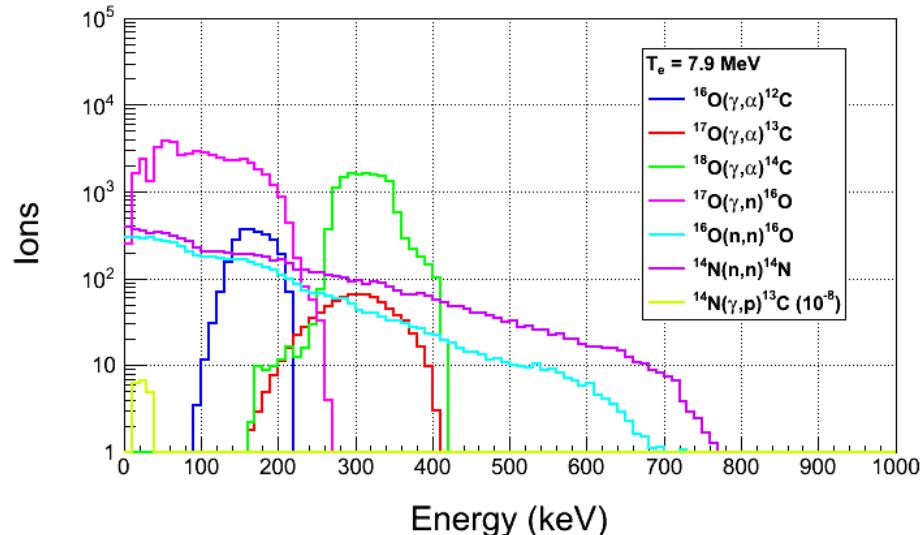
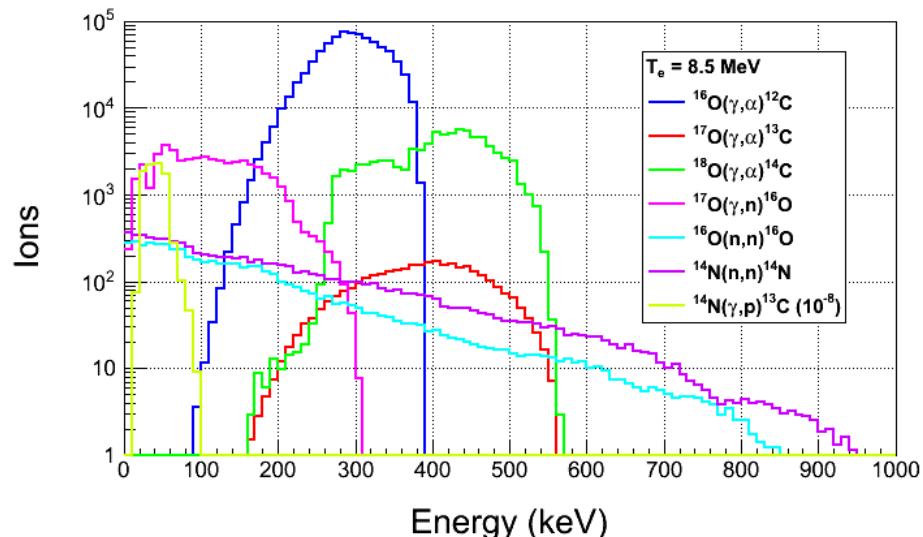
➤ Reject neutron background using acoustic signal (100 suppression factor)



# ION ENERGY DISTRIBUTIONS

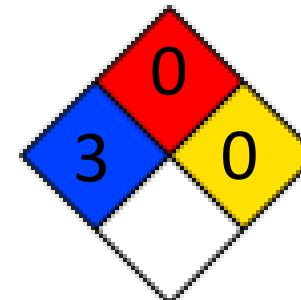
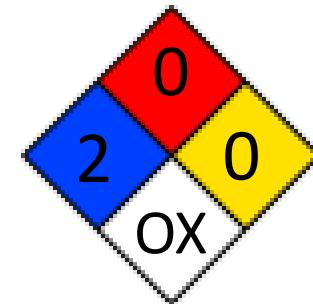
- Suppress background with Bubble Chamber threshold
  
  
  
- Calculated with Depletion:
  - I.  $^{17}\text{O}$  depletion = 5,000
  - II.  $^{18}\text{O}$  depletion = 5,000
  
  
  
- Threshold Efficiency (function of superheat):

| Particle            | Efficiency  |
|---------------------|-------------|
| $e^\pm$             | $<10^{-11}$ |
| $\gamma$            | $<10^{-11}$ |
| $(\gamma,n), (n,n)$ | $<10^{-2}$  |



# SAFETY

- Super heated liquid N<sub>2</sub>O, Nitrous oxide (laughing gas)
  - I. At room temperature, it is a colorless, non-flammable gas, with a slightly sweet odor and taste
- High pressure system:
  - I. Design Authority: Dave Meekins
  - II. T = 5°C
  - III. P = 60 atm
- Buffer liquid: Mercury
  - I. Closed system
  - II. Volume: 135 mL



# SUMMARY AND OUTLOOK

- Test N<sub>2</sub>O Bubble Chamber at HIGS (February 2014)
- Measure cross sections of <sup>18</sup>O(γ,α)<sup>14</sup>C and <sup>17</sup>O(γ,α)<sup>13</sup>C at HIGS (Summer 2014)
- Test Bubble Chamber at JLab with Bremsstrahlung beam (October 2014)
- If successful, run depleted N<sub>2</sub>O bubble chamber at JLab <sup>16</sup>O(γ,α)<sup>12</sup>C
- Beam issues:
  - Design radiator, collimator, and dumps with GEANT4
  - Simulate Photon Spectrum with GEANT4 and FLUKA
  - Deliver 8.5 MeV K.E. electron beam to 5D Spectrometer with <0.1% energy uncertainty
- Bubble Chamber issues:
  - Study acoustic signal and measure neutron suppression factor
  - Deadtime measurement (now  $\tau \pm d\tau = 10.0 \pm 0.9$  sec)
  - Measure O-isotopes depletion
- Background tests:
  - Measure cosmic-ray background
  - Study chamber efficiency vs. superheat and measure γ-rays suppression factor

# BACKUP SLIDES

# COST ESTIMATE

- I. New beamline components:
  - I. New Dipole Magnet and Hall Probe
  - II. 2 Super Harps
  - III. Fast Valve
- II. Summary of labor cost by group:

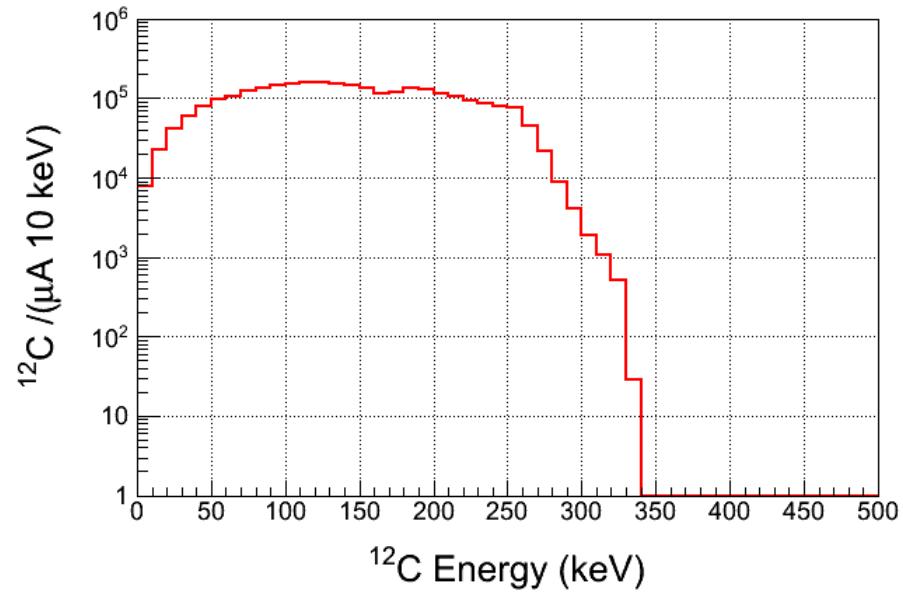
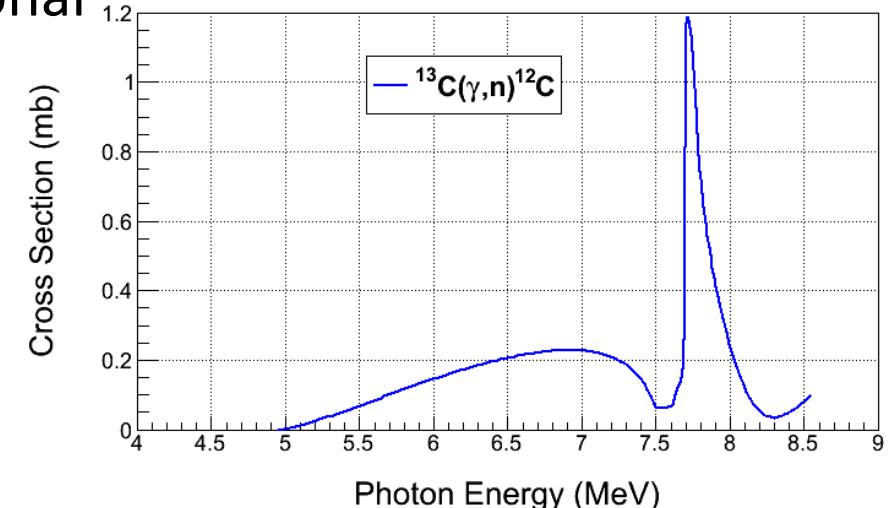
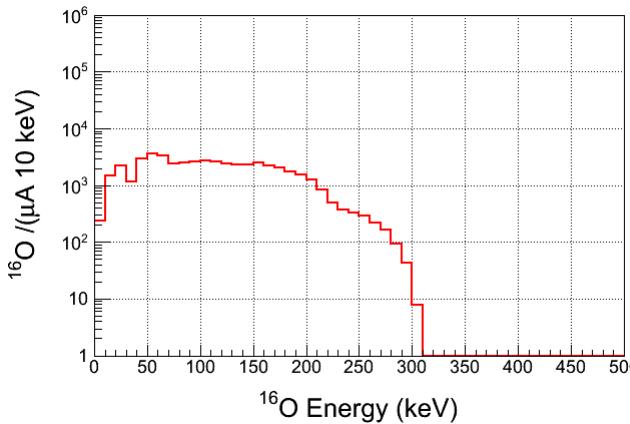
| Group              | Labor     |
|--------------------|-----------|
| Survey & Alignment | 3 wks x 2 |
| Magnet Test        | 1 wk x 2  |
| Engineering Design | 16 wks    |
| Software           | 3 wks x 2 |
| EES                | 6 wk x 2  |
| EH&Q               | 4 wks     |

| Item                            | Material Procurement                                    | Shop                         | Labor   |
|---------------------------------|---|------------------------------|---|
| New Dipole Magnet               | Dipole Magnet (\$8,000)<br>Hall Probe System (\$10,000) |                              | Design (2 week)<br>Mapping (1 week)<br>EESDC (1 week)<br>Alignment (2 days)   |
| New Beamline                    | 2 Super Harps (20,000)<br>Fast Valve (\$23,000)         | Pipes + Pedestals (\$20,000) | Design (6 weeks)<br>Alignment (1 week)<br>Software (6 weeks)<br>EES (6 weeks) |
| Radiator (cooled ladder, FSD)   | 0.02 and 0.10 mm Cu foils (\$2,000)                     | \$4,000                      | Design (2 week)<br>Alignment (2 days)   |
| Sweep Dipole                    |   |                              |   |
| Electron Dump                   | Pure Cu (\$5,000)                                       | Dump + Pipes (\$15,000)      | Design (4 weeks)<br>Alignment (1 day)   |
| Cu Collimator                   | Pure Cu (\$5,000)                                       | Collimator + Stand (\$5,000) | Design (1 week)<br>Alignment (1 day)  |
| Photon Dump & Stand             | Pure Al (\$3,000)                                       | \$4,000                      | Design (1 week)<br>Alignment (1 day)  |
| Safety Review                   |   |                              | 4 weeks   |
| Install                         |   |                              | 6 weeks   |
| Bubble Chamber                  |   |                              | Alignment (1 week)  |
| Total                           | \$76,000  | \$48,000                     | \$80,000  |
| Indirect G&A (55.65%)           | \$42,300  | \$26,400                     | \$42,500  |
| Indirect Stat & Fringe (57.15%) |   |                              | \$45,700  |
| <b>Total</b>                    | <b>\$118,300</b>  | <b>\$74,400</b>              | <b>\$168,200</b>  |

# $\text{CO}_2$ SUPERHEATED LIQUID?

- Similar Bubble Chamber operational parameters as  $\text{N}_2\text{O}$
- Natural Abundance:  $^{13}\text{C}$ : 1.07%
- Depletion:  $^{13}\text{C}$  depletion=1,000
- $^{13}\text{C}(\gamma, \text{n})^{12}\text{C}$  Background

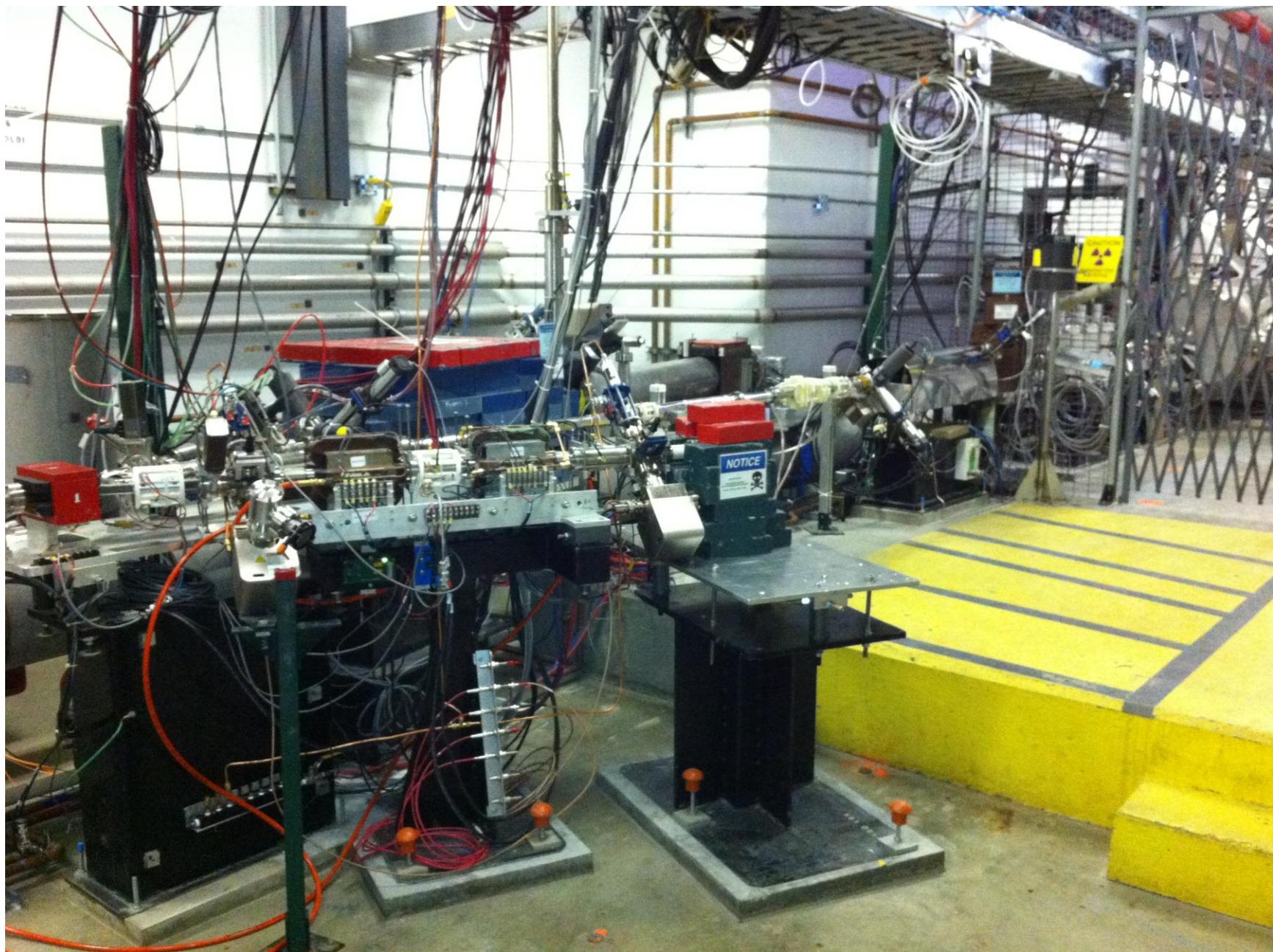
For comparison,  $^{17}\text{O}(\gamma, \text{n})^{16}\text{O}$

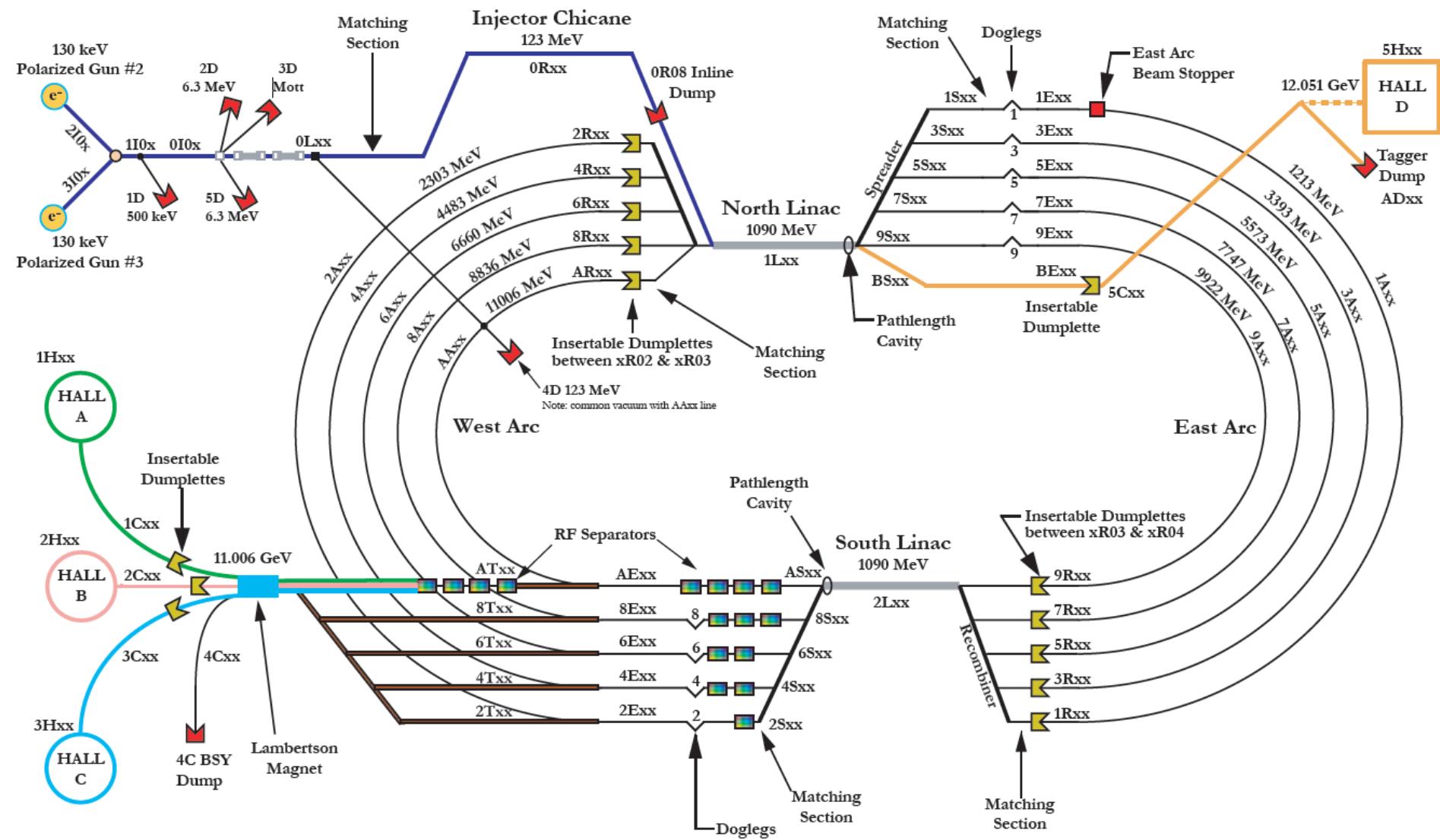


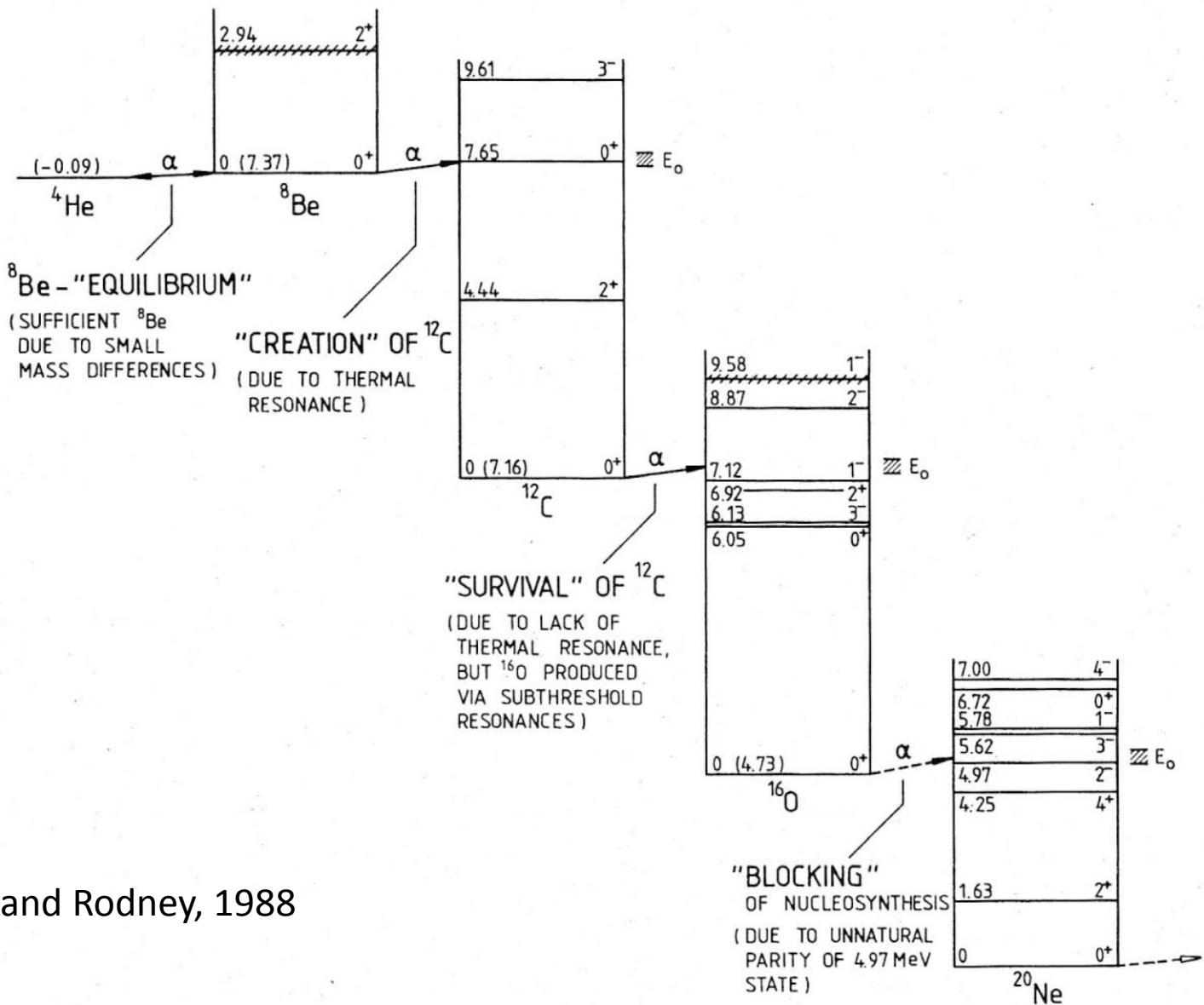
- $^{12}\text{C}(\gamma, 2\alpha)\alpha$  Background

# WATER SUPERHEATED LIQUID?

- Etching of glass vessel by superheated H<sub>2</sub>O
- T = 250°C
- P = 75 atm
- Background from secondary neutron–nucleus elastic scattering by neutrons from d( $\gamma$ ,n)p







Rolfs and Rodney, 1988