Josh Yoskowitz Ionization Cross Sections for "After 2 Days" RGA Spectrum

1 Purpose

To calculate electron-impact ionization cross sections for gas species found in the "After 2 Days" residual gas analyzer (RGA) spectrum taken on 5/21/18. The spectrum was analyzed using gnuplot and is shown below. Each substantial peak was identified and fit with a Gaussian function in order to determine the partial pressures of the various species of residual gas in the gun chamber.



Figure 1: Analysis of the RGA spectrum for the "After 2 Days" data

2 Calculation of the Ionization Cross Section

The equation for the calculation of the ionization cross section σ_i of the i^{th} gas species can be found in Reiser [1] and was originally developed by Slinker et. al. [4]:

$$\sigma_i = \frac{8a_0^2 \pi I_R A_1}{m_e c^2 \beta^2} f\left(\beta\right) \left(\ln \frac{2A_2 m_e c^2 \beta^2 \gamma^2}{I_R} - \beta^2 \right) \tag{1}$$

Numerically, this can be rewritten as:

$$\sigma_{i[m^{2}]} = \frac{1.872 \times 10^{-24} A_{1}}{\beta^{2}} f\left(\beta\right) \left[\ln\left(7.515 \times 10^{4} A_{2} \beta^{2} \gamma^{2}\right) - \beta^{2}\right]$$
(2)

In these two equations, $a_0 = 5.29 \times 10^{-11}$ m is the Bohr radius, $I_R = 13.6$ eV is the Rydberg energy, $m_e c^2$ is the rest mass energy of the electron, and β and γ are relativistic factors, A_1 and A_2 are empirical constants that depend on the type of gas species, and $f(\beta)$ is a function used when fitting data at low energies, i.e. $T_e \approx I_i$ where T_e is the kinetic energy of the electron and I_i is the ionization energy for the i^{th} gas species. Expressions for A_1 , A_2 , and $f(\beta)$ are given below:

$$f\left(\beta\right) = \frac{I_i}{T_e} \left(\frac{T_e}{I_i} - 1\right) = \frac{2I_i}{m_e c^2 \beta^2} \left(\frac{m_e c^2 \beta^2}{2I_i} - 1\right) \tag{3}$$

$$A_1 = M^2 \tag{4}$$

$$A_2 = \frac{e^{M^2}}{7.515 \times 10^4} \tag{5}$$

where C and M^2 are parameters given by Rieke and Prepejchal [2]. As an example calculation, for H₂, the values of C, $M^2 = A_1$, A_2 , and the ionization energy I_i from [2, 3] are given in the table below:

Gas Species	$A_1 = M^2$	C	A_2	$I_i(eV)$
H_2	0.695	8.115	1.5668	15.4

Table 1: Values for C, $M^2 = A_1$, and A_2 given by Rieke and Prepejchal and I_i given by Salancon for H₂ gas [2, 3]

Since at high energies, $\beta_e \gg \beta_{ion}$, we will assume that in the above equations, $\beta \approx \beta_e$. For a 200keV electron beam, $T_e = 200$ keV, $m_e c^2 = 511$ keV,

$$\begin{split} T_e &= (\gamma_e - 1) \, m_e c^2 = 200 \text{keV} \\ m_e c^2 &= 511 \text{keV} \\ \gamma_e &= 1 + \frac{T_e}{m_e c^2} = 1.39 \\ \beta_e &= \sqrt{1 - \frac{1}{\gamma_e^2}} = 0.695 \left(= 2.08 \times 10^8 \frac{\text{m}}{\text{s}} \right) \\ f\left(\beta_e\right) &= \frac{I_i}{T_e} \left(\frac{T_e}{I_i} - 1\right) \approx 1 \\ \sigma_i &= \frac{1.872 \times 10^{-24} A_1}{\beta_e^2} f\left(\beta_e\right) \left[\ln\left(7.515 \times 10^4 A_2 \beta_e^2 \gamma_e^2\right) - \beta_e^2 \right] \\ &\approx 2.994 \times 10^{-23} \text{m}^2 \end{split}$$

The change in density of the electron and gas molecules over time is given by Reiser[1]

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$$\frac{dn}{dt} = n_b n_g \sigma_i v = n_b n_g \sigma_i \beta_e c \tag{6}$$

At standard temperature ($T_0 = 273.15$ K) and pressure ($p_0 = 760$ torr = 1atm) the density of an ideal gas in a given volume is given by Loschmidt's number:

$$n_0 = \frac{p_0}{k_B T_0} \approx 2.687 \times 10^{25} \mathrm{m}^{-3} \tag{7}$$

Thus, for a given gas, its density is

$$n_g \left[m^{-3} \right] = \left(3.54 \times 10^{22} \right) p (\text{torr})$$

For H₂, p (torr) can be read off from Figure 1. In this case, $p_{H_2} = 1.11 \times 10^{-9}$ torr, so $n_{H_2} = 3.93 \times 10^{13} \text{m}^{-3}$. Knowing the electron density in the beam, n_b , one can calculate the ionization rate $\frac{dn}{dt}$.

References

- [1] Martin Reiser. Theory and Design of Charged Particle Beams. Wiley VCH Verlag GmbH, 2008.
- [2] Foster F. Rieke and William Prepejchal. Ionization cross sections of gaseous atoms and molecules for high-energy electrons and positrons. *Physical Review A*, 6(4):1507–1519, oct 1972.
- [3] E. Salançon, Z. Hammadi, and R. Morin. A new approach to gas field ion sources. Ultramicroscopy, 95:183-188, may 2003.
- [4] S. P. Slinker, R. D. Taylor, and A. W. Ali. Electron energy deposition in atomic oxygen. Journal of Applied Physics, 63(1):1–10, jan 1988.