

## 1 Purpose

To calculate electron-impact ionization cross sections for gas species found in the “After 2 Days” residual gas analyzer (RGA) spectrum taken on 5/21/18. The spectrum was analyzed using gnuplot and is shown below in Figure 1. Each substantial peak was identified and fit with a Gaussian function in order to determine the partial pressures of the various species of residual gas in the gun chamber. NOTE: The peak values must be divided by the correction factors listed here: <https://www.mksinst.com/docs/ur/GaugeGasCorrection.aspx>

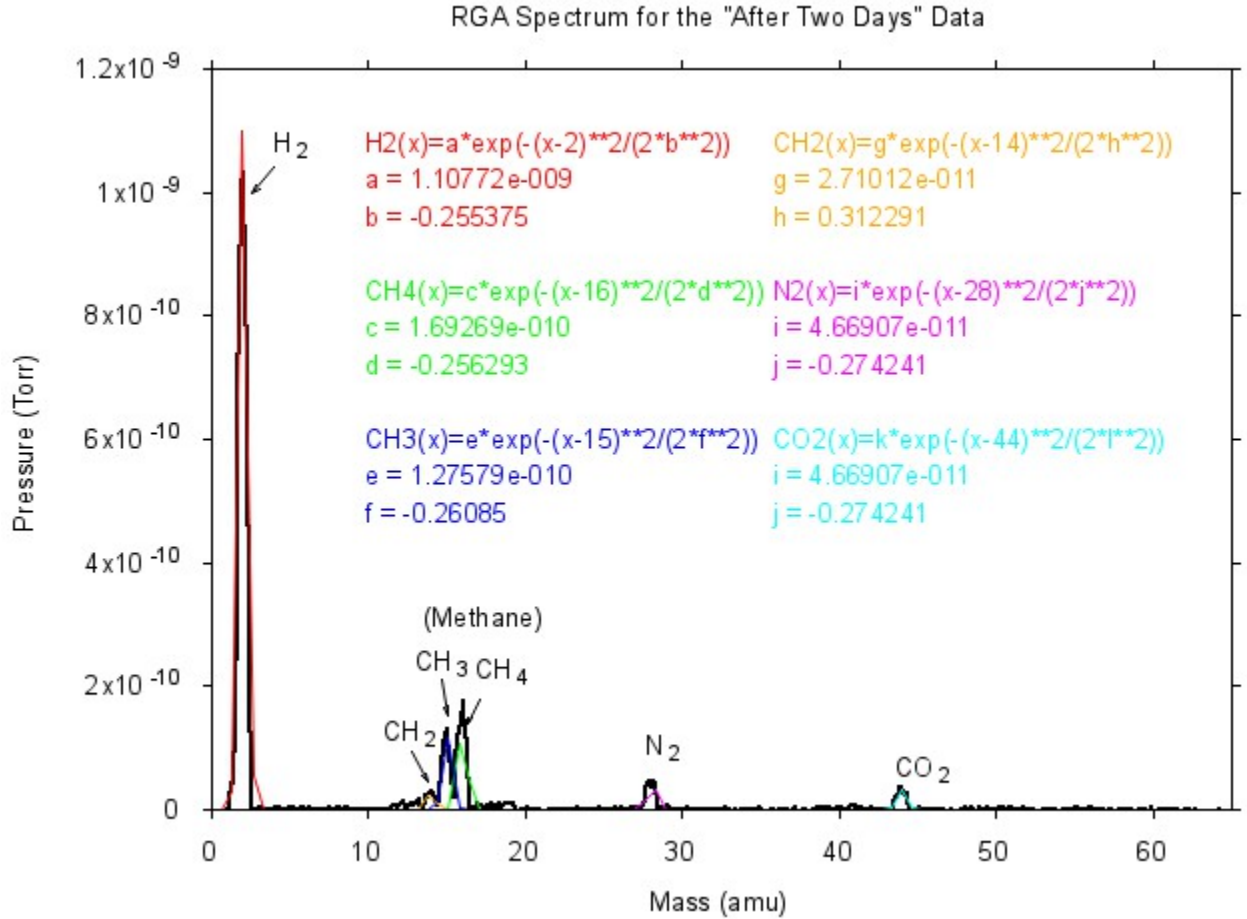


Figure 1: Analysis of the RGA spectrum for the “After 2 Days” data (before correction factor)

## 2 Calculation of the Ionization Cross Section

The equation for the calculation of the ionization cross section  $\sigma_i$  of the  $i^{th}$  gas species can be found in Reiser [1] and was originally developed by Slinker et. al. [3]:

$$\sigma_i = \frac{8a_0^2 \pi I_R A_1}{m_e c^2 \beta^2} f(\beta) \left( \ln \frac{2A_2 m_e c^2 \beta^2 \gamma^2}{I_R} - \beta^2 \right) \quad (1)$$

Numerically, this can be rewritten as:

$$\sigma_{i[\text{m}^2]} = \frac{1.872 \times 10^{-24} A_1}{\beta^2} f(\beta) [\ln(7.515 \times 10^4 A_2 \beta^2 \gamma^2) - \beta^2] \quad (2)$$

In these two equations,  $a_0 = 5.29 \times 10^{-11} \text{m}$  is the Bohr radius,  $I_R = 13.6 \text{eV}$  is the Rydberg energy,  $m_e c^2$  is the rest mass energy of the electron, and  $\beta$  and  $\gamma$  are relativistic factors,  $A_1$  and  $A_2$  are empirical constants that depend on the type of gas species, and  $f(\beta)$  is a function used when fitting data at low energies, i.e.  $T_e \approx I_i$  where  $T_e$  is the kinetic energy of the electron and  $I_i$  is the ionization energy for the  $i^{\text{th}}$  gas species. Expressions for  $A_1$ ,  $A_2$ , and  $f(\beta)$  are given below:

$$f(\beta) = \frac{I_i}{T_e} \left( \frac{T_e}{I_i} - 1 \right) = \frac{2I_i}{m_e c^2 \beta^2} \left( \frac{m_e c^2 \beta^2}{2I_i} - 1 \right) \quad (3)$$

$$A_1 = M^2 \quad (4)$$

$$A_2 = \frac{e^{\frac{C}{M^2}}}{7.515 \times 10^4} \quad (5)$$

where  $C$  and  $M^2$  are parameters given by Rieke and Prepejchal [2]. For  $\text{H}_2$ ,  $\text{CH}_4$ ,  $\text{CH}_3$ ,  $\text{N}_2$ , and  $\text{CO}_2$  the values of  $C$ ,  $M^2 = A_1$ ,  $A_2$ , and the ionization energy  $I_i$  from NIST (<https://webbook.nist.gov/>) are given in the table below:

Gas Species	$A_1 = M^2$	$C$	$A_2$	$I_i(\text{eV})$
$\text{H}_2$	0.695	8.115	1.5668	15.4
$\text{CH}_4$	4.23	41.85	0.2635	12.6
$\text{N}_2$	3.74	34.84	0.1478	15.6
$\text{CO}_2$	5.75	55.92	0.2227	13.8

Table 1: Values for  $C$ ,  $M^2 = A_1$ , and  $A_2$  given by Rieke and Prepejchal and  $I_i$  given by NIST for gas species found in the RGA spectrum.

Since at high energies,  $\beta_e \gg \beta_{ion}$ , we will assume that in the above equations,  $\beta \approx \beta_e$ . As an example calculation, for a 200keV electron beam,  $T_e = 200 \text{keV}$ ,  $m_e c^2 = 511 \text{keV}$ ,

$$\begin{aligned} T_e &= (\gamma_e - 1) m_e c^2 = 200 \text{keV} \\ m_e c^2 &= 511 \text{keV} \\ \gamma_e &= 1 + \frac{T_e}{m_e c^2} = 1.39 \\ \beta_e &= \sqrt{1 - \frac{1}{\gamma_e^2}} = 0.695 \left( = 2.08 \times 10^8 \frac{\text{m}}{\text{s}} \right) \\ f(\beta_e) &= \frac{I_i}{T_e} \left( \frac{T_e}{I_i} - 1 \right) \approx 1 \\ \sigma_i &= \frac{1.872 \times 10^{-24} A_1}{\beta_e^2} f(\beta_e) [\ln(7.515 \times 10^4 A_2 \beta_e^2 \gamma_e^2) - \beta_e^2] \\ &\approx 2.994 \times 10^{-23} \text{m}^2 \end{aligned}$$

The change in density of the electron and gas molecules over time is given by Reiser[1]

$$\frac{dn}{dt} = n_b n_g \sigma_i v = n_b n_g \sigma_i \beta_e c \quad (6)$$

At standard temperature ( $T_0 = 273.15 \text{K}$ ) and pressure ( $p_0 = 760 \text{torr} = 1 \text{atm}$ ) the density of an ideal gas in a given volume is given by Loschmidt's number:

$$n_0 = \frac{p_0}{k_B T_0} \approx 2.687 \times 10^{25} \text{m}^{-3} \quad (7)$$

Thus, for a given gas, its density is

$$n_g [\text{m}^{-3}] = (3.54 \times 10^{22}) p (\text{torr})$$

For  $\text{H}_2$ ,  $p (\text{torr})$  can be read off from Figure 1. In this case,  $p_{\text{H}_2} = 1.11 \times 10^{-9} \text{torr}$  with a correction factor of 0.46 for  $\text{H}_2$ , so  $n_{\text{H}_2} = 8.54 \times 10^{13} \text{m}^{-3}$ . Knowing the electron density in the beam,  $n_b$ , one can calculate the ionization rate  $\frac{dn}{dt}$ .

### 3 Ionization Cross Section vs. $T_e$

Starting from equation (2),

$$\sigma_i [\text{m}^2] = \frac{1.872 \times 10^{-24} A_1}{\beta^2} \frac{I_i}{T_e} \left( \frac{T_e}{I_i} - 1 \right) [\ln (7.515 \times 10^4 A_2 \beta^2 \gamma^2) - \beta^2]$$

we can rewrite  $\beta$  in terms of the electron beam kinetic energy  $T_e$ , which is proportional to the beam voltage:

$$\begin{aligned} T_e &= (\gamma - 1) m_e c^2 \\ \gamma &= 1 + \frac{T_e}{m_e c^2} \\ \frac{1}{\sqrt{1 - \beta^2}} &= 1 + \frac{T_e}{m_e c^2} \\ 1 - \beta^2 &= \left( \frac{1}{1 + \frac{T_e}{m_e c^2}} \right)^2 = \left( \frac{m_e c^2}{m_e c^2 + T_e} \right)^2 \\ \beta^2 &= 1 - \left( \frac{m_e c^2}{m_e c^2 + T_e} \right)^2 \end{aligned}$$

Thus,

$$\sigma_i = \frac{1.872 \times 10^{-24} A_1}{1 - \left( \frac{m_e c^2}{m_e c^2 + T_e} \right)^2} \frac{I_i}{T_e} \left( \frac{T_e}{I_i} - 1 \right) \left[ \ln \left( 7.515 \times 10^4 A_2 \left( 1 - \left( \frac{m_e c^2}{m_e c^2 + T_e} \right)^2 \right) \left( 1 + \frac{T_e}{m_e c^2} \right) \right) - \left( 1 - \left( \frac{m_e c^2}{m_e c^2 + T_e} \right)^2 \right) \right]$$

Using values in Table 1, a plot of  $\sigma_i$  vs.  $T_e$  for each of the gas species using Mathematica:

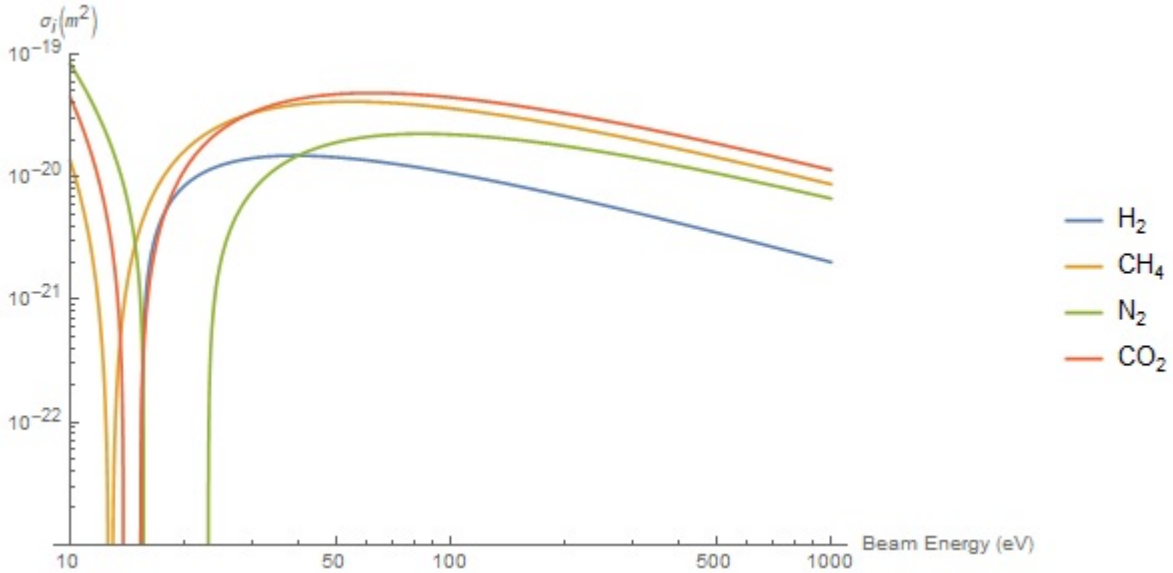


Figure 2: Plot of the ionization cross section  $\sigma_i$  vs. electron kinetic energy  $T_e$

## References

- [1] Martin Reiser. *Theory and Design of Charged Particle Beams*. Wiley VCH Verlag GmbH, 2008.
- [2] Foster F. Rieke and William Prepejchal. Ionization cross sections of gaseous atoms and molecules for high-energy electrons and positrons. *Physical Review A*, 6(4):1507–1519, oct 1972.
- [3] S. P. Slinker, R. D. Taylor, and A. W. Ali. Electron energy deposition in atomic oxygen. *Journal of Applied Physics*, 63(1):1–10, jan 1988.