

# Calibration of a Mott Polarimeter

Mexican Physical Society PROGRAM





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### Introduction to Mott scattering



When energetic electrons are scattered from massive nucleus, the scattering crosssection gets modified due to the spin-orbit interaction as follows:



We can measure the number of electrons scattered to the right or left

$$(1 + PS) \rightarrow \epsilon = \frac{N^r - N^l}{N^r + N^l}$$
$$\mathbf{Asymmetry}$$
$$\epsilon = PS$$

Considering flipping the polarization (**helicity**)

$$\epsilon = \frac{1-r}{1+r} \quad \mbox{with} \quad r = \sqrt{\frac{(N_+^l)(N_-^r)}{(N_-^l)(N_+^r)}} \label{eq:eq:elements}$$

The target has lot of atoms instead of a single massive nucleus, this lead us to possible multiple scatterings, when this happens the value of the Sherman functions gets reduced

The number of electrons that hit the target in a certain time, this is called the rate:

$$R(\theta) = N/T = \sigma(\theta)L\rho(I/e^{-})$$
thickness of beam current the target target density





 $\epsilon$  gets reduced for targets with more atoms (thicker)

### Objective

Calculate the polarization through the Sherman functions for this specific energy beam and through the asymmetries obtained from the experiment measuring the number of scattered electrons after having them hit gold foil targets of different thicknesses.

$$\epsilon = PS$$

## The experiment

#### **Electron gun**



Provides a polarized electron beam by illuminating a GaAs cathode with a circularly polarized laser and using high voltage to accelerate electrons to 180 keV



Rotates the direction of the spin by using perpendicular electric and magnetic fields. This effect is governed by the BMT equation:

#### **Wien Filter**

$$\frac{e}{m_0 c} \left[ aB_x + \left( \frac{(1+a)\beta^2 - 1}{\beta} \right) E_y \right]$$

### Mott polarimeter











## Data acquisition

The detectors measure 4 main things:

- **1** Value of the energy changing over time for the left and right detectors.
- 2 Logical sign: when to stop taking data.
- Helicity: takes high and low values for 33 ms
- **4** T-Settle: shows the helicity transition (0.5 ms)
- We will sum over the values of the energy to obtain **histograms** that organize the data.

energy

VS

no. of electrons



thickness.



Number of electrons with positive (red plot) and negative helicity (black plot) with a certain energy for the target of 60 nm thickness

### Measuring asymmetry and polarization

With the selected electrons for each one of the target foils a **target thickness** scan is obtained.

A fit of the following form is applied:

 $f(x) = \frac{a}{1+bx}$ 

a = 15.5305

b = 0.00196351

 $S_0 = -0.4261$ 



### Asymmetry and voltage applied to the Wien filter

The figure shows how the asymmetry factor gets modified as the voltage applied to the Wien filter increase in a target foil of 40 nm

A sinusoidal fit was performed as the yellow dashed line shows:

$$f(x) = a \sin (bx + c)$$
  

$$a = 14.2421$$
  

$$b = 0.0631088$$
  

$$c = -0.153606$$



## Asymmetry and current applied to the solenoids



A sinusoidal fit was performed as the yellow dashed line shows:

The beamline contains two solenoids with fields A and B for focusing and directing the electron beam.

We can observe how the asymmetry changes as the factor (A + B) changes for the target foil of 40 nm.

 $f(x) = a \sin(bx + c) - \begin{bmatrix} a = 14.5858 \\ b = 0.00367245 \\ c = -0.719592 \end{bmatrix}$ 

### Conclusions

It is important to select the data, to know which electrons will be used in the final measurement. The detectors will find more scattered electrons of one type than another depending on their helicity

The polarization has a value of  $\sim$  36%, which is congruent with the values normally found in the

The asymmetry changes when the solenoids current and the voltage of the Wien filter is modified.

interature.

Thicker foils have smaller values of asymmetry, the electrons get scattered multiple times.