Presentation on Ion Production and Clearing

(aka What have I been up to all this time???)

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Motivation and Questions to Answer

- Overarching question/goal: How can we increase the QE and thus charge lifetime of the photocathode?
- Current question: How can we reduce the number of ions reaching the photocathode and potentially damaging it, resulting in lower QE?
- Need to know:
  1. Ion production rates between gun and VWF
  2. Where newly formed ions go (are they trapped or do they leave the beam?)
  3. Distribution of ions at photocathode
- Possible solutions
  - Ion Clearing Gap (Bunch gap)
  - Clearing Electrode (Wien filter?)
  - Beam Shaking (Driving the beam at the oscillating ions' resonance frequency)

Ion production rate

- The ion production rate for a electron beam ionizing a certain gas is given by Reiser[3]

$$\frac{dn}{dt} = n_b n_g \sigma_i v$$

where n is the ion density, n_b is the electron beam density, n_g is the neutral gas density, σ_i is the ionization cross section for a given gas species, and v is the velocity of the electrons.

Ionization Cross Section

- The form for the ionization cross section σ_i for gas species i follows from Bethe's theory[1]. The general form used by Reiser is in the form from Slinker, Tayler and Ali's paper[5] shown below:

$$\sigma_i = \frac{8a_0^2\pi I_R A_1}{m_e c^2\beta^2} f(\beta)\left(\ln\frac{2A_2 m_e c^2\beta^2\gamma^2}{I_R} - \beta^2\right)$$

$$= \frac{1.872\times 10^{-24} A_1}{\beta^2} f(\beta)\left[\ln\left(7.515\times 10^4 A_2\beta^2\gamma^2\right) - \beta^2\right]$$

$$f(\beta) = \frac{I_i}{T_e}\left(\frac{T_e}{I_i} - 1\right) = \frac{2I_i}{m_e c^2\beta^2}\left(\frac{m_e c^2\beta^2}{2I_i} - 1\right)$$

Here, a_0 is the Bohr radius, $I_R = 13.6$ eV is the Rydberg energy, $\beta = \frac{v}{c}$, $\gamma = \left(1-\beta^2\right)^{-\frac{1}{2}}$, m_e and T_e are the electron's mass and kinetic energy respectively, I_i is the ionization energy of gas species i, $f(\beta)$ is a correction function for fitting the velocity data at low energies ($T_e \approx I_i$), and $A_1 = M^2$ & $A_2 = \frac{e^{C/M^2}}{7.515\times 10^4}$ are emperical constants that depend on the gas species. These constants are given by Rieke and Prepejchal[4].

Ionization Cross Section vs Beam Energy

- Rewriting σ_i as a function of beam energy T_e, we can plot σ_i for various gas species that are common in the accelerator vacuum:

$$\sigma_i = \frac{1.872 \times 10^{-24} A_1}{1 - \left(\frac{m_e c^2}{m_e c^2 + T_e}\right)^2} \frac{I_i}{T_e} \left(\frac{T_e}{I_i} - 1\right)$$

$$\times \left[\ln \left(7.515 \times 10^4 A_2 \left(1 - \left(\frac{m_e c^2}{m_e c^2 + T_e}\right)^2\right) \left(1 + \frac{T_e}{m_e c^2}\right)\right) - \left(1 - \left(\frac{m_e c^2}{m_e c^2 + T_e}\right)^2\right)\right]$$

Gas Species	$A_1 = M^2$	C	A_2	I_i(eV)
H_2	0.695	8.115	1.5668	15.4
CH_4	4.23	41.85	0.2635	12.6
N_2	3.74	34.84	0.1478	15.6
CO_2	5.75	55.92	0.2227	13.8

Table: Values for C, $M^2 = A_1$, and A_2 given by Rieke and Prepejchal and I_i given by NIST for some common gas species.

Ionization Cross Section vs Beam Energy cont'd

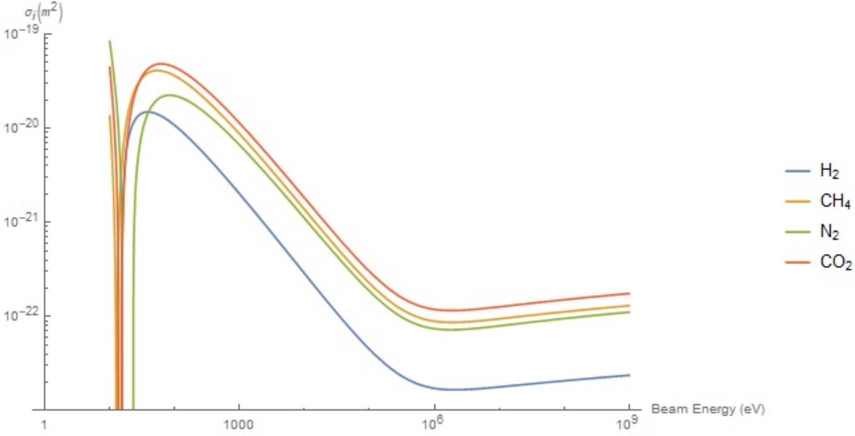

Figure: Plot of the ionization cross section σ_i vs. beam energy T_e

Example RGA Spectrum for Calculating Ion Production Rates

- For given values of n_g, n_b and v, we can calculate σ_i and the ion production rate $\frac{dn}{dt}$. As an example, we can use the RGA spectrum below to get the densities of the gas species n_g in the accelerator vacuum:

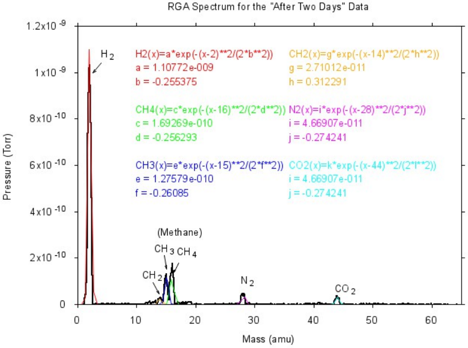

Figure: Analysis of the RGA spectrum for the "After 2 Days" data (before correction factor)

Gas Densities

We can assume that the residual gas behaves ideally (obeys Newton's laws, volume of gas molecules is much smaller than the gas volume, no external forces on the molecules, molecules in random motion). At standard temperature ($T_0 = 273.15\text{K}$) and pressure ($p_0 = 760\text{torr} = 1\text{atm}$) the density of an ideal gas in a given volume is given by Loschmidt's number:

$$n_0 = \frac{p_0}{k_B T_0} \approx 2.687 \times 10^{25} \text{m}^{-3}$$

Thus, for a given gas, its density is

$$n_g \left[\text{m}^{-3}\right] = \left(3.54 \times 10^{22}\right) p\,(\text{torr})$$

The partial pressures are calculated from the Gaussian fit functions (given in the RGA spectrum) using the Gaussian integral:

$$\int_{-\infty}^{\infty} e^{-a(x+b)^2} dx = \sqrt{\frac{\pi}{a}}$$

$$\int_{-\infty}^{\infty} Ae^{-\frac{(x+b)^2}{2\sigma^2}} dx = A\sqrt{2\pi\sigma^2}$$

Gas Densities cont'd

- These partial pressures then need to be corrected using correction factors that adjust the pressures of the gas species relative to nitrogen N_2 (from MKS website).
- The correction factor for each parent ion is assumed to be the same for each ion in its class (as in the case of CH_4, CH_3, and CH_2). Assuming an extractor gauge pressure of 2×10^{-12}torr, we can normalize these partial pressures by a normalization factor α that is equal to the sum of the corrected partial pressures divided by the extractor gauge pressure. Each partial pressure is then multiplied by α so that the sum of the partial pressures is 2×10^{-12}torr. In this case, $\alpha \approx 2.87 \times 10^{-3}$.
- From the normalized partial pressures, the number densities and ion production rates can be calculated.

Table of Values for n_g, σ_i and $\frac{dn}{dt}$

Gas species	Uncorrected Pressure (torr)	Correction factor	Corrected Pressure (torr)	Normalized Pressure (torr)
H_2	7.09085×10^{-10}	0.46	3.26×10^{-10}	9.28×10^{-13}
CH_4	1.08744×10^{-10}	1.40	1.52×10^{-10}	4.33×10^{-13}
CH_3	8.34180×10^{-11}	1.40	1.17×10^{-10}	3.33×10^{-13}
CH_2	2.12148×10^{-11}	1.40	2.97×10^{-11}	8.45×10^{-14}
N_2	3.20961×10^{-11}	1.00	3.21×10^{-11}	9.14×10^{-14}
CO_2	3.20961×10^{-11}	1.42	4.56×10^{-11}	1.30×10^{-13}

Gas species	Gas Density n_g $\left(\text{molecules/m}^3\right)$	Ionization Cross Section σ_i $\left(\text{m}^2\right)$	Ion Production Rate $\left(\text{ions/m}^3\text{s}\right)$
H_2	3.29×10^{10}	2.99×10^{-23}	4.06×10^{17}
CH_4	1.53×10^{10}	1.53×10^{-22}	9.66×10^{17}
CH_3	1.18×10^{10}	8.00×10^{-23} *	3.89×10^{17}
CH_2	2.99×10^{9}	9.00×10^{-23}*	1.11×10^{17}
N_2	3.24×10^{9}	1.27×10^{-22}	1.70×10^{17}
CO_2	4.60×10^{9}	2.04×10^{-22}	3.87×10^{17}

*Denotes values from NIST using the Binary-Encounter-Bethe (BEB) model here https://physics.nist.gov/PhysRefData/Ionization/intro.html

Ion Production Rate vs Beam Energy

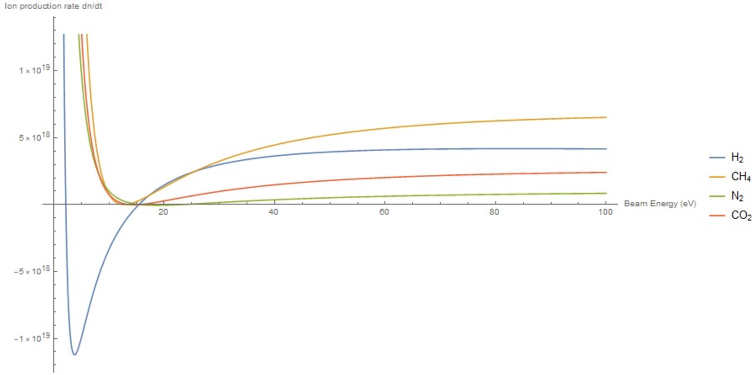

Figure: Ion production rate as a function of low beam energies for the gas species found in the RGA spectrum excluding CH_3 and CH_2.

Ion Production Rates at Various Beam Energies

Gas Species	IPR at 50eV $\left(m^{-3}\,s^{-1}\right)$	IPR at 100eV $\left(m^{-3}\,s^{-1}\right)$	IPR at 500eV $\left(m^{-3}\,s^{-1}\right)$	IPR at 1keV $\left(m^{-3}\,s^{-1}\right)$
H_2	3.93×10^{18}	4.15×10^{18}	3.03×10^{18}	2.45×10^{18}
CH_4	5.22×10^{18}	6.52×10^{18}	5.81×10^{18}	4.95×10^{18}
N_2	5.22×10^{17}	8.43×10^{17}	9.07×10^{17}	7.99×10^{17}
CO_2	1.80×10^{18}	2.41×10^{18}	2.26×10^{18}	1.94×10^{18}
Gas Species	IPR at 100keV $\left(m^{-3}\,s^{-1}\right)$	IPR at 130keV $\left(m^{-3}\,s^{-1}\right)$	IPR at 180keV $\left(m^{-3}\,s^{-1}\right)$	IPR at 1MeV $\left(m^{-3}\,s^{-1}\right)$
H_2	4.88×10^{17}	4.52×10^{17}	4.16×10^{17}	3.46×10^{17}
CH_4	1.15×10^{18}	‘1.07×10^{18}	9.91×10^{17}	8.46×10^{17}
N_2	2.00×10^{17}	1.87×10^{17}	1.73×10^{17}	1.50×10^{17}
CO_2	4.60×10^{17}	4.29×10^{17}	3.97×10^{17}	3.40×10^{17}

Table: Ion Production Rates (IPR) of each gas species for selected beam energies.

Where do newly formed ions go?

- According to Reiser, if the ions have the *same* charge as the beam particles (i.e. if the ions are negatively charged), then they are expelled to the walls by the beam's space charge in the absence of magnetic fields.
- If the ions have *opposite* charge to the beam particles, they are trapped in the beam and contribute to neutralization
- In order to consider the degree of ionization and neutralization, we have to consider several different cases and ask:
  - How does v_e compare with v_g?
  - How likely is ionization to occur? Qualitative ion production rate?
  - Are the collisions elastic or inelastic?
  - Under what conditions does recombination occur?

The Unlikely Case: $v_e \ll v_g$

- In this case, $T_g \gg T_e$ and the relative velocity u is large, thus all collisions are essentially *elastic*, provided the gas particles are relatively heavy (may not be true for hydrogen)
- For large ionization potentials I_i, the electron does not have enough energy to ionize the gas molecule and simply scatters off the molecule.
- In the unlikely case that the electron ionizes the gas molecule, the recombination cross section is given by Derbenev[2]:

$$\sigma_{rec} = \frac{16\pi z^2 e^6}{3\sqrt{3} m_e^2 v^2 c^3 \hbar} \ln \frac{2I}{m_e v^2} \text{ for } mv^2 \ll 2I$$

$$I = \frac{z^2 e^4 m}{2\hbar^2}$$

  We see that recombination is unlikely, since the relative velocity is too large
- Ion production rate is slim to none.
- This case is likely only important close to the photocathode. We can safely say that no ionization occurs there

The Possible Case: $v_e \approx v_g$

- In this case, $T_g \approx T_e$, so we must consider elastic and inelastic collisions:
- Elastic scattering occurs when:
  - Scattering angle is large (grazing collision)
  - Electron energy T_e is large
  - Ionization potential is large ($T_e \gg I$)
- Inelastic scattering occurs when
  - Scattering angle is small
  - Electron energy is large enough
  - Ionization potential is relatively small
- Uncertain when $T_e \approx I$. Better to restrict ourselves to above cases.
- Ionization will occur depending on how T_e and I compare.
- Recombination seems very likely (since u is small), thus ion production rate is irrelevant.

The Likely Case $v_e \gg v_i$

- In all cases, $T_e \gg T_g$, thus all collisions are elastic, provided the gas particles are relatively heavy (as with the first case)
- We can use Reiser's formula for σ_i, since $T_e \gg I$,
- Recombination is highly unlikely, since $m_e v^2 \gg I$
- Ion production rate is highest in this case.

Distribution of Ions at Photocathode

- Gas molecules lose electrons when ionized by electrons, thus all newly formed ions are positively charged (not sure how gas molecules can *gain* electrons from the electron beam), thus they are trapped in the electron beam.
- When electrons are accelerated forward, positively charged ions are accelerated backward. Since they are trapped by the beam, it stands to reason the ions would remain within the electron beam shape at the photocathode (beam shape determines distribution).
- We can use GPT to confirm these hypotheses. A more rigorous approach may be to use an approximate solution to the Boltzmann/Landau kinetic equation that would show exactly what the distribution of the ions are as a result of ionization (i.e. collisions).

References

[1] H. Bethe. Zur theorie des durchgangs schneller korpuskularstrahlen durch materie. *Annalen der Physik*, 397(3):325–400, 1930.

[2] Yaroslav Derbenev. Theory of electron cooling, 1978.

[3] Martin Reiser. *Theory and Design of Charged Particle Beams*. Wiley VCH Verlag GmbH, 2008.

[4] Foster F. Rieke and William Prepejchal. Ionization cross sections of gaseous atoms and molecules for high-energy electrons and positrons. *Physical Review A*, 6(4):1507–1519, oct 1972.

[5] S. P. Slinker, R. D. Taylor, and A. W. Ali. Electron energy deposition in atomic oxygen. *Journal of Applied Physics*, 63(1):1–10, jan 1988.