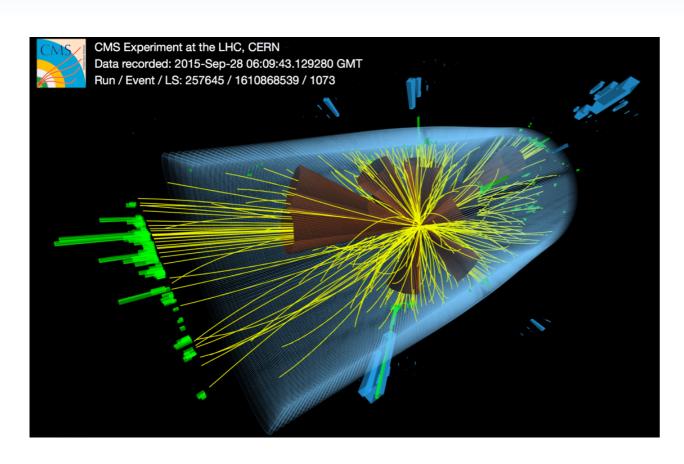
# Jet observable at the EIC

Kyle Lee Stony Brook University

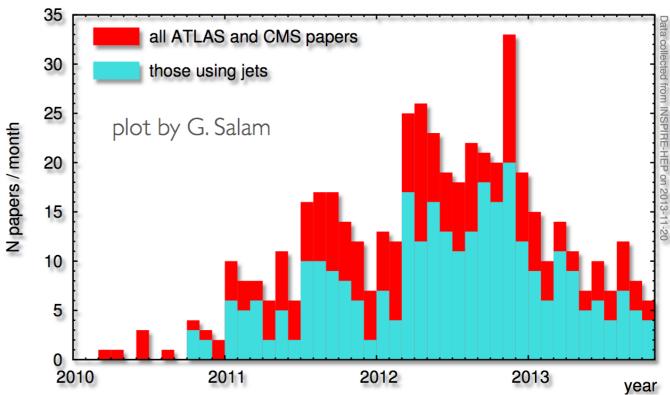
> EICUGM 07/30/18 - 08/02/18



# Jets at the LHC

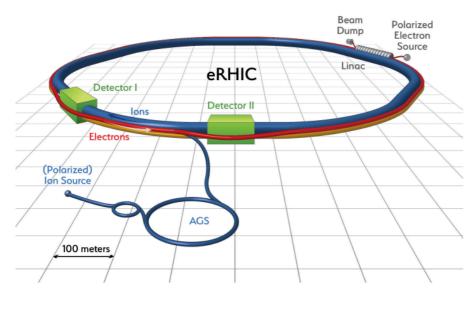


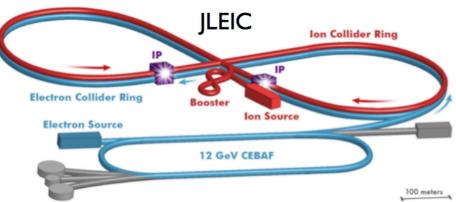
 Jets are produced copiously at the LHC



• At the LHC, 60 - 70 % of ATLAS & CMS papers use jets in their analysis!

# Jets at the EIC

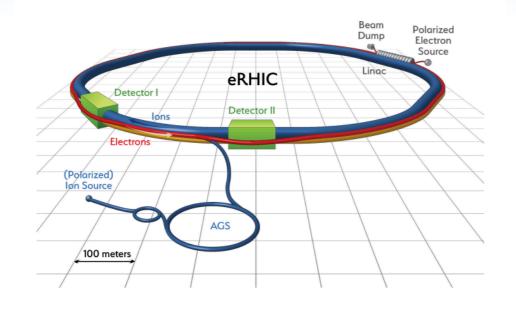




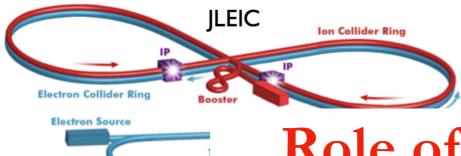
- $\sqrt{S_{\rm EIC}} \ll \sqrt{S_{\rm LHC}} \Leftrightarrow \sqrt{p_{T_J,\rm EIC}} \ll \sqrt{p_{T_J,\rm LHC}}$ Lower  $p_{T,J}$  for EIC
- $N_{J,EIC} \ll N_{J,LHC}$ Smaller jet multiplicity for EIC
- Less contamination from underlying events and pileups

• Different circumstances compared with the LHC and New opportunities

# Jets at the EIC



- $\sqrt{S_{\rm EIC}} \ll \sqrt{S_{\rm LHC}} \Leftrightarrow \sqrt{p_{T_J,\rm EIC}} \ll \sqrt{p_{T_J,\rm LHC}}$ Lower  $p_{T,J}$  for EIC
- $N_{J,EIC} \ll N_{J,LHC}$ Smaller jet multiplicity for EIC
- Less contamination from underlying events and pileups

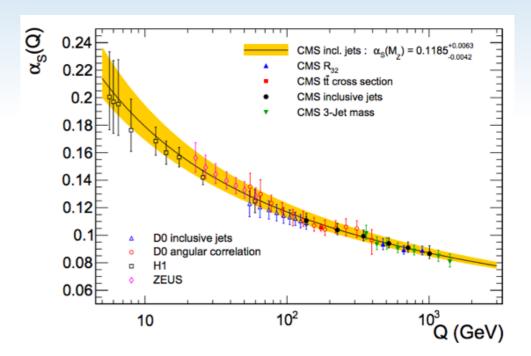


Role of higher power corrections?

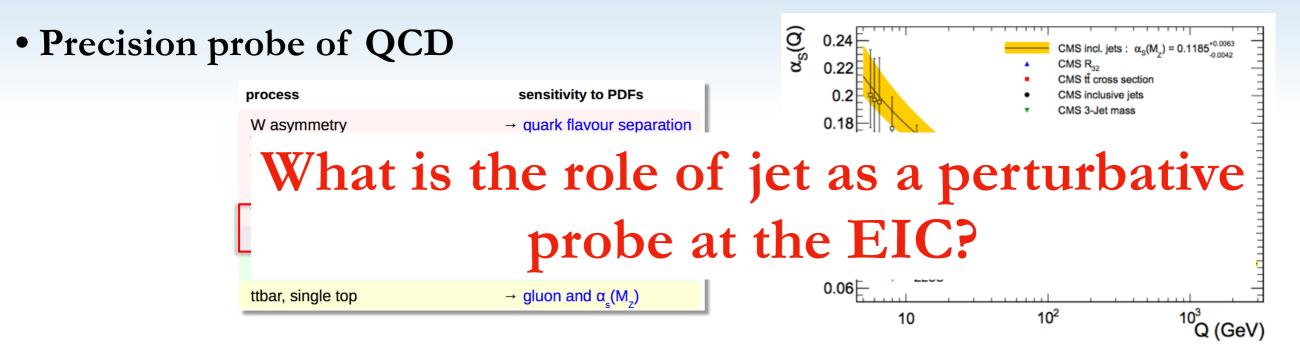
• Different circumstances compared with the LHC and New opportunities

• Precision probe of QCD

process	sensitivity to PDFs
W asymmetry W and Z production (differential) W+c production Drell-Yan (DY): high invariant mass Drell-Yan (DY): low invariant mass W,Z +jets	<ul> <li>→ quark flavour separation</li> <li>→ valence quarks</li> <li>→ strange quark</li> <li>→ sea quarks, high-x</li> <li>→ low-x</li> <li>→ gluon medium-x</li> </ul>
Inclusive jet and di-jet production	→ gluon and $\alpha_s(M_Z)$
Direct photon ttbar, single top	<ul> <li>→ gluon medium, high-x</li> <li>→ gluon and α<sub>s</sub>(M<sub>z</sub>)</li> </ul>



#### Inclusive jets - perturbative probe



Inclusive jets - perturbative probe

sensitivity to PDFs

 $\rightarrow$  gluon and  $\alpha_{1}(M_{2})$ 

Precision probe of QCD

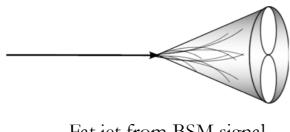
process

ttbar, single top

Wasymmetry old quark flavour separation old let as a perturbative probe at the EIC?

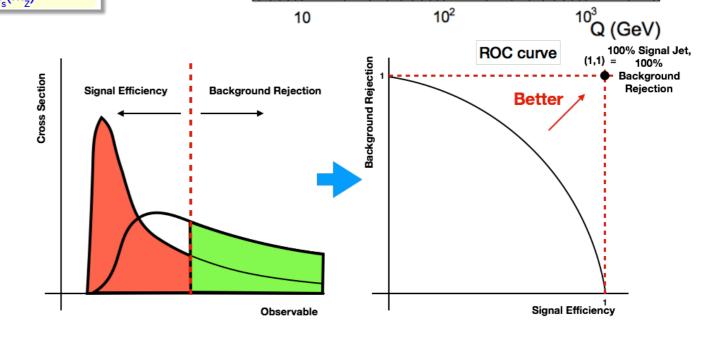
0.22

Constrain BSM Models



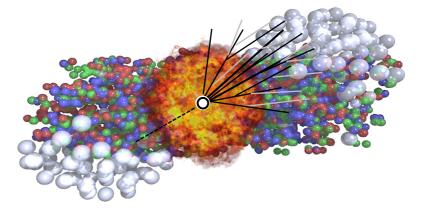
Fat jet from BSM signal

• Probe of quark gluon plasma

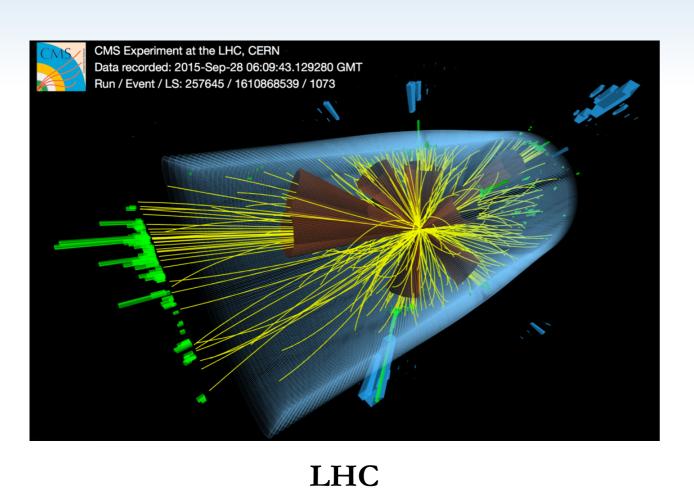


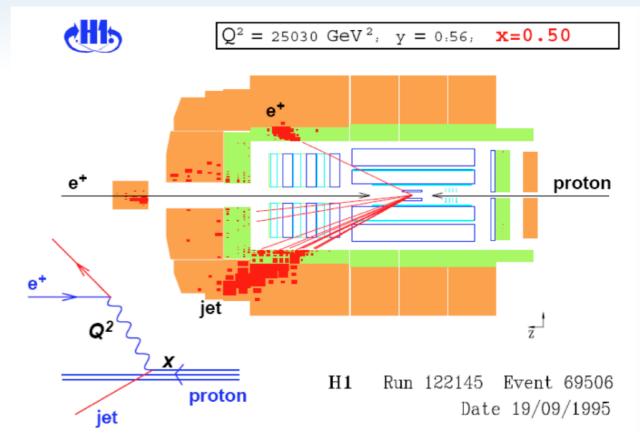
CMS tt cross section

CMS inclusive jets



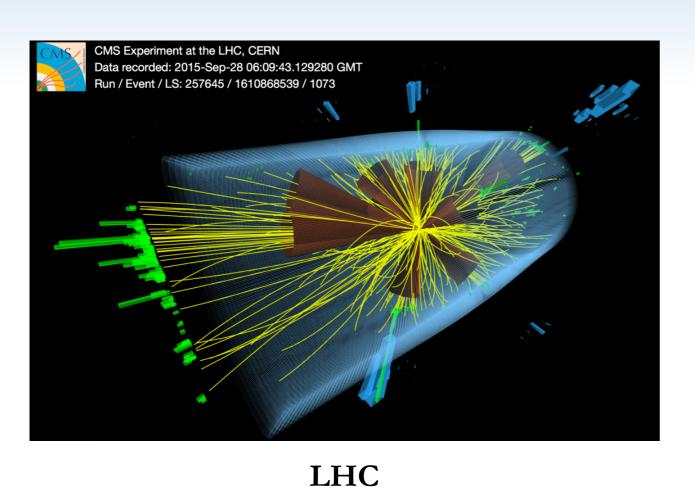
 Precision probe of QCD CMS tt cross section process sensitivity to PDFs CMS inclusive jets → quark flavour separation W asymmetry What is the role of jet as a perturbative probe at the EIC?  $\rightarrow$  gluon and  $\alpha_{1}(M_{1})$ ttbar, single top  $10^{2}$ 'Q (GeV) Constrain BSM Models 100% Signal Jet, **ROC** curve 100% Background Signal Efficiency Rejection **Background Rejection Better** Classification of different type of jets? Cold Nuclear Modification in e+A • Probe of quark gluon plasma

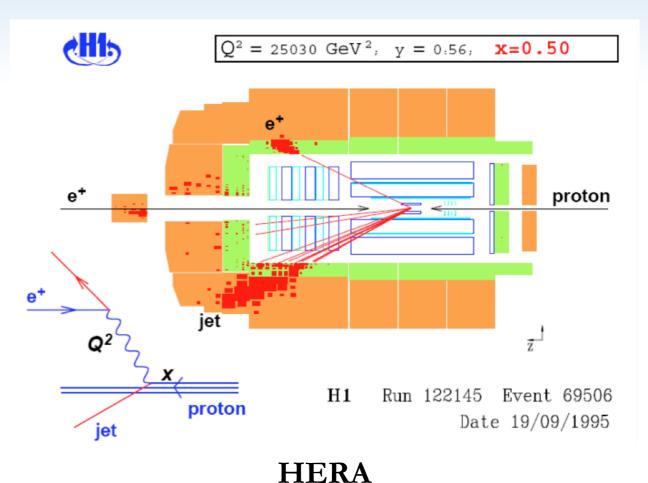




**HERA** 

• Typical event at the LHC and HERA

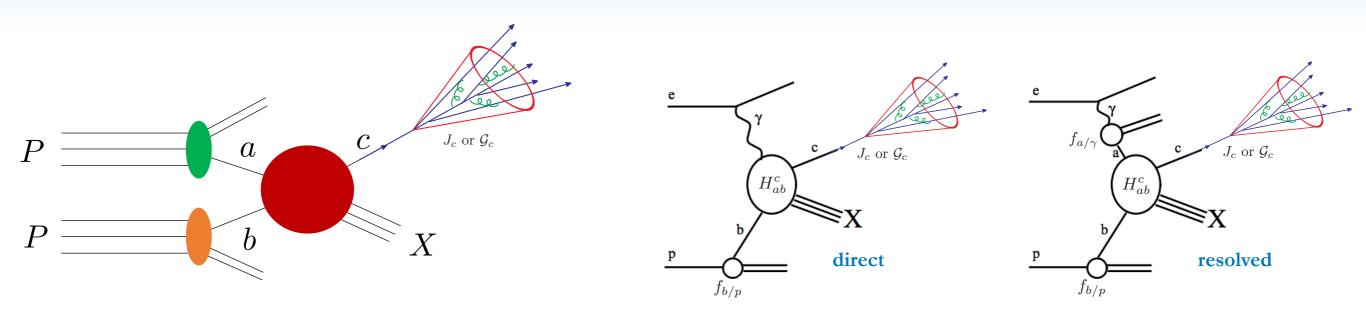




• Typical event at the LHC and HERA

# What is the role of NP physics at the EIC?

#### Processes of Interest



We want to study semi-inclusive jet production  $p + p \rightarrow Jet((with/without) substructure) + X$ 

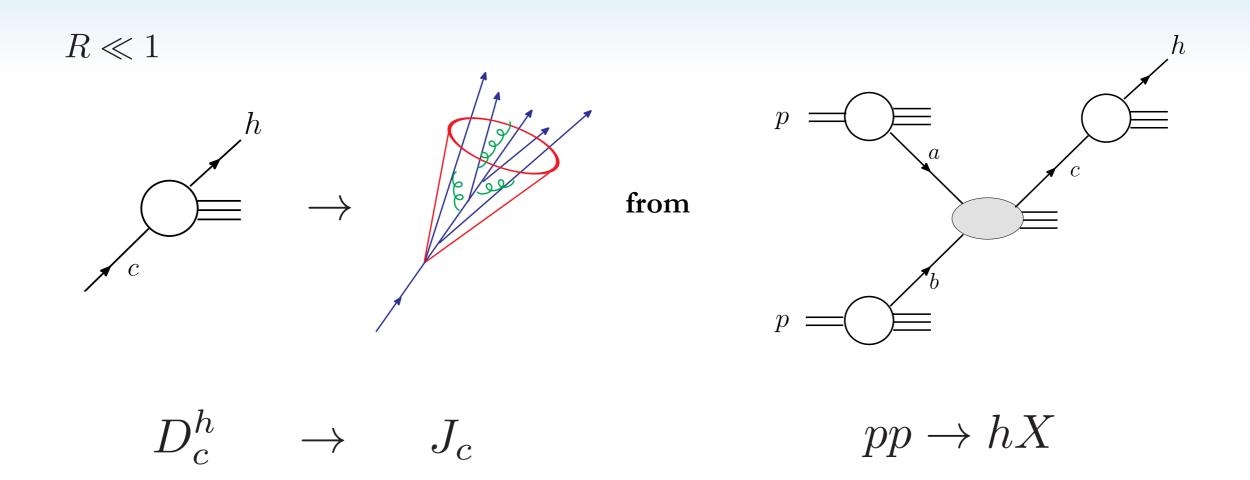
photoproduction at the EIC e + p  $\rightarrow$  e + Jet((with/without) substructure) + X

More statistics. No veto on additional jets.

#### Plans of this talk

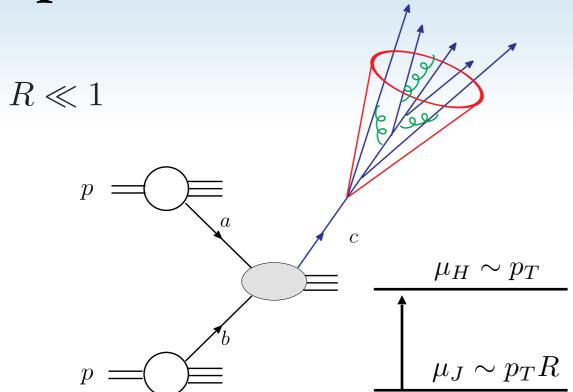
- Inclusive jet production at the LHC
- Jet substructure measurements at the LHC
- Role of non-perturbative effects
- Study of the EIC case
- Conclusions

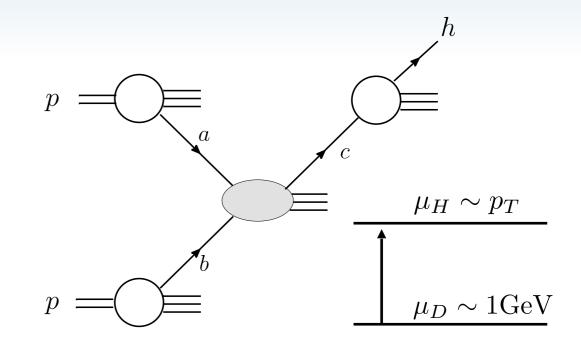
# Factorization of Inclusive Jet Production



• Simple replacement of the fragmentation function by "semi-inclusive jet function".

Comparison with the inclusive hadron production case





Factorization

Inclusive Jet 
$$\frac{d\sigma^{pp\to {\rm jet}X}}{dp_T d\eta} = \sum_{a,b,c} f_a \otimes f_b \otimes H^c_{ab} \otimes J_c + \mathcal{O}(R^2) + \mathcal{O}(\frac{\Lambda_{QCD}}{p_T R})$$

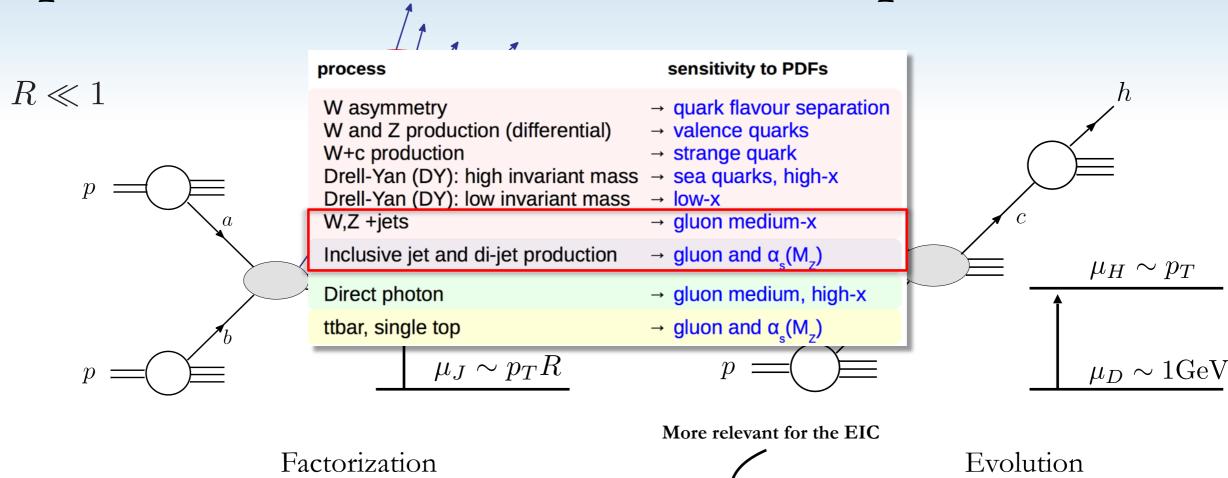
$$\frac{d\sigma^{pp\to hX}}{dp_T d\eta} = \sum_{a,b,c} f_a \otimes f_b \otimes H^c_{ab} \otimes D^h_c$$

Evolution

$$\mu \frac{d}{d\mu} J_i = \sum_j P_{ji} \otimes J_j$$

$$\mu \frac{d}{d\mu} D_i^h = \sum_j P_{ji} \otimes D_j^h$$

# Comparison with the inclusive hadron production case



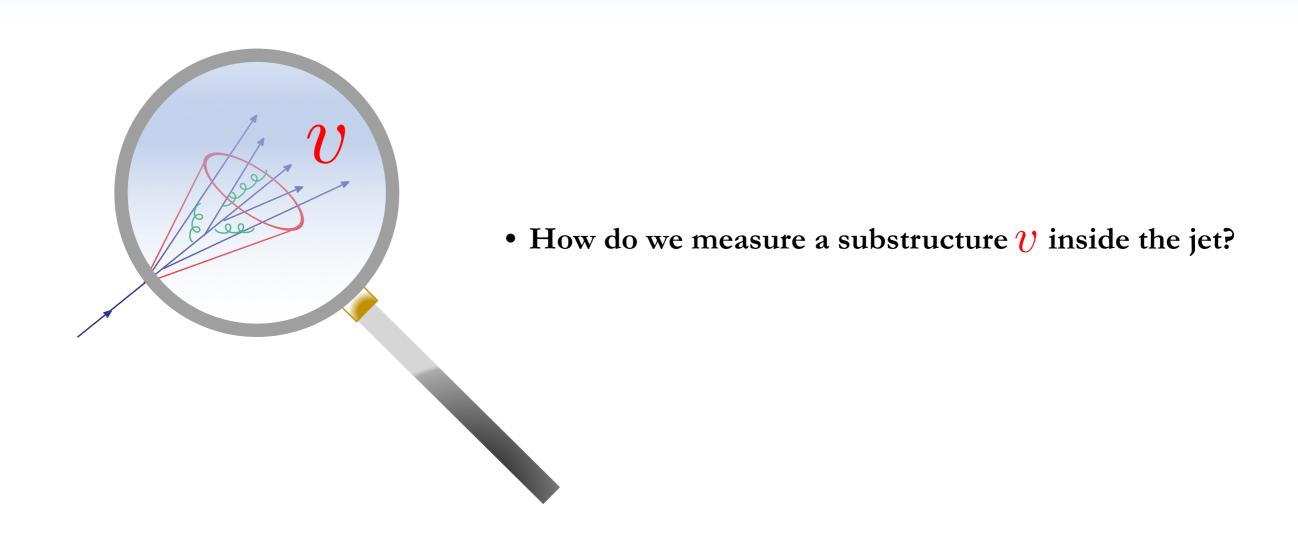
Factorization

Factorization

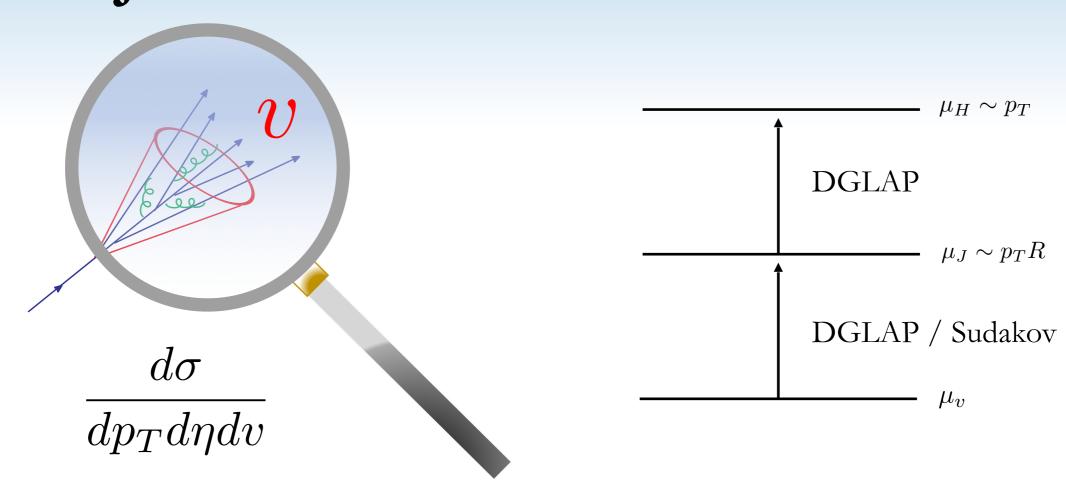
Evolution

Inclusive Jet  $\frac{d\sigma^{pp\to jet X}}{dp_T d\eta} = \sum_{a,b,c} f_a \otimes f_b \otimes H^c_{ab} \otimes J_c + \mathcal{O}(R^2) + \mathcal{O}(\frac{\Lambda_{QCD}}{p_T R})$   $\mu \frac{d}{d\mu} J_i = \sum_j P_{ji} \otimes J_j$ Hadron  $\frac{d\sigma^{pp\to hX}}{dp_T d\eta} = \sum_{a,b,c} f_a \otimes f_b \otimes H^c_{ab} \otimes D^h_c$   $\mu \frac{d}{d\mu} D^h_i = \sum_j P_{ji} \otimes D^h_j$ 

# Jet Substructure Measurements



#### Jet Substructure Measurements



- When we measure a substructure v from the jet, once we evolve to  $\mu_J$  the remaining evolution to  $\mu_H$  is given by DGLAP evolution!
- Two step factorization:
  - a) production of a jet
  - b) probing the internal structure of the jet produced.

# Jet angularity

• A generalized class of IR safe observables, angularity (applied to jet):

$$\tau_a^{e^+e^-} = \frac{1}{E_J} \sum_{i \in J} E_i \theta_{iJ}^{2-a} \qquad \text{More relevant for the EIC}$$
 
$$\tau_a^{pp} = \frac{1}{p_T} \sum_{i \in J} p_{T,i} (\Delta R_{iJ})^{2-a} = \left(\frac{2E_J}{p_T}\right)^{2-a} \tau_a^{e^+e^-} + \mathcal{O}((\tau_a^{pp})^2)$$
 
$$\tau_0^{pp} = \frac{m_J^2}{p_T^2} + \mathcal{O}((\tau_0^{pp})^2)$$

- More sensitive to collinear radiation as 'a' gets larger. (factorization breaks at a=2).
- a=0 related to thrust (jet mass)
- a=1 related to jet broadening (sensitive to rapidity divergence)
- Many studies done for exclusive case:

Sterman et al. `03, `08, Hornig, C. Lee, Ovanesyan `09, Ellis, Vermilion, Walsh, Hornig, C.Lee `10, Chien, Hornig, C. Lee `15, Hornig, Makris, Mehen `16

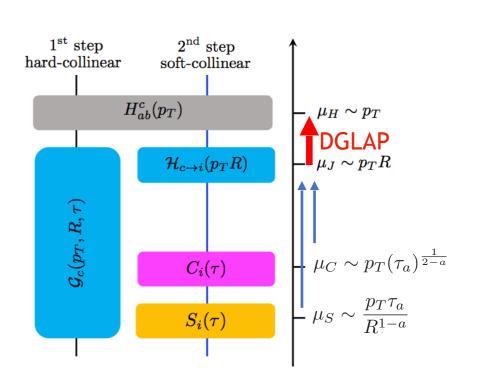
# Jet angularity

- Replace  $J_c(z, p_T R, \mu) \to \mathcal{G}_c(z, p_T R, \tau_a, \mu)$
- When  $au_a \ll R^2$ , Refactorize  $\mathcal{G}_c$  as

$$\mathcal{G}_{c}(z, p_{T}R, \tau_{a}, \mu) = \sum_{i} \mathcal{H}_{c \to i}(z, p_{T}R, \mu)$$

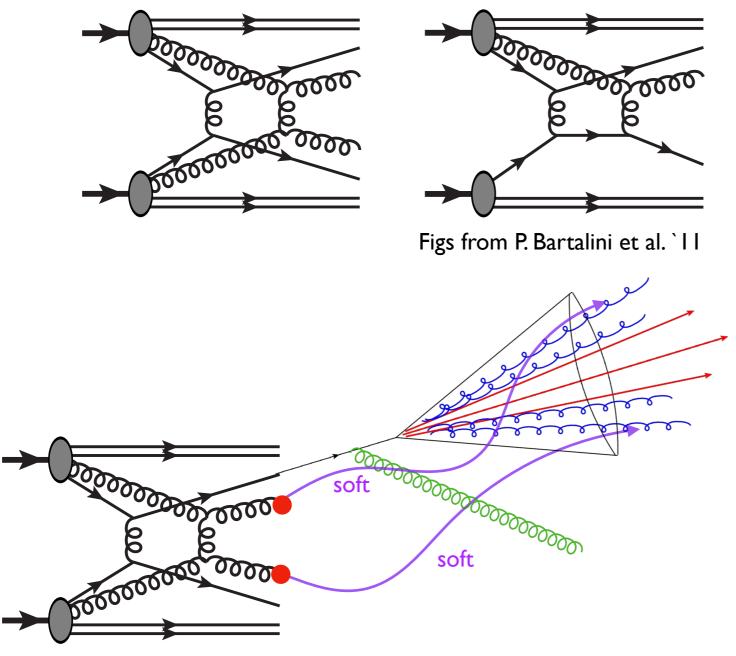
$$\times \int d\tau_{a}^{C_{i}} d\tau_{a}^{S_{i}} \delta(\tau_{a} - \tau_{a}^{C_{i}} - \tau_{a}^{S_{i}}) C_{i}(\tau_{a}^{C_{i}}, p_{T}\tau_{a}^{\frac{1}{2-a}}, \mu) S_{i}(\tau_{a}^{S_{i}}, \frac{p_{T}\tau_{a}}{R^{1-a}}, \mu) + \mathcal{O}\left(\frac{m^{2}}{p_{T}^{2}R^{2}}\right)$$

- Each pieces describe physics at different scales.
- $\mu_J \rightarrow \mu_H$  evolution follows DGLAP evolution equation again
- Resums  $(\alpha_s \ln R)^n$  and  $(\alpha_s \ln^2 \frac{R}{\tau_a^{1/(2-a)}})^n$



## Non-perturbative Effects

• Non-perturbative effects:

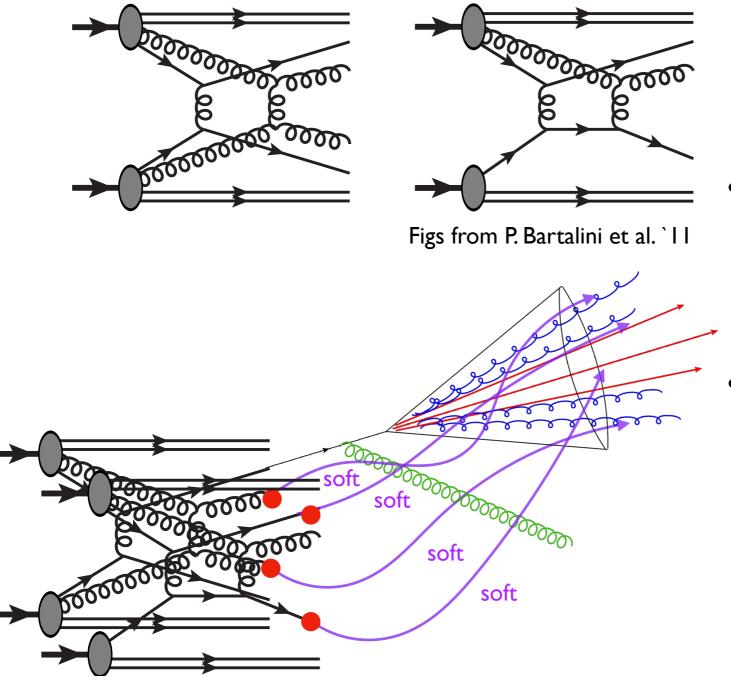


• Multi-Parton Interactions (MPI) (Underlying Events (UE)) Multiple secondary scatterings of

partons within the protons may enter and contaminate jet.

## Non-perturbative Effects

• Non-perturbative effects:



• Multi-Parton Interactions (MPI) (Underlying Events (UE)) Multiple secondary scatterings of

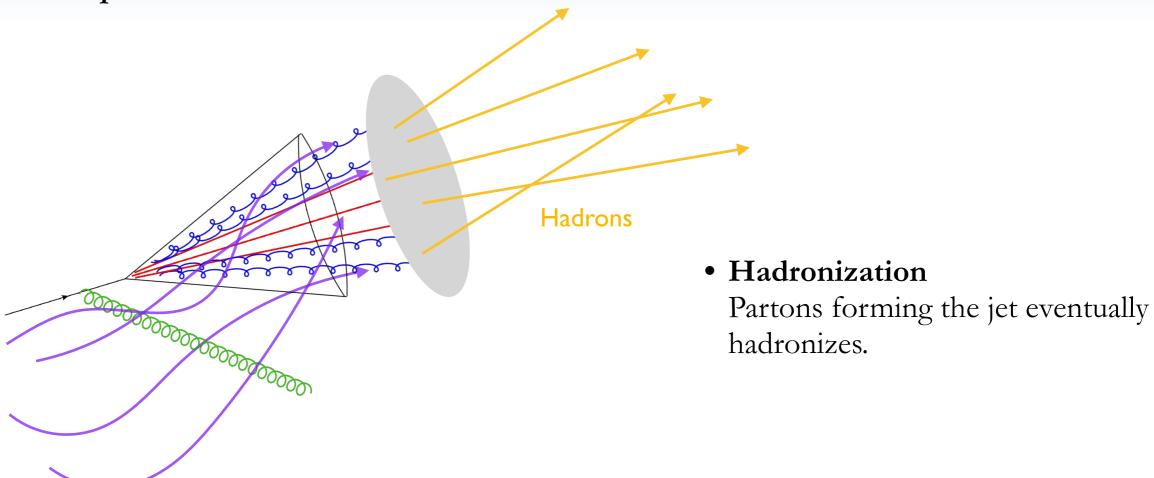
partons within the protons may enter and contaminate jet.

• Pileups

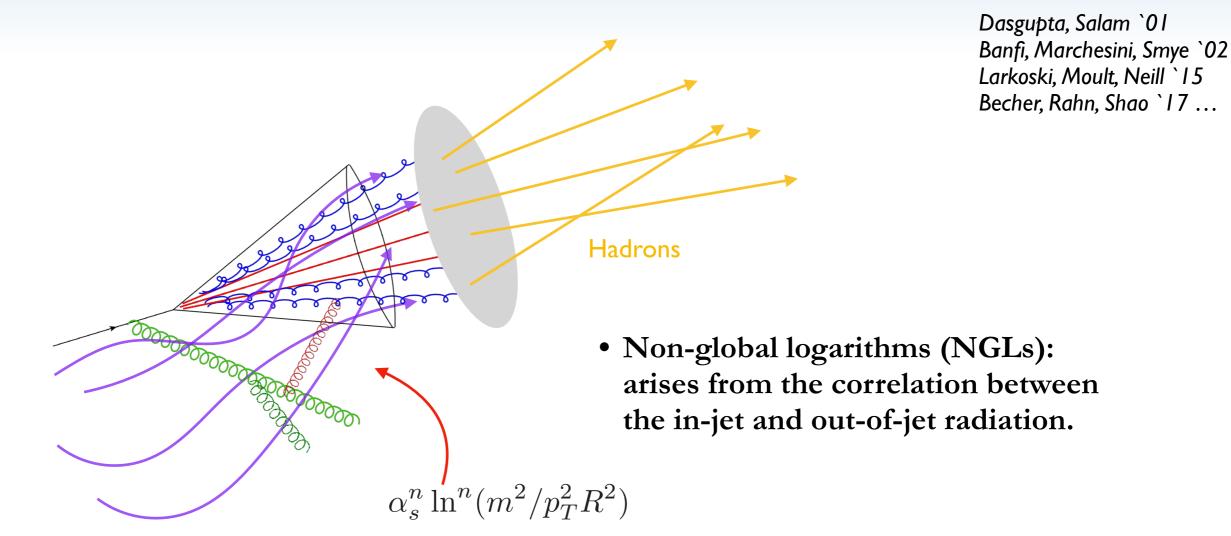
Secondary proton collisions in a bunch may enter and contaminate jet.

# Non-perturbative Effects

• Non-perturbative effects:



# Non-global logarithms



rather small effect for jet mass can be larger for the EIC

# Non-perturbative Model

• As  $\tau$  gets smaller,  $\mu_S \sim \frac{p_T \tau}{R}$  (smallest scale) can approach a non-perturbative scale.

We shift our perturbative results by convolving with non-perturbative shape function to smear

$$\frac{d\sigma}{d\eta dp_T d\tau} = \int dk F_{\kappa}(k) \frac{d\sigma^{\text{pert}}}{d\eta dp_T d\tau} \left(\tau - \frac{R}{p_T}k\right)$$

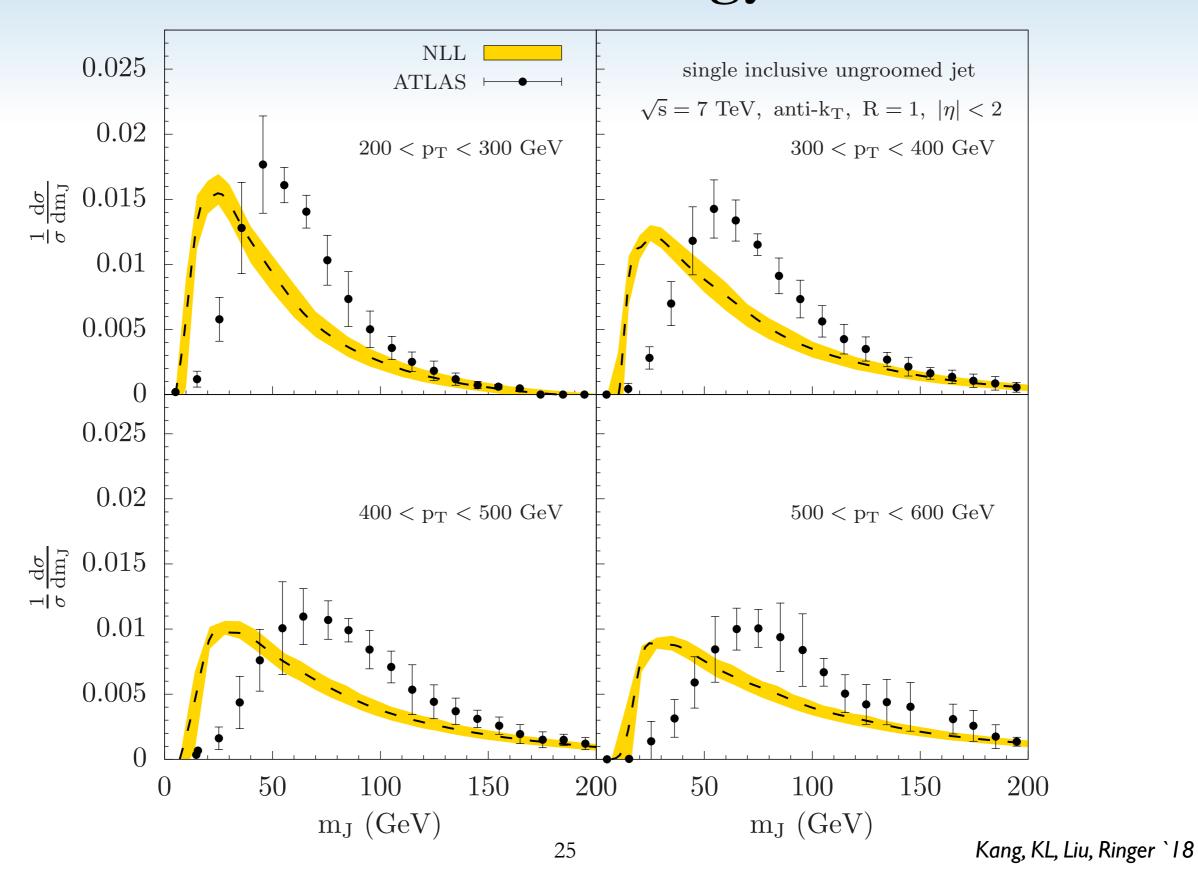
• Single parameter NP soft function:

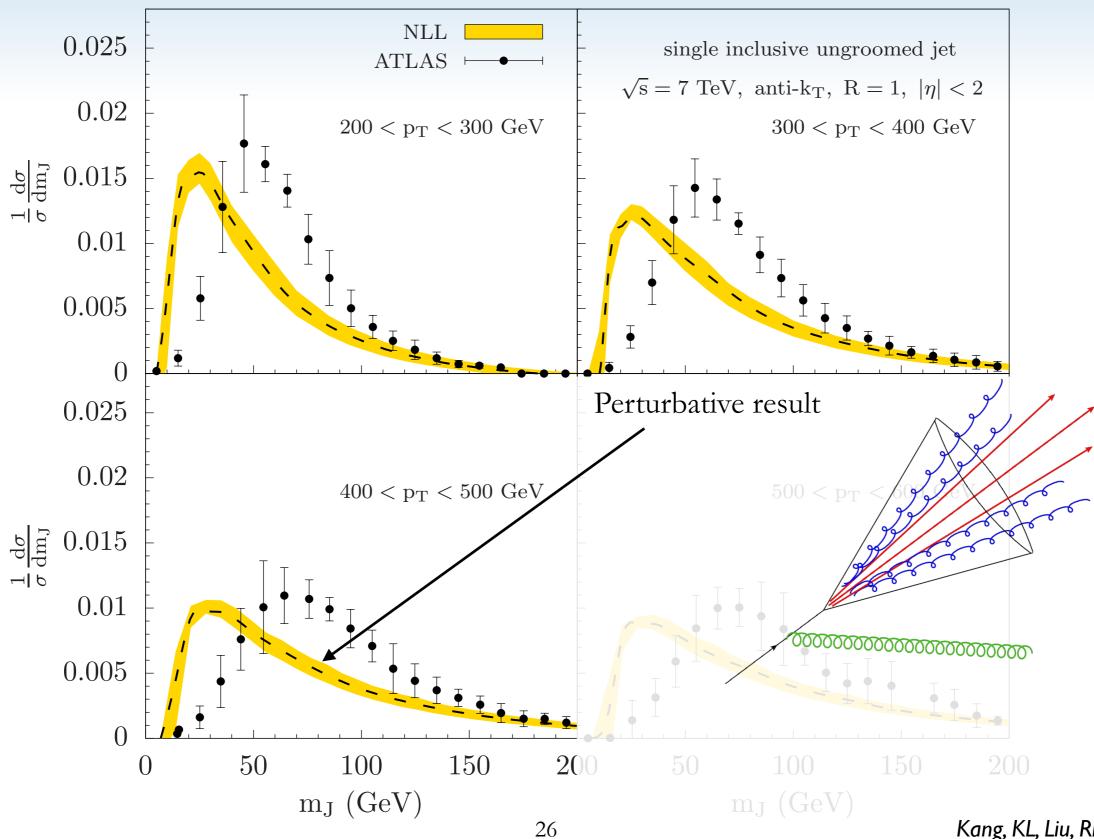
$$F_{\kappa}(k) = \left(\frac{4k}{\Omega_{\kappa}^2}\right) \exp\left(-\frac{2k}{\Omega_{\kappa}}\right) \qquad \text{Stewart, Tackmann, Waalewijn `15}$$

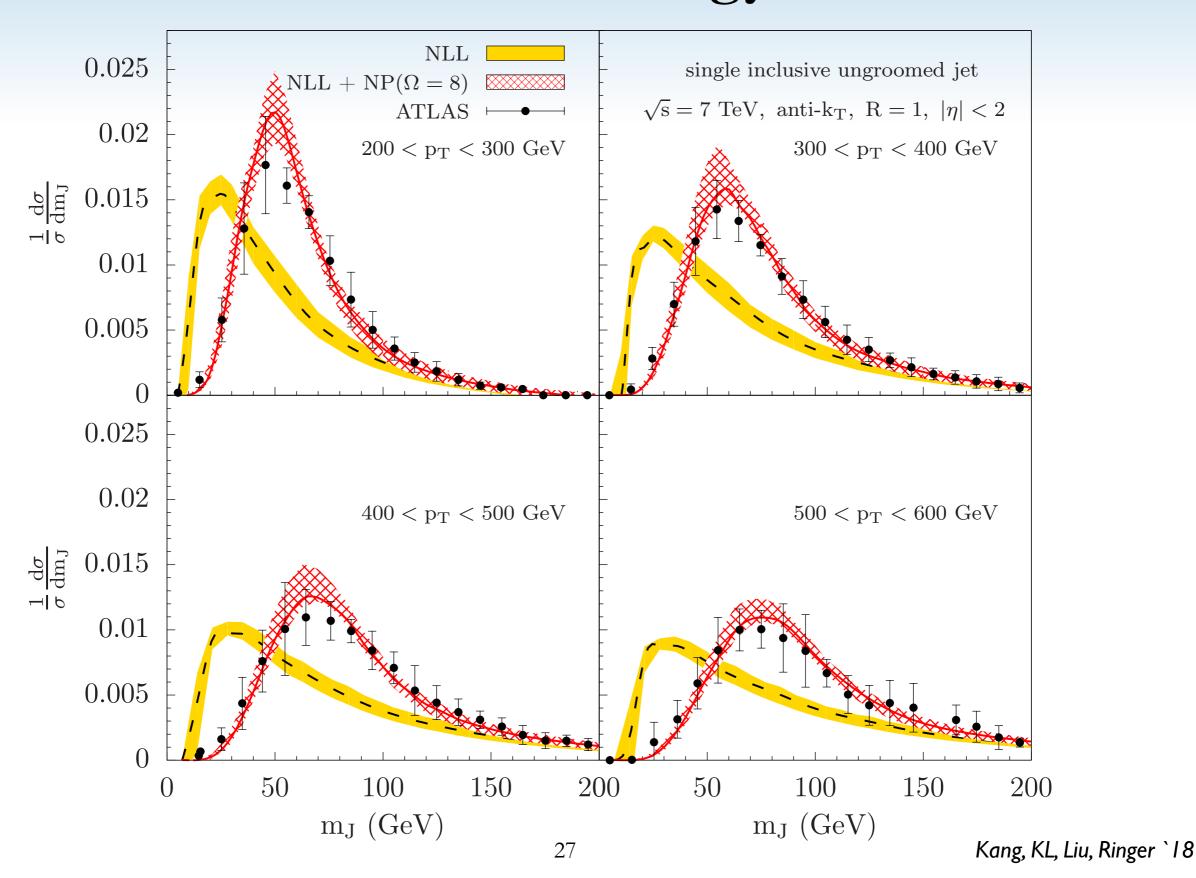
- Both hadronization and MPI effects in jet mass is well-represented by just shifting first-moments.
- The parameter  $\Omega_{\kappa}$  is related to shift in the distribution:

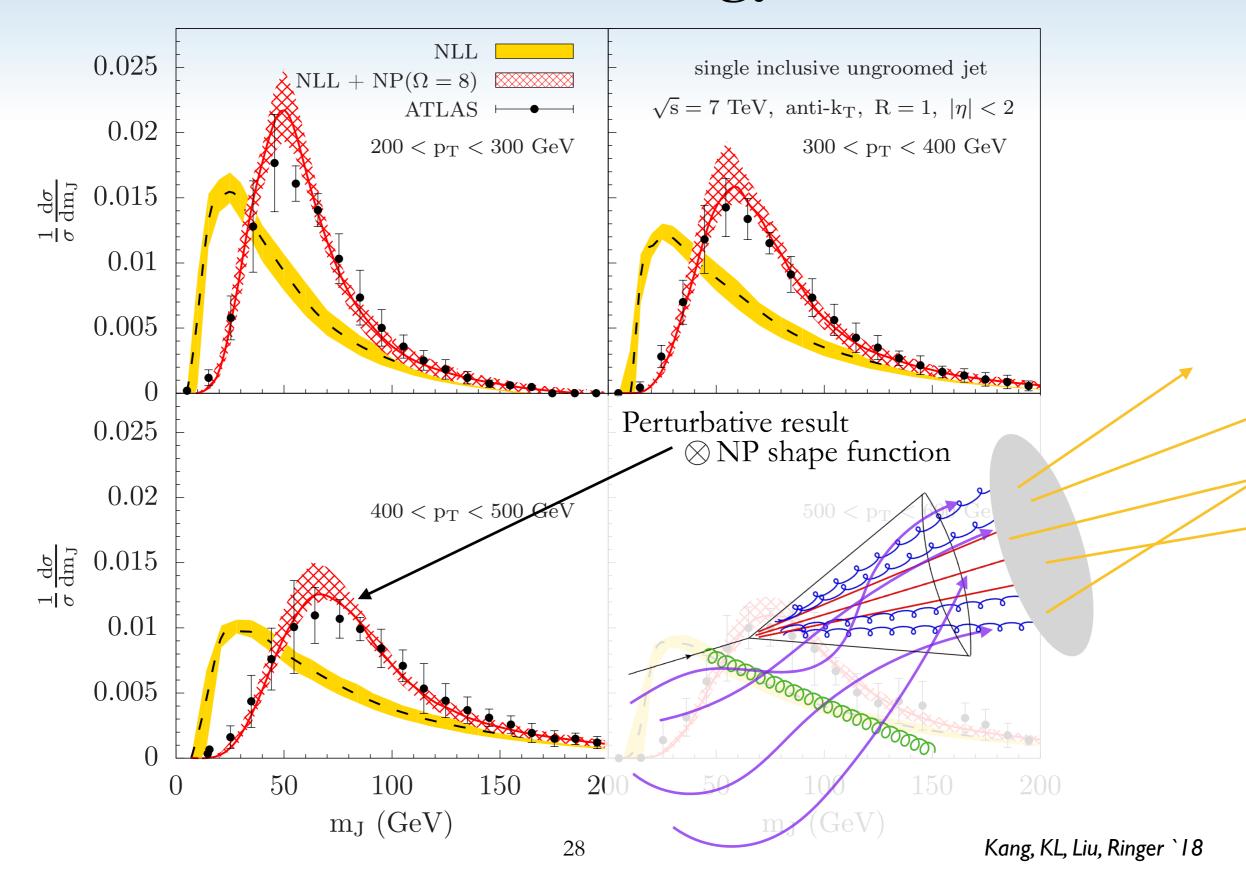
$$\tau = \tau_{\text{pert}} + \tau_{\text{NP}} = \tau_{\text{pert}} + \frac{R\Omega_{\kappa}}{p_T} = \tau_{\text{pert}} + \frac{R\left(\Lambda_{\text{hadro.}} + \Lambda_{\text{MPI}}\right)}{p_T}$$

 $\Omega_{\kappa} \sim \Lambda_{had} \sim 1 \, \mathrm{GeV}$  corresponds to non-perturbative effects coming primarily from the hadronization alone.









# Soft Drop Grooming

• Underlying Events (UE) are difficult to understand.

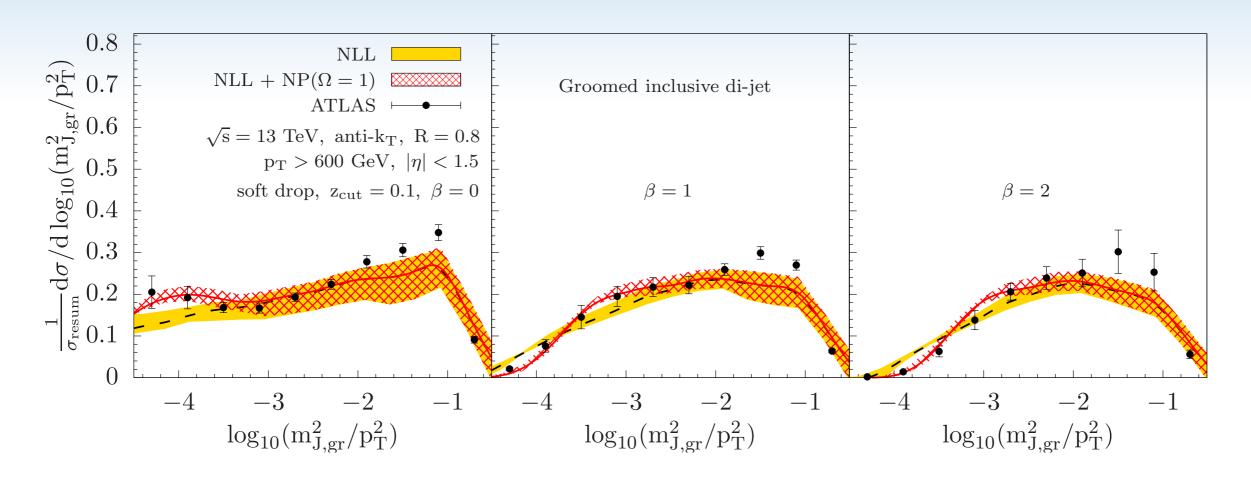
How do we get a better hold of these contaminations in the jet?

• Hint: contamination generally from soft radiations.

Groom jets to reduce sensitivity to wide-angle soft radiation.

Hadrons

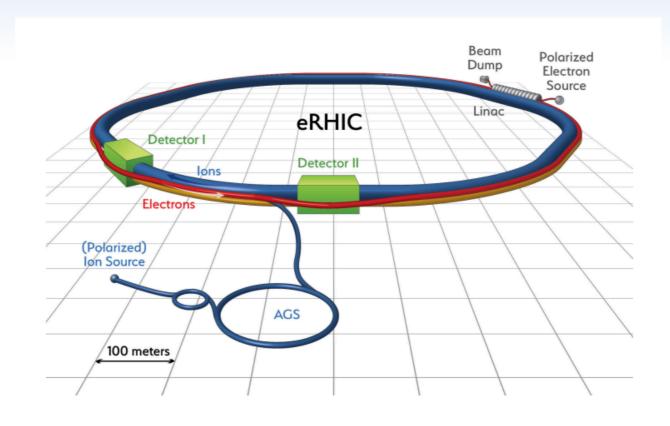
# Phenomenology (groomed jet mass)

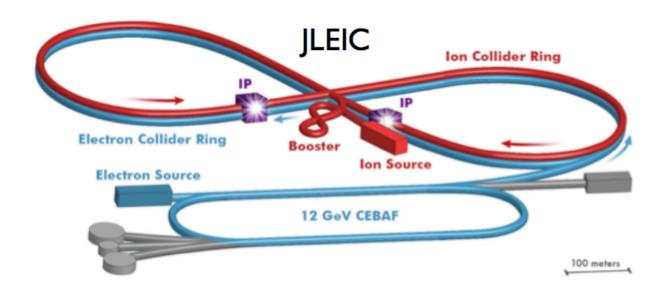


- Developed the formalism for single inclusive groomed jet mass cross-section.
- Shows very good agreement with the data.
- $\Omega_k = 1 \text{ GeV} \Longrightarrow \text{Reduced contamination as expected.}$ NP effects mostly from hadronization.

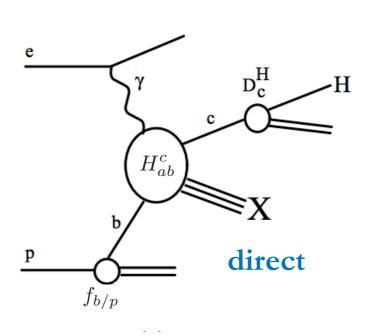
See also ATLAS, arXiv:1711.08341 Larkoski, Marzani, Soyez, Thaler `14 Frye, Larkoski, Schwartz, Yan `16

# Jets at the EIC





#### Photoproduction at the EIC



hadron

$$\frac{d\sigma^{ep\to ehX}}{dp_T d\eta} = \sum_{a,b,c} f_{a/l} \otimes f_{b/p} \otimes H^c_{ab} \otimes D^h_c$$
 
$$\frac{d\Delta\sigma^{ep\to ehX}}{dp_T d\eta} = \sum_{a,b,c} \Delta f_{a/l} \otimes \Delta f_{b/p} \otimes \Delta H^c_{ab} \otimes D^h_c$$
 Weizsäcker-Williams spectrum 
$$f_{a/l} = P_{\gamma l} \otimes f_{a/\gamma}$$

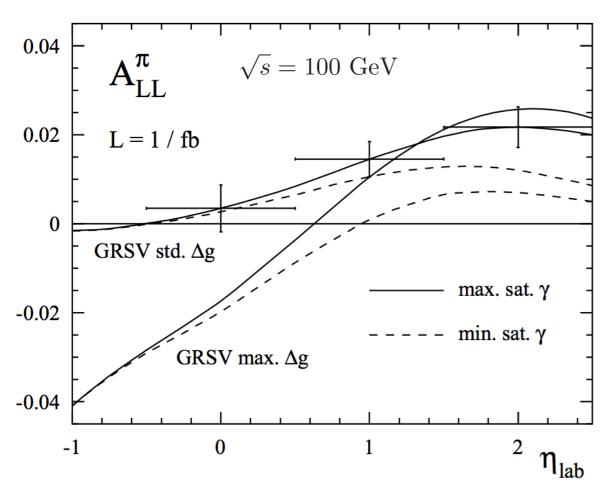
- For the direct process,  $f_{a/\gamma} = \delta(1 x_{\gamma})$ .
- Observe outgoing lepton to tag  $Q^2$
- Require high  $p_T$  and  $Q^2 < 0.1 \text{ GeV}^2$  (near on-shell photon)

See Jäger, Stratmann, Vogelsang `03

#### Polarized Gluon and Photon PDF

Study in 2003,

Jäger, Stratmann, Vogelsang `03



$$A_{LL} = \frac{d\Delta\sigma}{d\sigma} = \frac{d\sigma_{++} - d\sigma_{+-}}{d\sigma_{++} + d\sigma_{+-}}$$
$$\Delta f_{\text{max}} = f \qquad \Delta f_{\text{min}} = 0$$

- Sensitivity to polarized gluon pdf at low  $\eta_{lab}$
- Sensitivity to polarized photon pdf at high  $\eta_{lab}$

Assumptions:  $D_c^{\pi^0}$  has been well-determined.

Use inclusive jets as a perturbative probe!

For using dijet process, see Xiaoxuan's talk

• Study of polarized pdfs

$$\frac{d\Delta\sigma^{ep\to e\pi^0X}}{dp_Td\eta} = \sum_{a,b,c} \Delta f_{a/l} \otimes \Delta f_{b/p} \otimes \Delta H_{ab}^c \otimes D_c^{\pi^0}$$

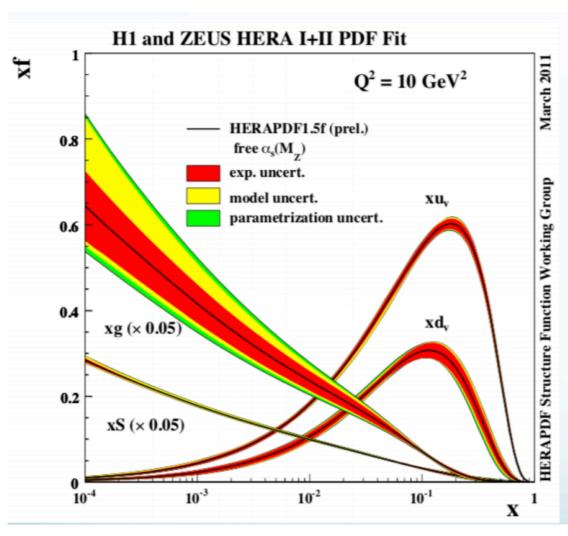
# HERA PDF fit with and without jets

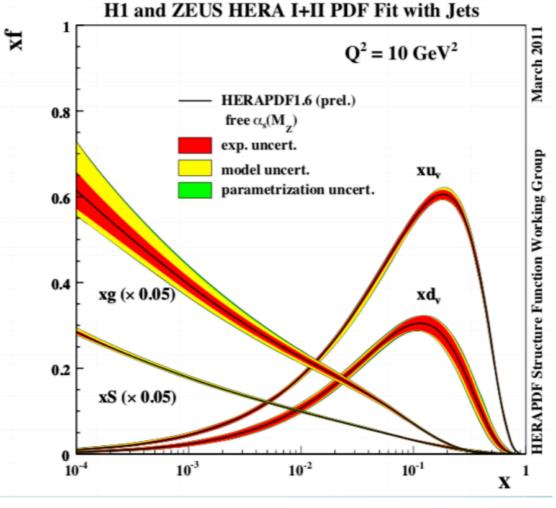
• Important for constraining gluon PDF

Nuclear Physics B (Proc. Suppl.) 222—224 (2012) January—March 2012

#### HERA 2011

Proceedings of the Ringberg Workshop New Trends in HERA Physics 2011

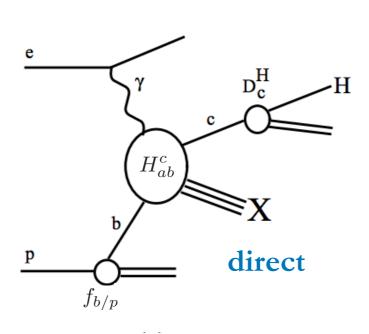




Without jets

With jets

#### Photoproduction at the EIC

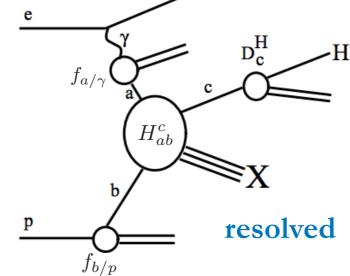


hadron  $\frac{aa}{a}$ 

$$rac{d\sigma^{ep o ehX}}{dp_T d\eta} = \sum_{a,b,c} f_{a/l} \otimes f_{b/p} \otimes H^c_{ab} \otimes D^h_c$$
 \quad \text{Weizsäcker-Williams spectrum} \, \frac{d\Delta \sigma^{ep}}{dp\_T} \, \frac{d\Delta \si

- For the direct process,  $f_{a/\gamma} = \delta(1 x_{\gamma})$ .
- Observe outgoing lepton to tag  $Q^2$
- Require high  $p_T$  and  $Q^2 < 0.1 \text{ GeV}^2$  (near on-shell photon)

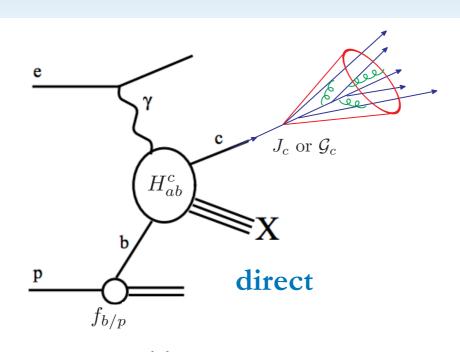
See Jäger, Stratmann, Vogelsang `03

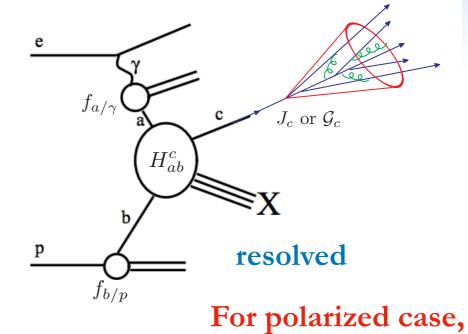


For polarized case,

$$\frac{d\Delta\sigma^{ep\to ehX}}{dp_T d\eta} = \sum_{a,b,c} \Delta f_{a/l} \otimes \Delta f_{b/p} \otimes \Delta H_{ab}^c \otimes D_c^h$$

## Photoproduction at the EIC





hadron

$$\frac{d\sigma^{ep\to ehX}}{dp_T d\eta} = \sum_{a,b,c} f_{a/l} \otimes f_{b/p} \otimes H^c_{ab} \otimes D^h_c$$

 $\frac{d\Delta\sigma^{ep\to ehX}}{dp_T d\eta} = \sum_{a,b,c} \Delta f_{a/l} \otimes \Delta f_{b/p} \otimes \Delta H^c_{ab} \otimes D^h_c$ 

$$\frac{d\sigma^{ep\to e \text{jet}X}}{dp_T d\eta} = \sum_{a,b,c} f_{a/l} \otimes f_{b/p} \otimes H^c_{ab} \otimes J_c + \mathcal{O}(R^2)$$

$$\frac{d\sigma^{ep\to e \text{jet}(m_J)X}}{dp_T d\eta dm_J} = \sum_{a,b,c} f_{a/l} \otimes f_{b/p} \otimes H^c_{ab} \otimes \mathcal{G}_c(m_J) + \mathcal{O}(R^2)$$

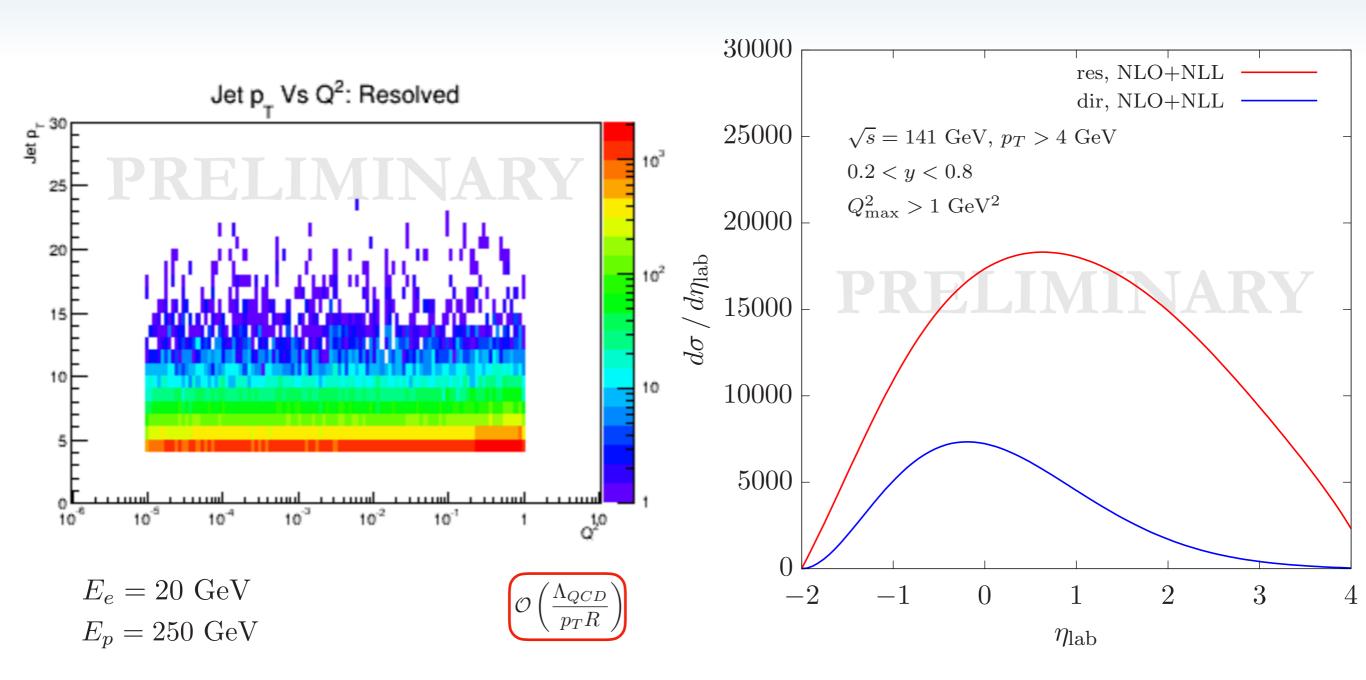
- Sensitivity to the photon pdfs. Can be done for polarized and unpolarized case.
- Quark and gluon discrimination with jet mass observed.

Jäger, Stratmann, Vogelsang `03

• Role of NP physics?

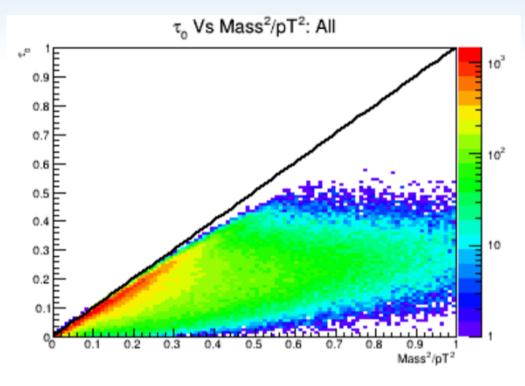
Chu, Aschenauer, Lee, Zheng `I7 In collaboration with Elke Aschenauer and Brian Page

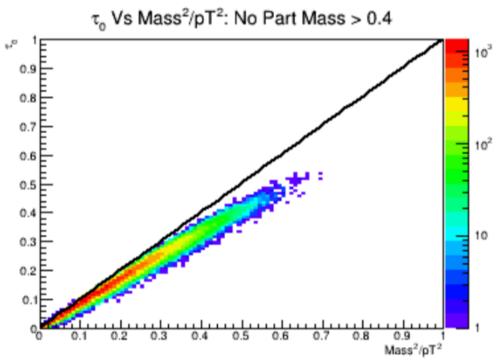
# $p_T$ distribution for the jets in the EIC

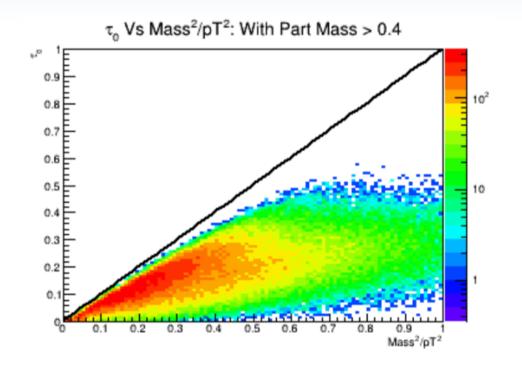


• 4 GeV  $< p_T < 15$  GeV for  $Q^2 < 1$  GeV, contribution mostly from resolved.

#### Power corrections



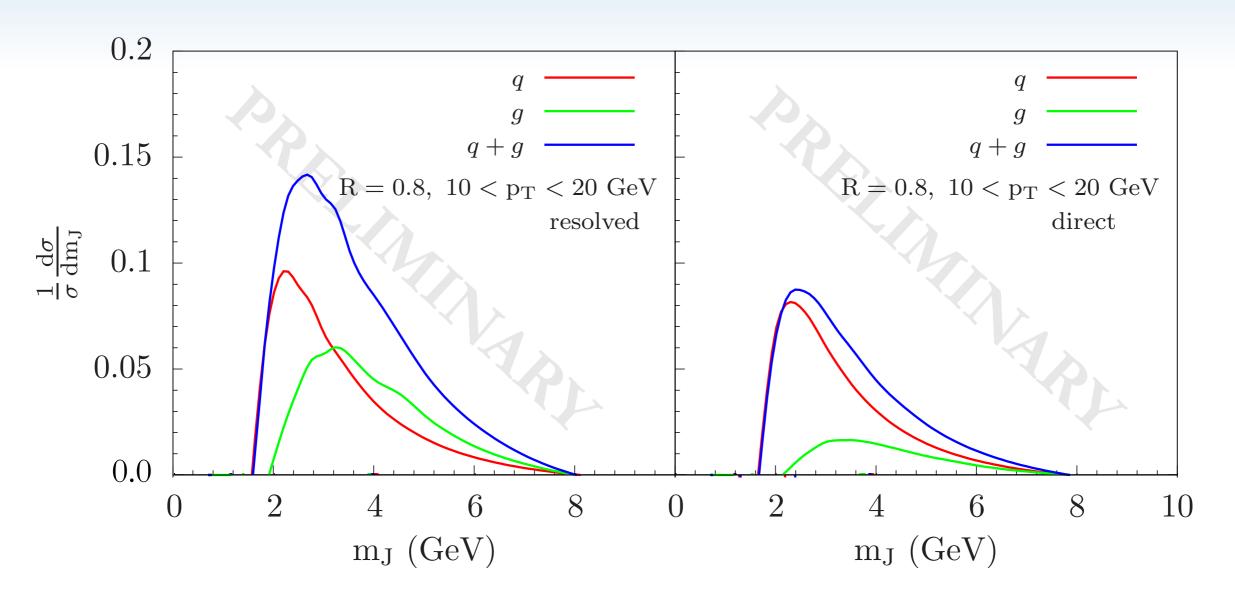




$$\tau_0^{pp} = \frac{m_J^2}{p_T^2} + \mathcal{O}((\tau_0^{pp})^2)$$

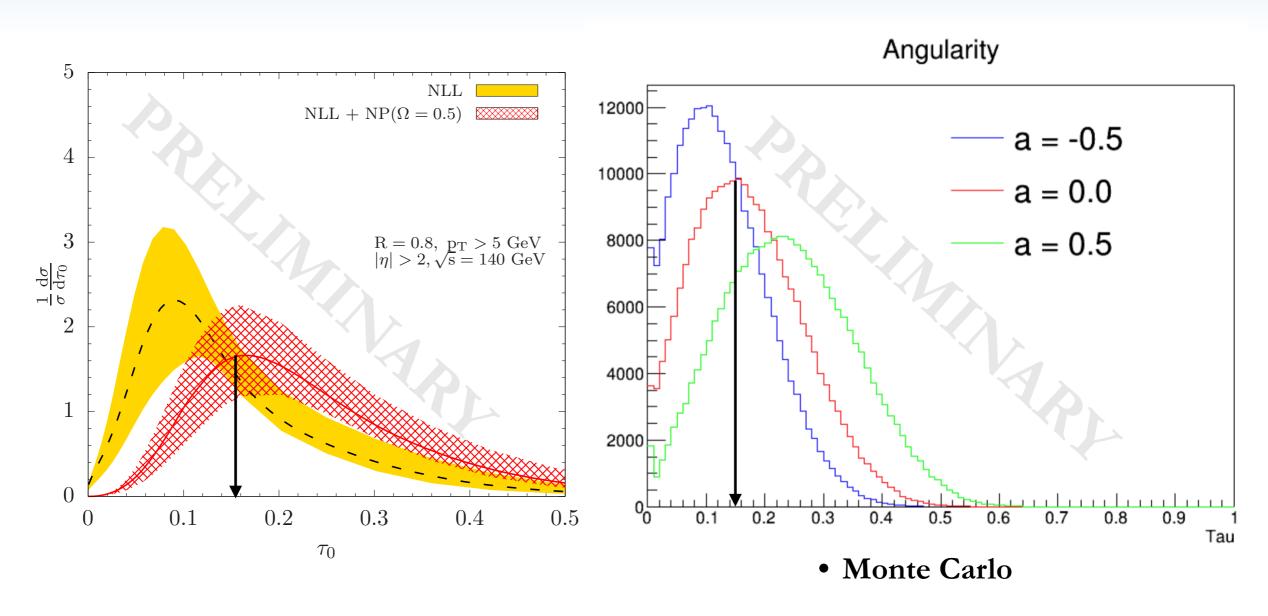
• Mass effects: important challenge and opportunity

## **Preliminary Plots**



• Fraction of gluon contribution is reduced for the direct process relative to the resolved process.

## **Preliminary Plots**



•  $\Omega_{\kappa} = 0.5 \text{ GeV}$ , assumption that NP effects only come from the hadronization gives the right peak value  $\Longrightarrow$  less contamination from UE than LHC

#### **Conclusions**

- Formalisms for studying semi-inclusive jet production with and without a substructure measurement were introduced.
- Discussed phenomenology of jet mass in the LHC.
- Discussed various non-perturbative effects.
- Jets have tremendous success at the LHC, interesting opportunities and new challenges at the EIC:

power corrections, reduced contaminations, constraining unpolarized and polarized pdfs.