

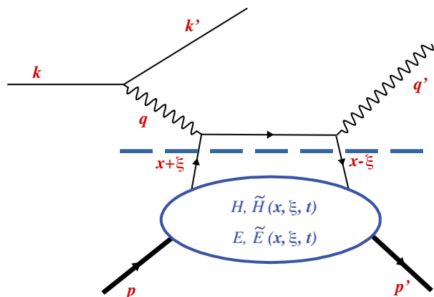
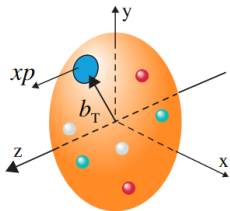
# Recoil proton polarization in DVCS

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# Proton structure and GPDs

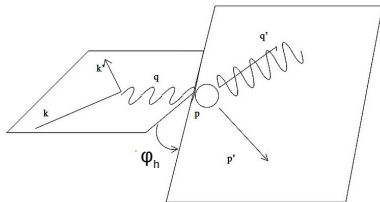
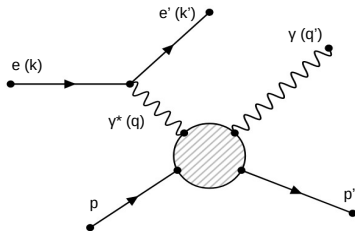


Generalized Parton Distributions (GPDs): nucleon structure in terms of longitudinal momentum & transverse position.

- Measured in exclusive processes like Deeply Virtual Compton Scattering (DVCS).

**Factorization:** splitting into perturbative hard part + non-perturbative soft part (GPDs).

GPDs are accessible through **Compton Form Factors (CFFs)**.

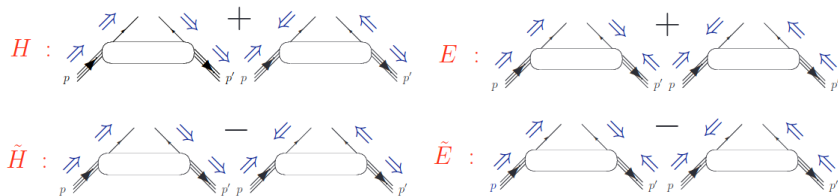


## DVCS parameters

- $E_k$ : beam energy.
- $Q^2$ : virtuality of the photon,  $Q^2 = -q^2 = -(k - k')^2$ .
- $x_B = \frac{Q^2}{2p \cdot q}$ .
- $t$ : 4-momentum transfer to the proton,  $t = (p' - p)^2$
- $\phi_h$ : angle between the leptonic and hadronic planes.

These parameters determine the kinematics of the scattered particles.

# GPD meanings



Connection of  $E$  to the total orbital angular momentum of the quarks through Ji's sum rule:  $J^q = \frac{1}{2} \int dx x [H^q(x, \xi, t = 0) + E^q(x, \xi, t = 0)]$

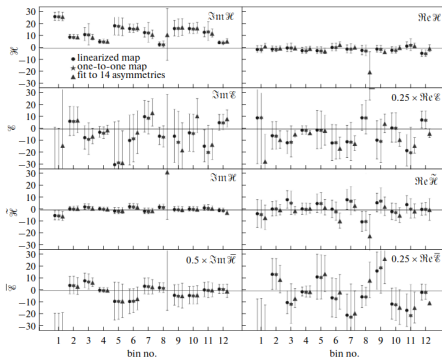
$E$  describes the process when the proton changes helicity.

Can we introduce a new observable to determine  $E$  to higher precision?

# Measuring CFFs

## The common setup to measure CFFs

- $\mathcal{H}$ : unpolarized target
- $\tilde{\mathcal{H}}$ : longitudinally polarized target
- $\mathcal{E}$ : transversely polarized target - challenge!
  - Gaseous transversely polarized target by HERMES (low luminosity).



- $\mathcal{E}$  is much less constrained than  $\mathcal{H}$ ,  $\tilde{\mathcal{H}}$ .
- $\mathcal{E}$  can also be measured through DVCS on a deuterium target, but theoretical complications.
- Can we extract more info on  $\mathcal{E}$  by measuring the **polarization** of the **recoil proton**?

# Polarization

- Code by Pierre Guichon, using the exact mathematical expressions.
- The polarization  $\vec{P} = (P_x, P_y, P_z)$  is computed for the DVCS+Bethe-Heitler process, including their interference.

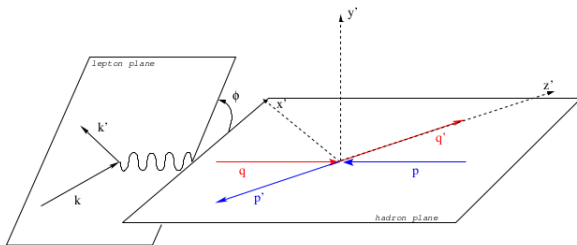
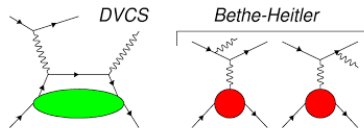
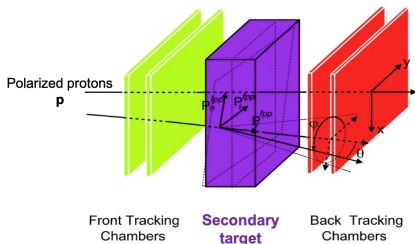


Figure 1: Rotated CM frame  $(x'y'z')$

- $P_y$  (normal to the hadronic plane) is sensitive to CFF  $\mathcal{E}$ .
- Additionally,  $P_x$  is sensitive to  $\tilde{\mathcal{H}}$ .

# Proton polarimeter



- Rescatter the proton with  $\theta_{pol}$ ,  $\phi_{pol}$  inside a carbon analyzer.
- A set of trackers before and after the analyzer detect the incoming and outgoing protons.

Transverse polarization will induce a modulation:

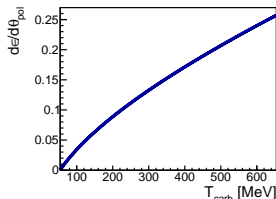
$$N(\theta_{pol}, \phi_{pol}) = N_0 \frac{d\epsilon}{d\theta_{pol}} [1 + A_p(\theta_{pol})(P_y \cos \phi_{pol} - P_x \sin \phi_{pol})]$$

- Polarization  $\propto$  amplitude.
- At  $\phi_h = 180^\circ$ :  $P_x \rightarrow -P_x$ ,  $P_y \rightarrow P_y$  when beam helicity switches.
  - Subtract over different beam helicity  $\rightarrow P_x^m (\tilde{\mathcal{H}})$ .
  - Add all events  $\rightarrow P_y (\mathcal{E})$ .
  - Two experiments in one.

# Polarimeter performance

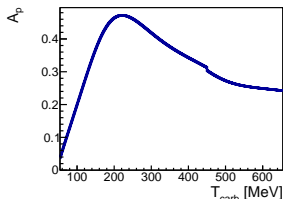
## Efficiency

Probability to have a useful scattering in the analyzer, Bonin et al.



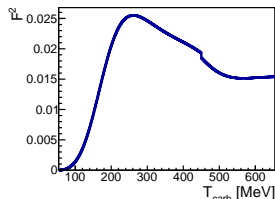
## Analyzing power

Sensitivity of the scattering to the polarization, McNaughton et al.



## Figure of merit

$$F^2 = \int_{\theta_{min}}^{\theta_{max}} A_p(\theta)^2 \epsilon(\theta) d\theta$$



The polarimeter performance depends on the proton momentum, scattering angle in the analyzer  $\theta_{pol}$  and analyzer thickness (here 15 cm).



## Challenges

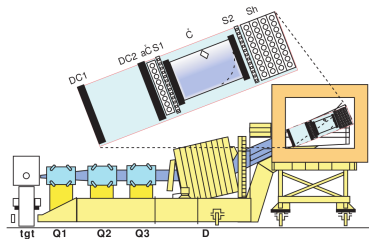
- DVCS has a low cross section compared to (semi-)inclusive processes.
- We need to rescatter the recoil proton, expecting a polarimeter efficiency of order 0.1.
- To achieve high statistics we need a high luminosity.

Candidate facility: **Hall C** at **Jefferson Lab**, using a **polarized electron beam** and an **unpolarized liquid hydrogen target**.

## Example settings

- Target length: 15 cm.
- Beam current: 10  $\mu\text{A}$  during 3 weeks of data taking.

# Electron and photon detection



## Electron detection: HMS

- Focusing spectrometer.
- Scattering angle acceptance:  $10.5\text{--}80^\circ$ .
- Angular acceptance:  $\pm 1.8^\circ$  in-plane,  $\pm 4.9^\circ$  out-of-plane.
- Momentum acceptance  $\pm 10\%$ .



## Photon detection: NPS calorimeter

- Angular acceptance:  $\pm 5.3^\circ$  horizontally,  $\pm 6.7^\circ$  vertically.
- $30 \times 36$   $\text{PbWO}_4$  crystals. Position resolution: 2-3 mm.
- Sweeping magnet reduces low energy electron background.

## Procedure

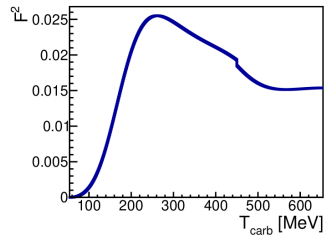
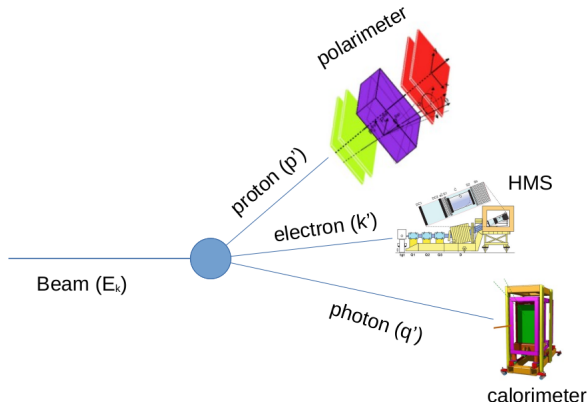
- Consider different DVCS configurations (5 parameters fully determine the particle kinematics).
  - Evaluate the CFFs using GK model and compute the differential cross section and polarization.
  - Construct a figure of merit  $\mathcal{F}'$  proportional to  $\Delta P_y = P_y(\mathcal{E}) - P_y(0)$ :
$$\mathcal{F}' = F \cdot \sqrt{N_{\text{protons}}} \cdot \Delta P_y .$$
  - Take detector effects into account using Geant4 simulation.
- 
- Require isolation of all particles from each other ( $10^\circ$  in general,  $10.5^\circ$  electron-beam,  $30^\circ$  proton-electron).
  - Imposing constraint  $|t|/Q^2 \leq 0.25$  for validity of leading twist.

# Optimization result

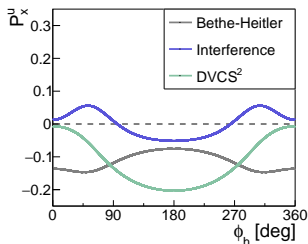
Maximizing  $\mathcal{F}'$  gives:

$$E_k = 10.6 \text{ GeV}, Q^2 = 1.8 \text{ GeV}^2, x_B = 0.17, t = -0.45 \text{ GeV}^2, \phi_h = 180^\circ$$

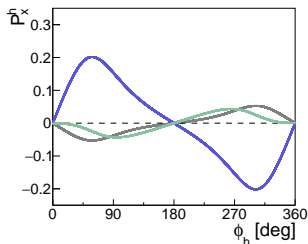
electron $ k' $	$\theta_{k'}$	photon $ q' $	$\theta_{q'}$	proton $ p' $	$E_{carb}$	$\theta_{p'}$
4.96 GeV/c	$10.6^\circ$	5.40 GeV/c	$-15.1^\circ$	0.71 GeV/c	0.19 GeV/c	$44^\circ$



# Decomposition as a function of $\phi_h$



Beam helicity  
independent ( $P_i^u$ )

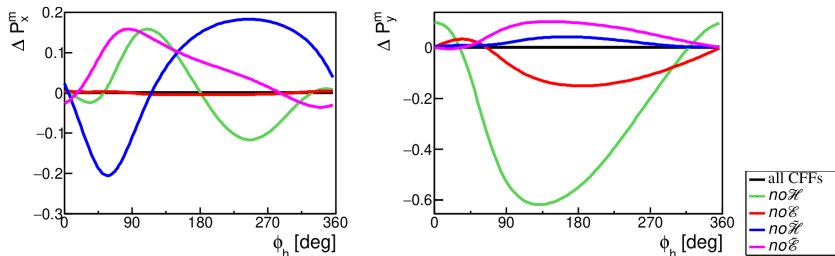


Beam helicity  
dependent ( $P_i^h$ )

$$P_x = h_e(P_x^u + h_e P_x^h)$$

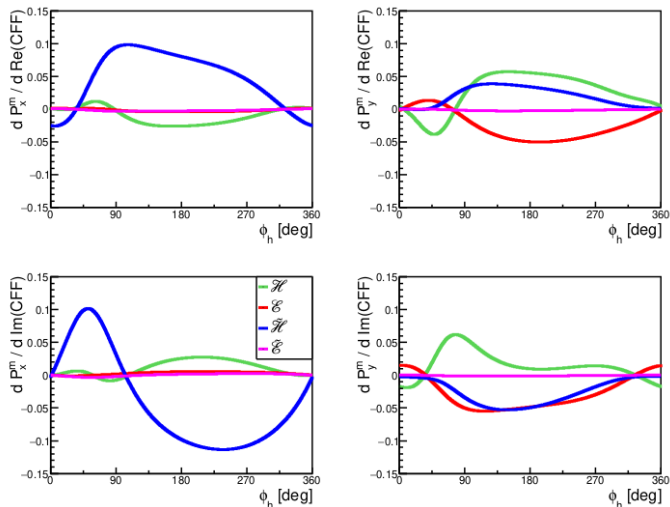
$$P_y = P_y^u + h_e P_y^h$$

# $\phi_h$ dependence



- Looking at the difference in polarization  $\Delta P_i$  between the GK model and GK model with one coefficient set to 0.
- Non-trivial  $\phi_h$  dependence.
- The polarization appears sensitive to all CFFs.
- The difference depends on the values of the CFFs in the GK model.

# dP/dCFF derivatives, $\phi_h$



Derivatives w.r.t.  $\Re(CFF)$ ,  $\Im(CFF)$  removes the model dependence of the interference terms and reduces that of the (bilinear) DVCS<sup>2</sup> terms.

# Local expression

While the full expression is super lengthy, we have derived a local one, showing the kinematical coefficients.

At  $E_k=10.6$  GeV,  $Q^2=1.8$  GeV<sup>2</sup>,  $x_B=0.17$ ,  $t=-0.45$  GeV<sup>2</sup>,  $\phi_h=180^\circ$ :

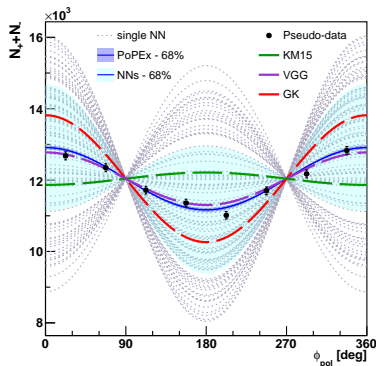
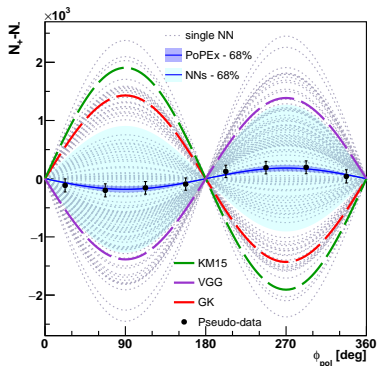
$$\begin{aligned}\mathcal{M}(x) = & -20.42 + 19.06 \operatorname{Re}\tilde{\mathcal{H}} + 7.15 \operatorname{Re}\mathcal{H} - 1.04 \operatorname{Re}\mathcal{E} - 0.56 \operatorname{Re}\tilde{\mathcal{E}} \\ & - 2.93 (\mathcal{H}\tilde{\mathcal{H}}^* + \mathcal{H}^*\tilde{\mathcal{H}}) + 0.16 (\mathcal{E}\tilde{\mathcal{H}}^* + \mathcal{E}^*\tilde{\mathcal{H}}) \\ & + 0.04 (\mathcal{H}\tilde{\mathcal{E}}^* + \mathcal{H}^*\tilde{\mathcal{E}}) + 0.03 (\mathcal{E}\tilde{\mathcal{E}}^* + \mathcal{E}^*\tilde{\mathcal{E}})\end{aligned}\quad (1)$$

$$\begin{aligned}\mathcal{M}(y) = & 15.50 \operatorname{Im}\mathcal{H} - 10.05 \operatorname{Im}\mathcal{E} + 3.44 \operatorname{Im}\tilde{\mathcal{H}} - 0.44 \operatorname{Im}\tilde{\mathcal{E}} \\ & + 1.51 \operatorname{Im}(\mathcal{E}\mathcal{H}^* - \mathcal{E}^*\mathcal{H}) + 0.14 \operatorname{Im}(\tilde{\mathcal{E}}\tilde{\mathcal{H}}^* - \tilde{\mathcal{E}}^*\tilde{\mathcal{H}})\end{aligned}\quad (2)$$

$$P^i = \mathcal{M}(i)/\mathcal{M}(0).$$



# Fitting the polarization



## Toy simulation of the polarimeter

- We assume a  $\pm 20^\circ \times \pm 30^\circ$  ( $\theta_{p'} \times \phi_{p'}$ ) polarimeter and simulate.
- Extract  $P_x$ ,  $P_y$  from the difference, sum over beam helicity.
- Statistical precision excludes  $\sim 90\%$  of the NN replicas from the PARTONS team and discriminate between GPD models.

# Optimizing wrt derivatives pre-simulation

Target	$E_k$	$Q^2$	$x_B$	t	$dP_y/dCFF$	dxs [pb/..]	F'
$\Re(E)$	10.6	1.9	0.19	-0.47	0.061	2.7	1.45
$\Im(E)$	8.8	1.9	0.15	-0.47	0.045	5.5	1.53
$\Im(E)$	10.6	2.2	0.17	-0.55	0.067	1.9	1.46
Ref ( $\Re$ )	10.6	1.8	0.17	-0.45	0.050	3.8	1.33
Ref ( $\Im$ )	"	"	"	"	0.049	"	1.31

$$F' = d\sigma \cdot (dP_y/dCFF)^2 \cdot \text{FoM}_{pol} \cdot 100$$

# Optimizing wrt derivatives pre-simulation

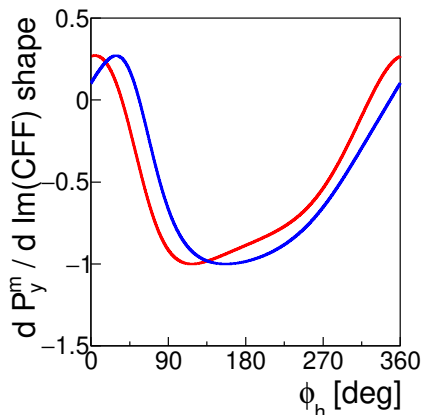
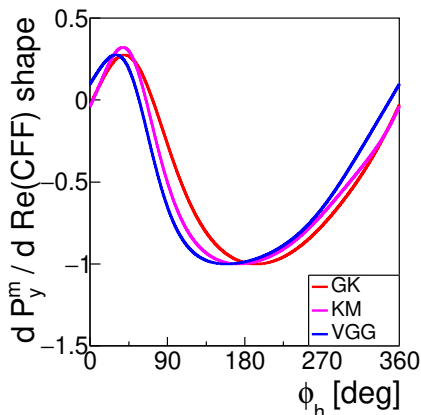
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$$F' = d\sigma \cdot (dP_y/dCFF)^2 \cdot \text{FoM}_{pol} \cdot 100$$

DVCS parameters	$p'_e, \theta'_e$	$p'_\gamma, \theta'_\gamma$	$p'_p, \theta'_p$
10.6,1.9,0.19,-0.47	5.27, 10.6	5.08, -16.6	0.730, 41
8.8,1.9,0.15,-0.47	2.05, 18.7	6.50, -10.8	0.730, 50
10.6,2.2,0.17,-0.55	3.70, 13.6	6.60, -12.6	0.797, 46
10.6,1.8,0.17,-0.45	4.96, 10.6	5.40, -15.1	0.712, 44

Note: The HMS has a higher acceptance at higher electron momentum.

# Derivative for different input models



The value at 180° is near the highest absolute value of the derivatives for the VGG and KM models too.

## Summary

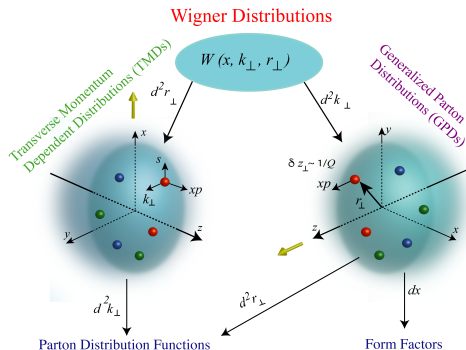
- We have explored a new way of measuring  $\mathcal{E}$  - by looking at the **polarization** of the **recoil proton** in DVCS.
- $P_y$  is highly sensitive to  $\mathcal{E}$  and to different models.
- **Very high statistics** for 1 sr polarimeter, 3 weeks of data-taking,  $10\mu A$ .
- **Good discrimination** between the GK, VGG and KM15 models in the statistical analysis. Potential to exclude many NN replicas.
- A starting point for a proposal has been identified.
- Robustness checks:
  - The optimization result does not change by much when optimizing w.r.t. derivatives rather than  $P_y - P_y(E=0)$  in GK.
  - For VGG and GK models,  $\phi_h = 180^\circ$  is also a good choice.

Publication in PRD: [PhysRevD.107.014020](https://arxiv.org/abs/PhysRevD.107.014020)

- Consider the **background** and how to reduce it.
- Develop a polarimeter design and determine polarimeter dimensions.
  - The Super Big-bite Spectrometer will be moved to Hall C. Can the GEM detectors run without magnetic field at 44 degrees 1.5-2m away from the target?
  - Another possibility: scintillating fibers active polarimeter.
- Prepare a proposal for Hall C.



# GPDs and proton structure



## GPDs

- GPDs and TMDs can both be obtained from Wigner distributions.

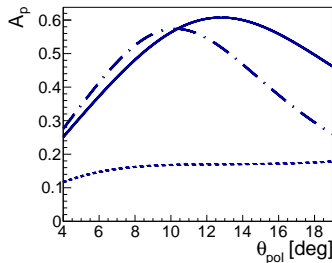
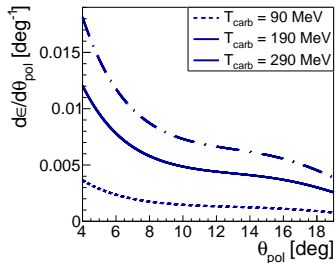


From helicity-dependent cross-section difference:

$$[\mathcal{C}_n^I]^{exp} \simeq [\mathcal{C}_n^I] = F_1 \mathcal{H} + \xi(F_1 + F_2) \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E}$$

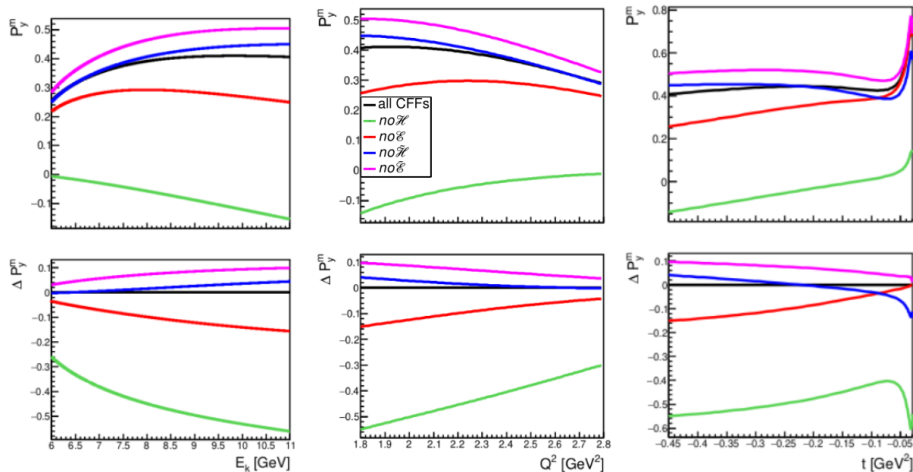
- Dominated by  $\mathcal{H}$ ,  $\tilde{\mathcal{H}}$  for a proton.
- In neutron:
  - $F_1$  is small.
  - Cancellation between  $u$ ,  $d$  polarized distributions in  $\tilde{\mathcal{H}}$ .

# Polarimeter properties in our region of interest



- The efficiency increases with proton energy.
- The analyzing power peaks around a kinetic energy of  $\sim 200$  MeV.

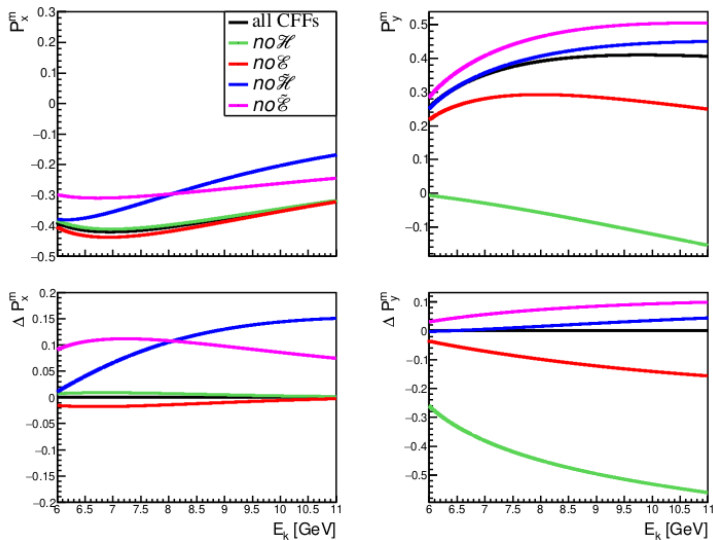
# Variations around the proposed kinematics



The figure of merit favours high beam energy  $E_k$ , low photon virtuality  $Q^2$  and high momentum transfer to the proton  $|t|$ .

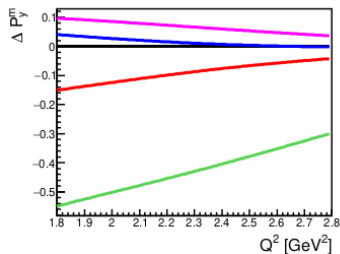
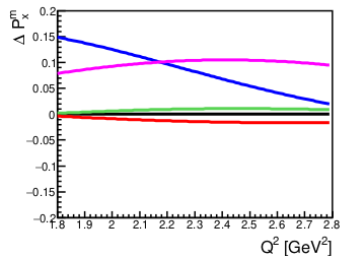
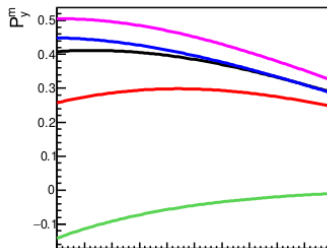
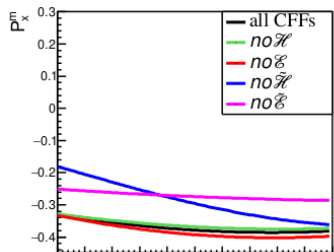
Imposing constraint  $|t|/Q^2 \leq 0.25$  for validity of leading twist.

# Variations: beam energy



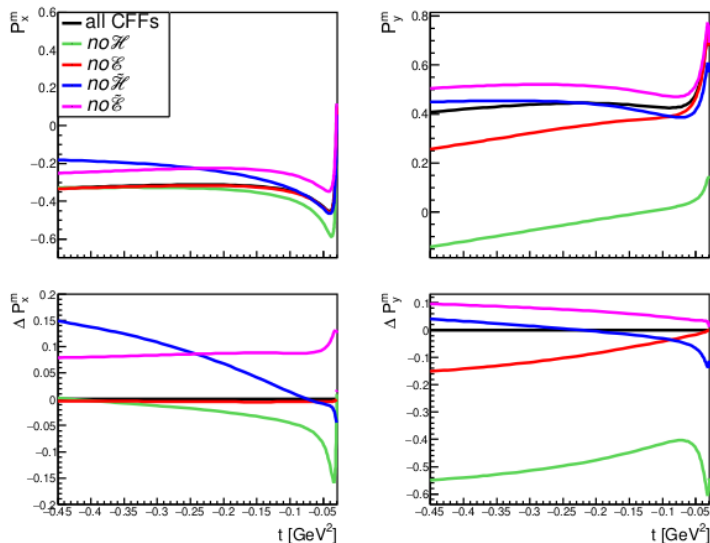
A high beam energy is preferred.

# Variations: $Q^2$



A low photon virtuality is favoured.

# Variations: momentum transfer to the proton $t$



A high  $|t|$  is favoured, but we also impose  $|t|/Q^2 \leq 0.25$ .

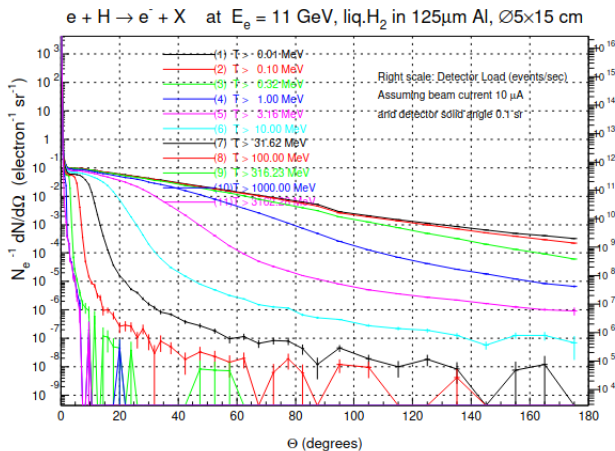
# CFF inputs

Model	$\mathcal{H}$	$\mathcal{E}$	$\tilde{\mathcal{H}}$	$\tilde{\mathcal{E}}$	$P_x^m$	$P_y^m$
GK	-1.1 +5.4i	-2.4 -0.4i	0.7 +1.8i	38 +19i	-0.33	0.41
VGG	-2.2 +4.8i	-1.0 +1.5i	0.5 +1.4i	51 +0i	0.32	0.17
KM15	-2.9 +3.2i	1.6 +0i	0.5 +1.5i	124 +0i	-0.44	-0.04
ANN	$-1.6^{+2.6}_{-3.3} + 3.3^{+3.3}_{-3.3}i$	$-4.1^{+9.2}_{-0} + 0^{+18}_{-18}i$	$0.6^{+3.0}_{-1.1} + 1.1^{+2.5}_{-2.5}i$	$-25^{+265}_{-4} + 4^{+72}_{-72}i$	$0.04^{+0.25}_{-0.25}$	$0.20^{+0.41}_{-0.41}$

- Larger  $|t| \rightarrow$  larger  $\Delta P_y$ , smaller  $d\sigma$ .
- Larger  $Q^2 \rightarrow$  smaller  $\Delta P_y$ , smaller  $d\sigma$ .
- Larger  $x_B \rightarrow$  smaller  $d\sigma$ .
- Large  $|k'|$  (small  $Q^2/x_B$ )  $\rightarrow$  larger acceptance.
- Small  $E_k$  or large  $Q^2$  - the angle between the electron and the proton decreases.



# Accidental background



- The background of accidental electrons is largest at small angles from the beam.