Pion Parton Distributions at 22 GeV

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Studying hadron structure

Should we only study the structure of protons to learn about QCD? *NO!*

- Mesons offer importance of emergent phenomena of QCD such as
 - How is mass generated?
 - How do quarks and gluons arrange themselves within hadrons?
 - Why is there confinement?
- Allows for another probe of confinement scales in quark-gluon bound systems

Questions in meson structure

- 1. How to test universality of PDFs?
- 2. What is the gluonic content of pions compared with protons?
- 3. What is the transverse momentum structure of pions, and how does it relate to confinement scales?

Accessing pions indirectly

- Exchange of pions among nucleons keep the nucleus intact
- Model the nucleon as having a "cloud of pions"



Leading Neutron (LN) electroproduction



Available datasets for pion structures



Pion PDFs in JAM



1. How to test universality of pion PDFs?

- The tagged deep inelastic scattering (TDIS) program at JLab offers overlapping kinematics with DY
- But how reliable are these kinematics with 11 GeV? resonance region



Meson	Mass~(MeV)	Decay width (MeV)
ρ	775.26 ± 0.23	149.1 ± 0.8
b_1	1229.5 ± 3.2	142 ± 9
a_1	1230 ± 40	425 ± 175
a_2	1318.2 ± 0.6	107 ± 5
$\pi(1670)$	$1670.6\substack{+2.9\\-1.2}$	258^{+8}_{-9}

Current 11 GeV TDIS kinematics

• Plotting available 11 GeV TDIS kinematics with a few representative W_{π} curves



Kinematics with 22 GeV

• We keep some data points above the $W_{\pi} = 2$ GeV cut



Impact on pion PDFs with 22 GeV

- Knowledge of pion PDFs increases dramatically with 22 GeV beam
- Assuming 1.2% systematic uncertainty and 200 days of data-taking



2. Gluonic content of the pion

- The gluon has sensitivity to F_2^{π} at next-to-leading order (NLO)
- However, it comes in at leading order (LO) for F_L^{π}
- If we can perform L-T separation in regions of kinematics, we may be able to access g_π
- Because the ρ meson does not contribute to F_L , we analyze the region in $2m_\pi < W_\pi < 1~{\rm GeV}$

Impact of F_L studies

- We look only at 11 GeV kinematics that overlap with 8.8 GeV beam kinematics
- Reduction in the gluon uncertainty at large x



3. Transverse momentum structure

Relation to k_T -space TMD

 $ilde{f}_{q/\mathcal{N}}(x,b_T) = (2\pi)^2 \int d^2 oldsymbol{k}_T e^{-ioldsymbol{b}_T\cdotoldsymbol{k}_T} f_{q/\mathcal{N}}(x,k_T)$

$$ilde{f}_{q/\mathcal{N}}(b_T|x;Q,Q^2) \equiv rac{ ilde{f}_{q/\mathcal{N}}(x,b_T;Q,Q^2)}{\int \mathrm{d}^2 oldsymbol{b}_T ilde{f}_{q/\mathcal{N}}(x,b_T;Q,Q^2)} \cdot \cdot$$

- Broadening in b_T -space appearing as x increases \Rightarrow Narrowing in k_T -space
- Up quark in pion is narrower than up quark in proton in b_T -space \Rightarrow Broader in k_T -space



Pion SIDIS: access to TMDs

 $eN \rightarrow e'N'\pi X$

- Measure an outgoing pion in the TDIS experiment
- Gives us another observable sensitive to pion TMDs
 - Needed for tests of universality



Available kinematics for JLab

Can only use 22 GeV data for any TMD analysis



Conclusion

- Impacts from the 11 GeV TDIS experiment on pion PDFs will be limited
- Tests of universality in "clean" DIS regions are needed at 22 GeV
- The 11 GeV TDIS can map out the low- W_{π} resonance region and may allow for F_L constraints
- SIDIS at 22 GeV can offer another observable for pion TMDs

Backup Slides

EIC vs JLab 22 GeV

 JLab measurements will be much more precise with a 200 day beam run – luminosity plays a big role



Testing systematics of the Sullivan process

• We look to $pp \rightarrow nX$ data as well



- Here, sensitive as well to the $f_{\pi N}$ splitting function
 - Additional observable to test the universality

Data and theory comparisons

• Perform cut on $|t| < 0.1 \text{ GeV}^2$



Resulting splitting function

• Agrees with the prior within the uncertainty bands



Resulting χ^2 for the $pp \rightarrow nX$ data

- All models as a function of the cut on |t|
- |t|_{max} = 0.1 GeV² is ideal as it gives good description of data for all models



JAM analysis with threshold resummation



Pion PDFs from lattice + experimental data



• The inclusion of lattice QCD data along with experimental data can also help us to reveal pion structure

Check the resonance regions



$$I^{G}(J^{P}) = 1^{-}(0^{-})$$

Mass $m = 139.57039 \pm 0.00018$ MeV (S = 1.8) Mean life $\tau = (2.6033 \pm 0.0005) \times 10^{-8}$ s (S = 1.2) $c\tau = 7.8045$ m



$$I(J^{PC}) = 0.1(1^{-1})$$

 $\begin{array}{ll} \text{Mass } m < \ 1 \times 10^{-18} \ \text{eV} \\ \text{Charge } q < \ 1 \times 10^{-46} \ e & (\text{mixed charge}) \\ \text{Charge } q < \ 1 \times 10^{-35} \ e & (\text{single charge}) \\ \text{Mean life } \tau = \text{Stable} \end{array}$

Full width Γ Γ _{ee} = 7.04 ±	$=$ 149.1 \pm 0. = 0.06 keV).8 MeV
<i>b</i> ₁ (1235)	I G	$G(J^{PC}) = 1^+(1^{+})$
Mass $m = 122^{\circ}$ Full width $\Gamma =$	9.5 ± 3.2 MeV 142 ± 9 MeV	(S = 1.6) (S = 1.2)
a2(1	320)	$I^{G}(J^{PC}) = 1^{-}(2^{+})$

ρ(770)

See the note in $\rho(770)$ Particle Listings.

Mass $m = 775.26 \pm 0.25$ MeV

Mass $m = 1316.9 \pm 0.9$ MeV (S = 1.9) Full width $\Gamma = 107 \pm 5$ MeV ^[j]

 $I^{G}(J^{PC}) = 1^{+}(1^{-})$

The quantum numbers of a charged π and photon result in specific outgoing mesons

Resonances

• Possible low-lying resonances from $\gamma^*\pi$



Kinematics of Sullivan-variables

• Cuts on the $pp \rightarrow nX$ data at $|t| > 0.1 \text{ GeV}^2$



Definition of W_{π}^2

• Derived from kinematics

$$W_{\pi}^{2} = t - Q^{2} \left(1 - \frac{\bar{x}_{L}}{x} \right) = t - Q^{2} \left(1 - \frac{1}{x_{\pi}} \right).$$

Impact study details

• We created pseudodata in the form of

$$R^{\rm T} = \frac{\mathrm{d}^4 \sigma(eN \to e'N'(\Lambda)X)}{\mathrm{d}x \mathrm{d}Q^2 \mathrm{d}x_L \mathrm{d}t} / \frac{\mathrm{d}^2 \sigma(eN \to e'X)}{\mathrm{d}x \mathrm{d}Q^2} \Delta x_L \Delta t$$

• We used a luminosity of: $d\mathcal{L}/dt - 5 \times 10^{38}/\text{cm}^2/\text{s}$

Use of
$$W^2$$
 for SIDIS

The unobserved invariant mass-squared in inclusive DIS is

$$W_{\rm tot}^2 = M^2 + \frac{Q^2(1 - x_{\rm Bj})}{x_{\rm Bj}}.$$
 (6.26)

In SIDIS it is

$$W_{\text{SIDIS}}^{2} = M^{2} + M_{\text{B}}^{2} + \frac{Q^{2}(1 - x_{\text{Bj}} - z_{\text{h}})}{x_{\text{Bj}}} + \frac{Q^{4}z_{\text{h}}\left(\sqrt{1 + \frac{4M^{2}x_{\text{Bj}}^{2}}{Q^{2}}}\sqrt{1 - \frac{4M^{2}x_{\text{Bj}}^{2}M_{\text{B},\text{T}}}{z_{\text{h}}^{2}Q^{4}}} - 1\right)}{2M^{2}x_{\text{Bj}}^{2}}$$
$$\stackrel{M,M_{\text{B}} \to 0}{=} \frac{Q^{2}(1 - x_{\text{Bj}})(1 - z_{\text{h}})}{x_{\text{Bj}}} - \frac{\mathbf{P}_{\text{B},\text{T}}^{2}}{z_{\text{h}}}.$$
(6.27)

• Replace M^2 with t

Average
$$b_T$$

• The conditional expectation value of b_T for a given x

$$\langle b_T | x \rangle_{q/\mathcal{N}} = \int \mathrm{d}^2 \boldsymbol{b}_T \, b_T \, \tilde{f}_{q/\mathcal{N}}(b_T | x; Q, Q^2)$$

 Shows a measure of the transverse correlation in coordinate space of the quark in a hadron for a given x



Brief words on kaon TDIS

- Sullivan process applies, but a hyperon must be tagged
- Consider again, not only inclusive W^2 but W_K^2



Kinematics for 11 GeV Kaon TDIS

• Beware of such large |t| further away from kaon pole



Kinematics for 22 GeV Kaon TDIS

Accepting of more points at smaller |k|



Resonance from K^*

• The K^* resonance is much more narrow than for ρ meson

•
$$W_{K,\max}^2 = 1 \text{ GeV}^2$$

