



Phenomenology of helicity TMD: status and challenges

Matteo Cerutti

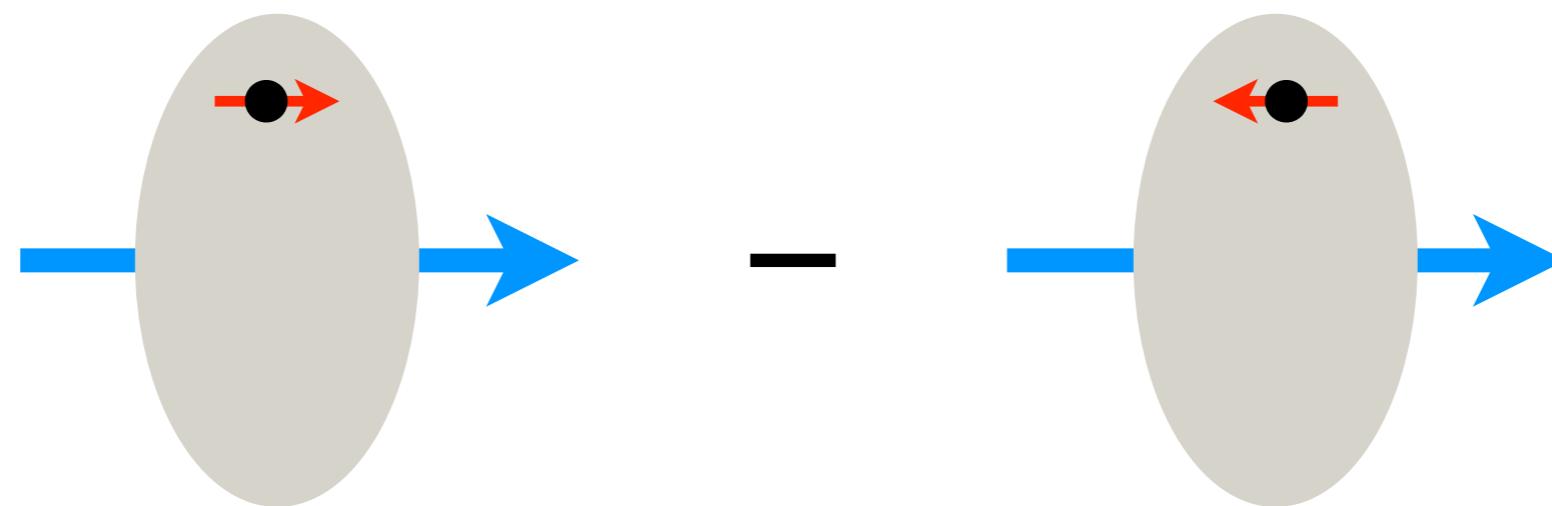
MAP Collaboration

A. Bacchetta, A. Bongallino, M. Radici, L. Rossi

arXiv: 2409.18078

Helicity distribution

$$g_1^q(x) = q^+ - q^-$$

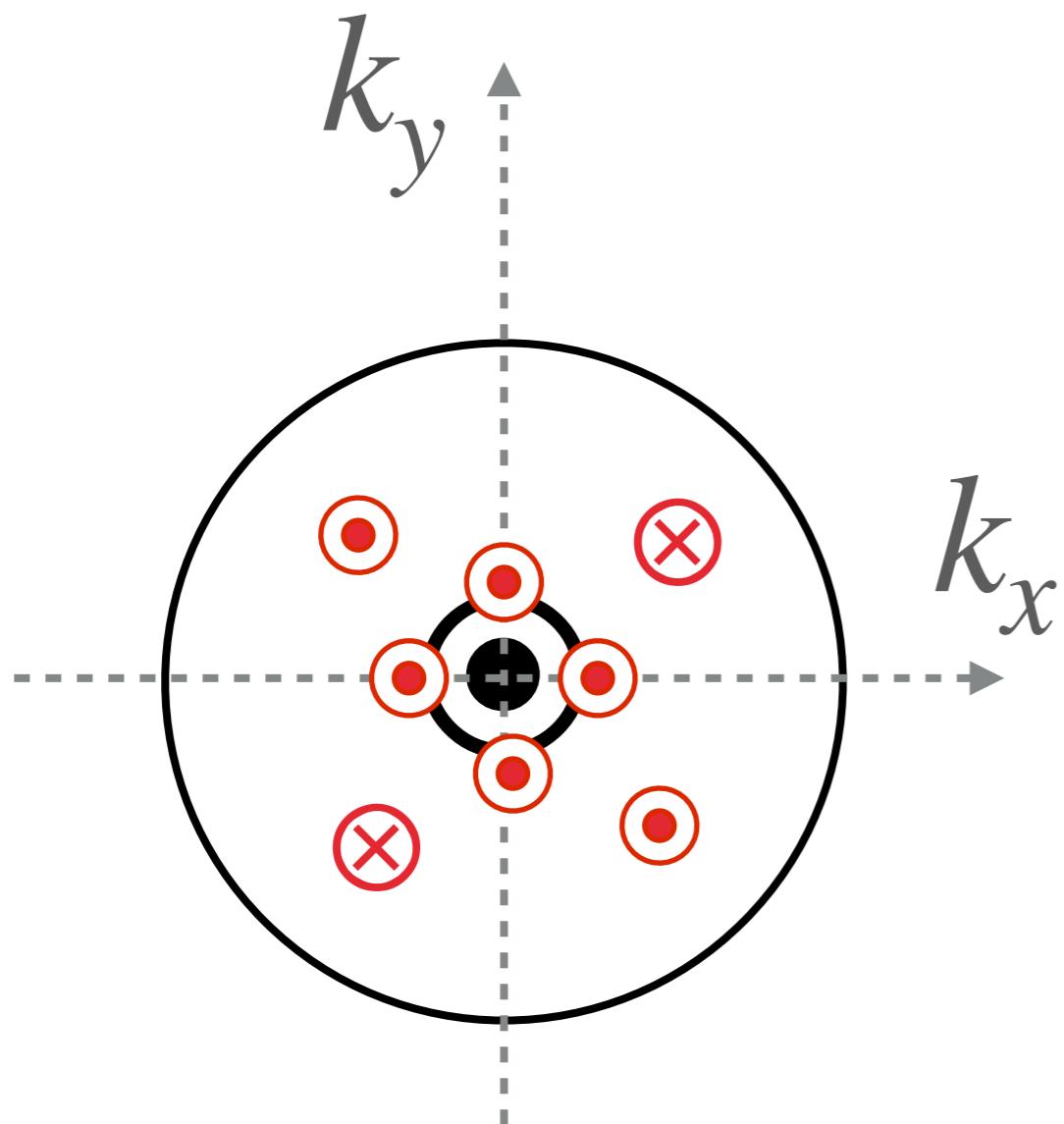


Quark Polarization

	U	L	T
U	$f_1(x)$		
L		$g_1(x)$	
T			$h_1(x)$

Helicity TMD distribution

$$g_1^q(x, k_\perp) = q^+ - q^-$$

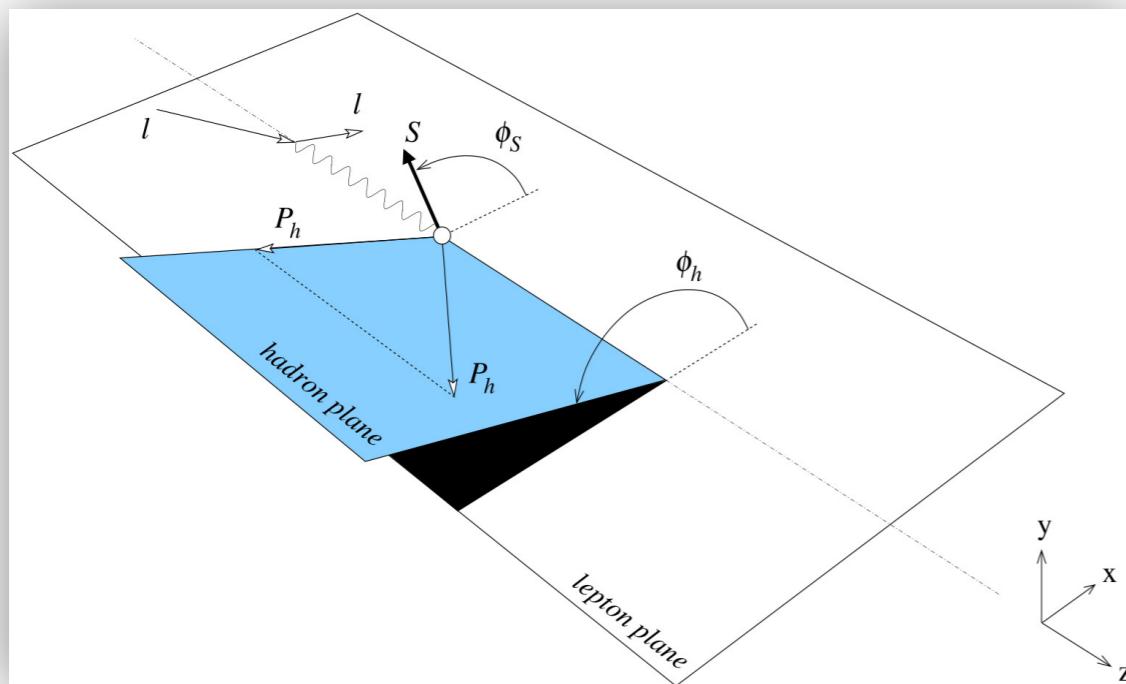


- ♦ How the polarization of the proton reflects on its internal structure in **3 dimensions?**
- ♦ How the polarization of the quark distorts their **transverse momentum?**
- ♦ Do quarks with spin parallel to the proton's spin have **smaller** or **larger** transverse momentum?

Experimental Observables

Analysis of longitudinally polarized process

SIDIS



DOUBLE SPIN ASYMMETRY

$$A_1 = \frac{d\sigma^{\rightarrow\leftarrow} - d\sigma^{\rightarrow\rightarrow} + d\sigma^{\leftarrow\rightarrow} - d\sigma^{\leftarrow\leftarrow}}{d\sigma^{\rightarrow\leftarrow} + d\sigma^{\rightarrow\rightarrow} + d\sigma^{\leftarrow\rightarrow} + d\sigma^{\leftarrow\leftarrow}}$$

A. Bacchetta et al., Phys.Rev.D 70 (2004), 117504

M. Diehl and S. Sapeta, Eur. Phys. J. C 41, 515 (2005)

Interpretation in terms of TMDs

TMD factorization

$$A_1(x, z, Q, |\mathbf{P}_{hT}|) = \frac{\sum_{a=q,\bar{q}} e_a^2 \int_0^{+\infty} d|\mathbf{b}_T|^2 J_0\left(\frac{|\mathbf{b}_T| |\mathbf{P}_{hT}|}{z}\right) \hat{g}_1^a(x, |\mathbf{b}_T|^2, Q) \hat{D}_1^{a \rightarrow h}(z, |\mathbf{b}_T|^2, Q)}{\sum_{a=q,\bar{q}} e_a^2 \int_0^{+\infty} d|\mathbf{b}_T|^2 J_0\left(\frac{|\mathbf{b}_T| |\mathbf{P}_{hT}|}{z}\right) \hat{f}_1^a(x, |\mathbf{b}_T|^2, Q) \hat{D}_1^{a \rightarrow h}(z, |\mathbf{b}_T|^2, Q)}$$

- ♦ Large energy scale $Q^2 \gg M^2$
 - ♦ Small transverse momentum $q_T^2 \ll Q^2$
- ⇒ Experimental observables in terms of universal objects

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unpolarized TMDs from MAP22 extraction

MAP Collaboration, Bacchetta et al., JHEP 10 (2022)

Parameterization

Collins, Soper, Sterman, Nucl. Phys. B (1985)

Collins, *Foundations of perturbative QCD* (2011)

The evolution of the TMDs follows the CSS approach **consistently**:

$$\hat{f}_1(x, |\mathbf{b}_T|^2, Q) = [C^f \otimes f_1](x, b_\star(|\mathbf{b}_T|^2)) f_{NP}(x, |\mathbf{b}_T|^2, Q_0) e^{S(\mu_{b_\star}^2, Q^2)} e^{g_K(\mathbf{b}_T) \ln(Q^2/Q_0^2)}$$

Analogously for $D_1(z, |\mathbf{b}_T|^2, Q)$

$$\hat{g}_1(x, |\mathbf{b}_T|^2, Q) = [C^g \otimes g_1](x, b_\star(|\mathbf{b}_T|^2)) g_{NP}(x, |\mathbf{b}_T|^2, Q_0) e^{S(\mu_{b_\star}^2, Q^2)} e^{g_K(\mathbf{b}_T) \ln(Q^2/Q_0^2)}$$

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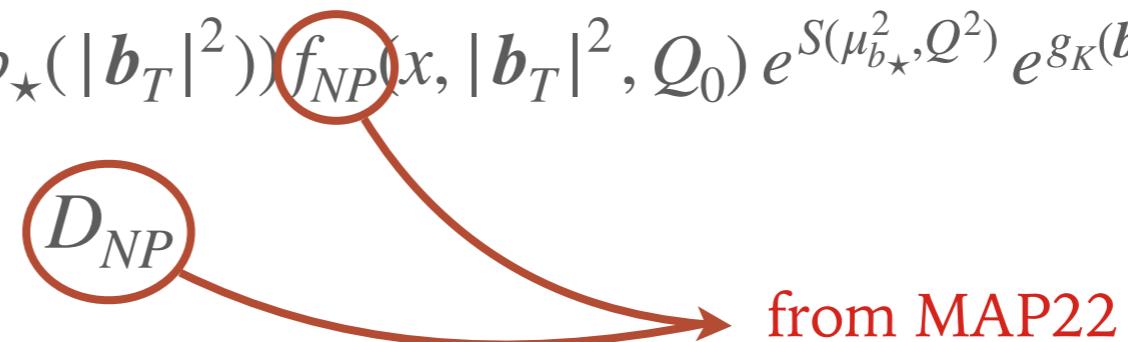
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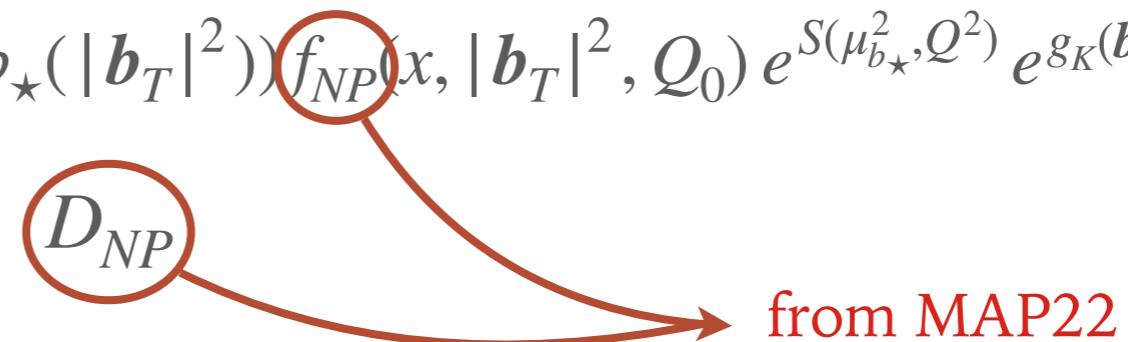
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$$g_{NP}(x, k_\perp^2, Q_0) = f_{NP}^{MAP22}(x, k_\perp^2, Q_0) \frac{e^{-\frac{k_\perp^2}{\omega_1(x)}}}{k_{norm}(x)}$$

Simple Gaussian \times MAP22

Positivity constraint

$$g_{NP}(x, k_\perp^2, Q_0) = f_{NP}^{MAP22}(x, k_\perp^2, Q_0) \frac{e^{-\frac{k_\perp^2}{\omega_1(x)}}}{k_{norm}(x)}$$

- Proportional to f_{NP}^{MAP22}
- x-dependent

$$k_{\text{norm}}(x) \rightarrow \int d^2 k_\perp g_{NP} = 1$$

$$\omega_1(x) \rightarrow \text{crucial to satisfy } |g_1| \leq f_1$$

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At $Q_0 = 1$ GeV, the ratio g_1/f_1 reads:

$$\frac{g_1(x, \mathbf{k}_\perp^2, Q_0)}{f_1(x, \mathbf{k}_\perp^2, Q_0)} = \frac{g_1(x, Q_0)}{f_1(x, Q_0)} \frac{e^{-\frac{\mathbf{k}_\perp^2}{\omega_1(x)}}}{k_{norm}(x)}$$

$\underset{|\mathbf{k}_\perp| \rightarrow 0}{\circlearrowleft} \infty$

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$$\frac{g_1(x, Q_0)}{f_1(x, Q_0)} \frac{1}{k_{norm}(x)} \leq 1 \quad \longrightarrow$$

$$\omega_1(x) = f_{pos.}(x) + N_{1g}^2 \frac{(1-x)^{\alpha_{1g}^2} x^{\sigma_{1g}}}{(1-\hat{x})^{\alpha_{1g}^2} \hat{x}^{\sigma_{1g}}}$$

Phenomenology: STATUS

Airapetian et al. (HERMES), Phys. Rev. D (2019)

Experiment	N_{dat}	$\chi^2_{\text{NLL}}/N_{\text{dat}}$	$\chi^2_{\text{NNLL}}/N_{\text{dat}}$
HERMES ($d \rightarrow \pi^+$)	47	1.34	1.30
HERMES ($d \rightarrow \pi^-$)	47	1.10	1.08
HERMES ($d \rightarrow K^+$)	46	1.26	1.25
HERMES ($d \rightarrow K^-$)	45	0.93	0.89
HERMES ($p \rightarrow \pi^+$)	53	1.17	1.21
HERMES ($p \rightarrow \pi^-$)	53	0.86	0.86
Total	291	1.11	1.09

- ◆ MAP22 kinematic cuts
- ◆ 291 fitted data points
- ◆ Perturbative order: NLO

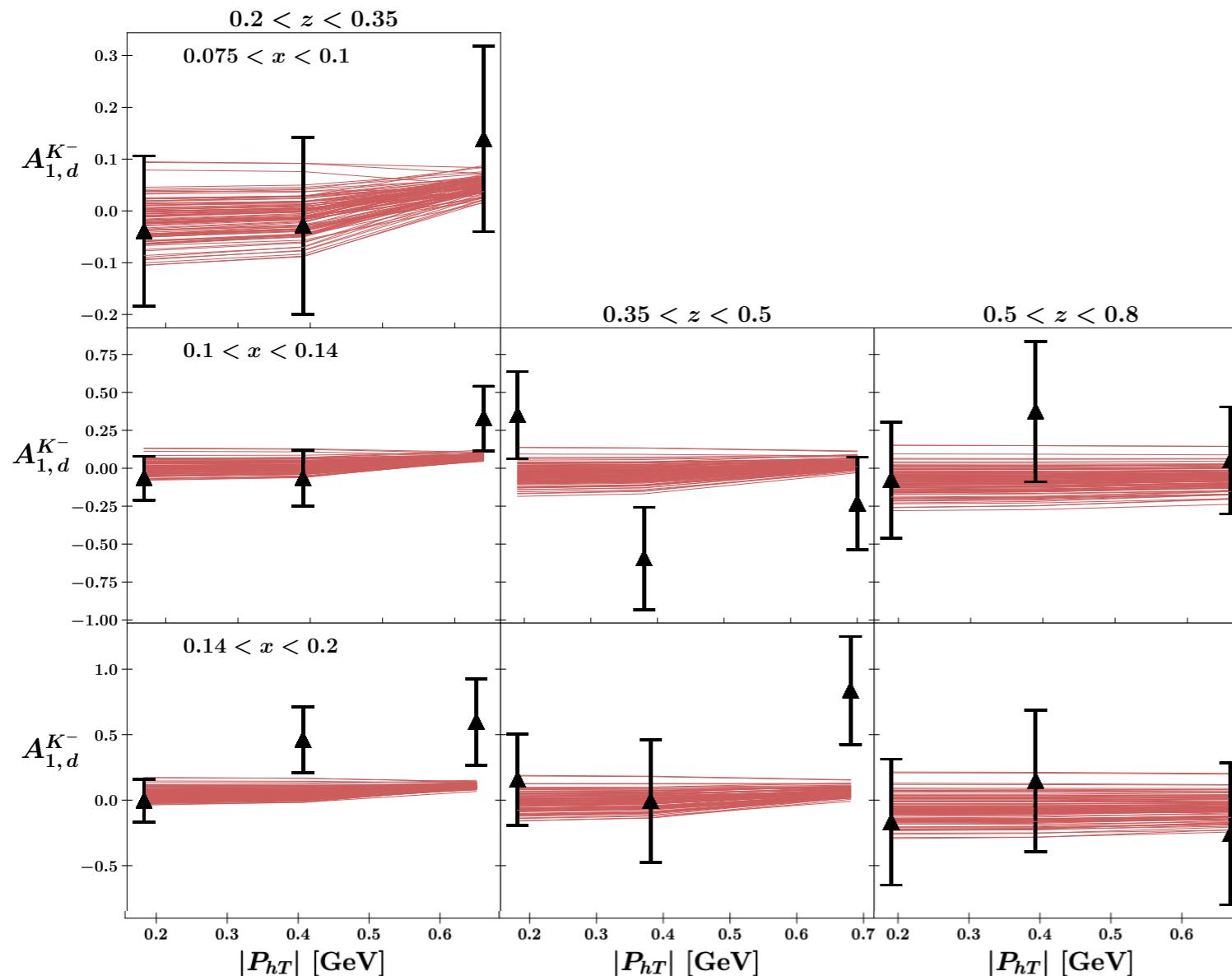
Highest possible
since C^g known up to NLO

- ◆ Collinear PDFs: NNPDFPol, MMHT, DSS
- ◆ Perturbative accuracy: **NLL & N2LL**
- ◆ 3 fitted parameters
- ◆ Error analysis with bootstrap method

Gutiérrez-Reyes et al., Phys. Lett. B (2017)

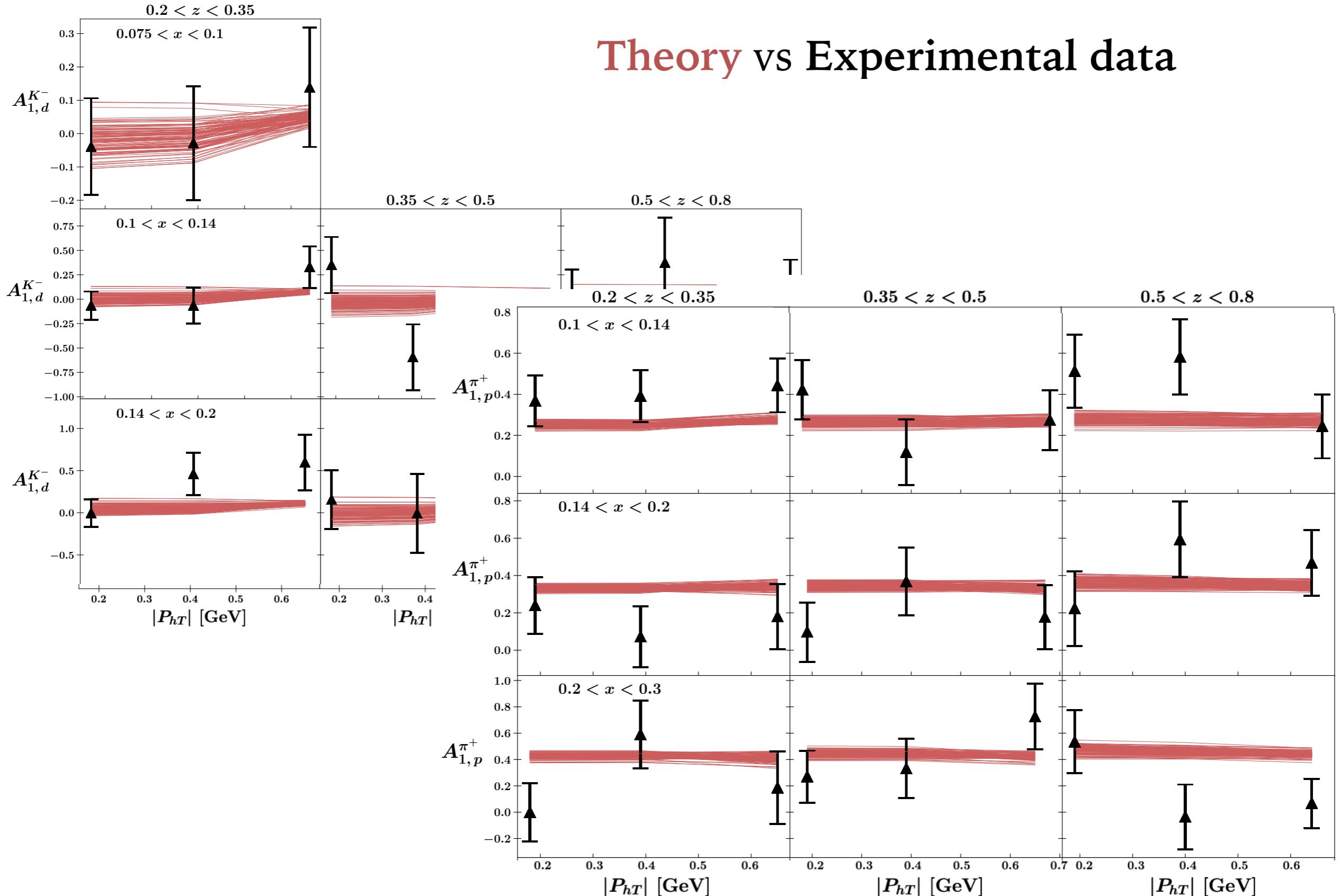
Phenomenology: STATUS

Experimental data

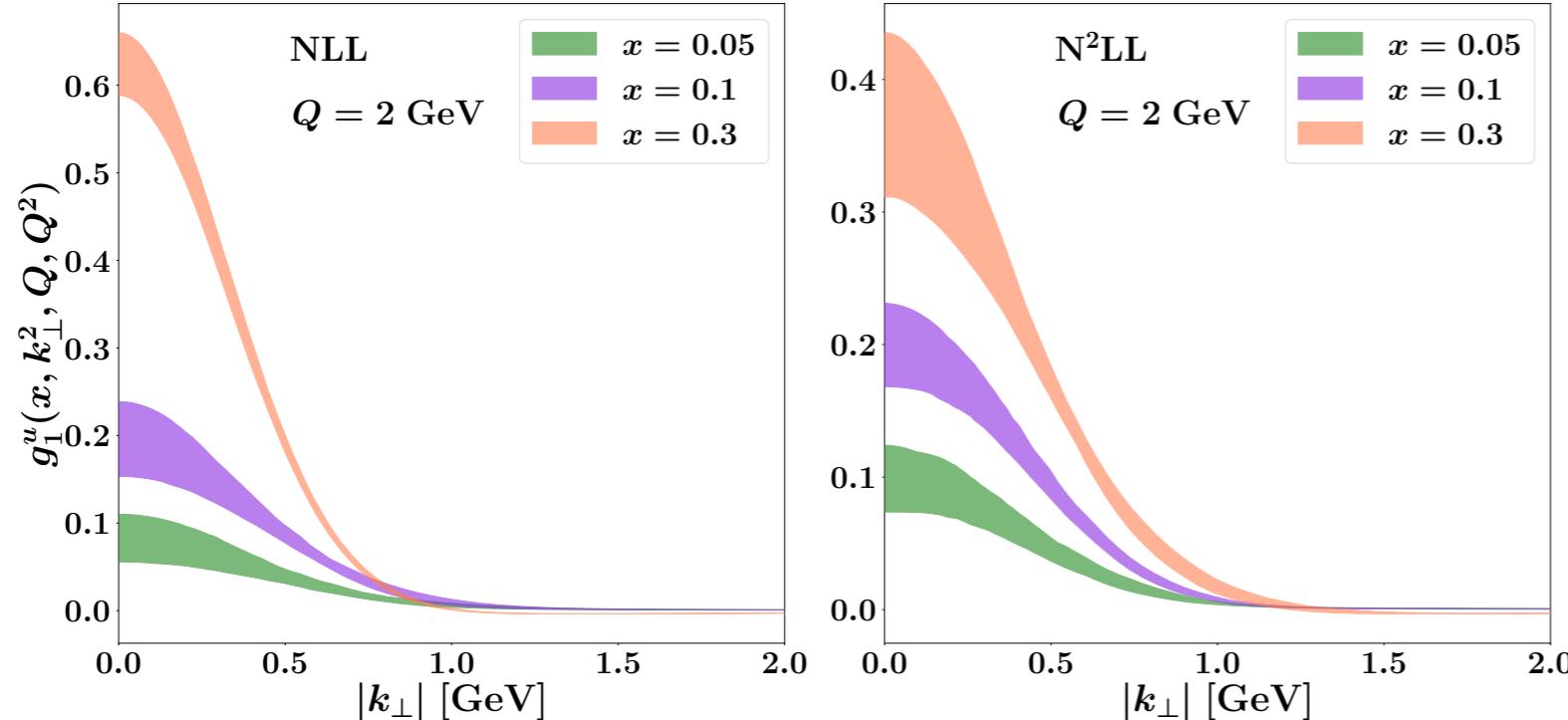


Phenomenology: STATUS

Theory vs Experimental data

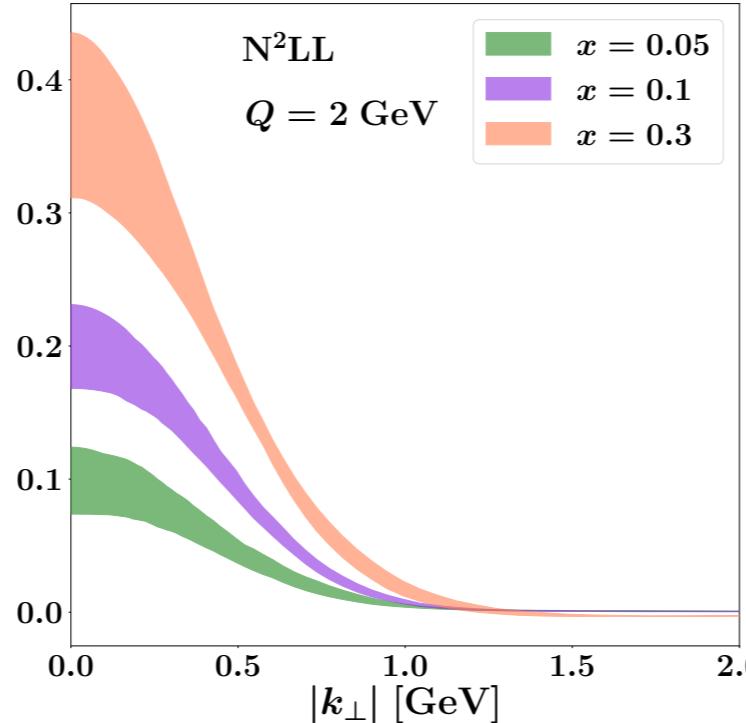
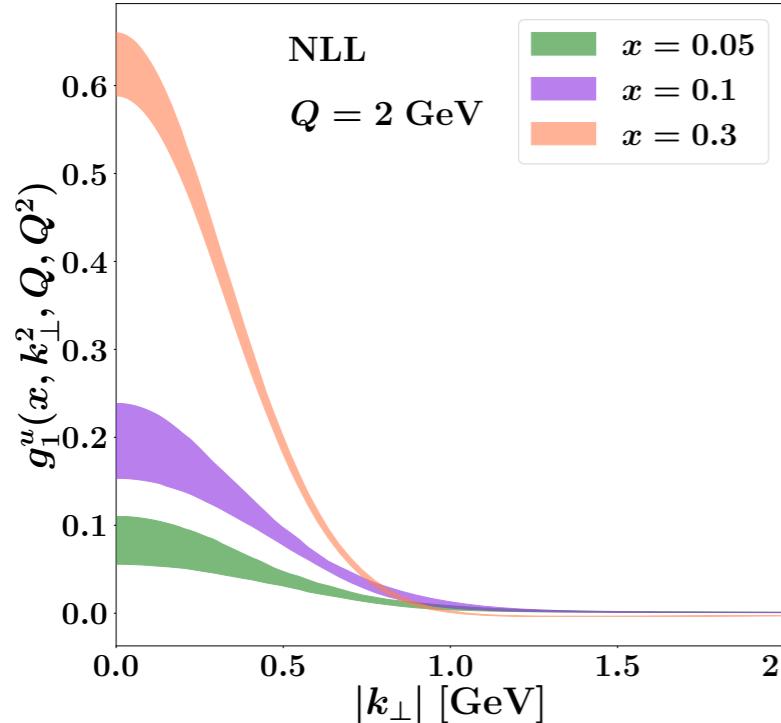


Phenomenology: STATUS



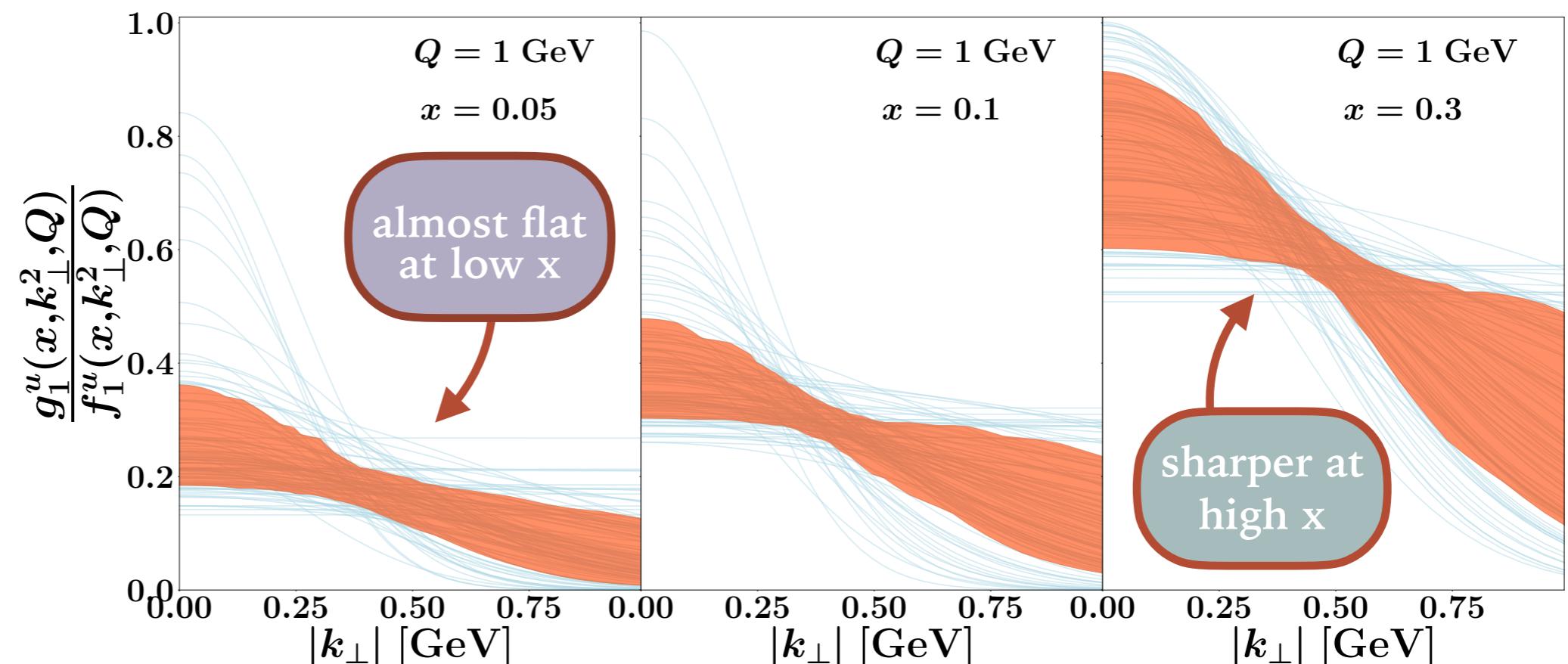
Extracted helicity
TMDs

Phenomenology: STATUS

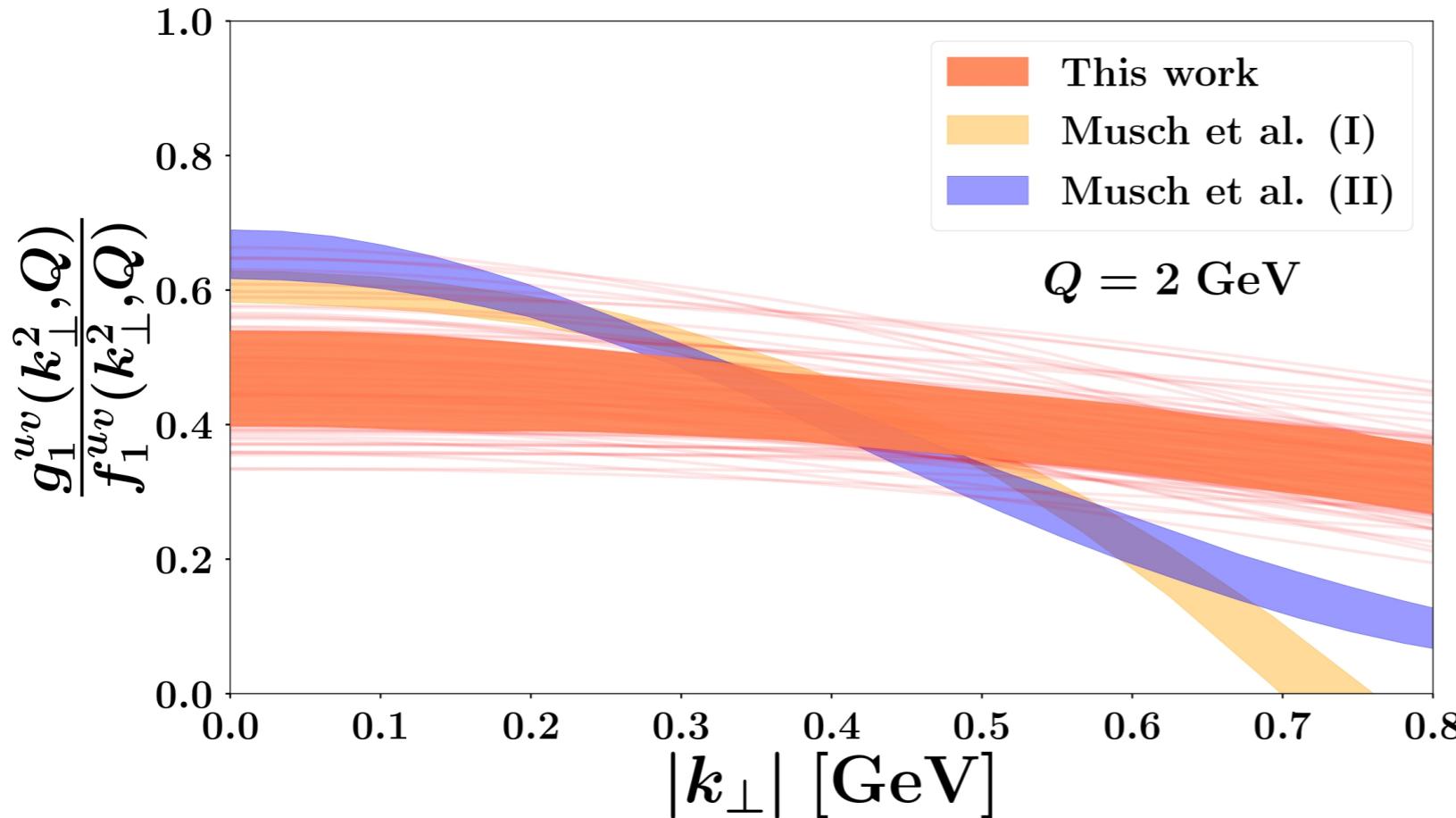


Extracted helicity
TMDs

Helicity
ratio



Phenomenology: STATUS



Comparison to
lattice calculation

Musch et al., Phys. Rev. D (2011)

- ♦ g_1/f_1 TMD ratio for u_v integrated over x , at NNLL
- ♦ Yellow and blue bands correspond to two lattice predictions
- ♦ Milder slope but fair agreement

Phenomenology: CHALLENGES

- Not many constraints on fitted parameters

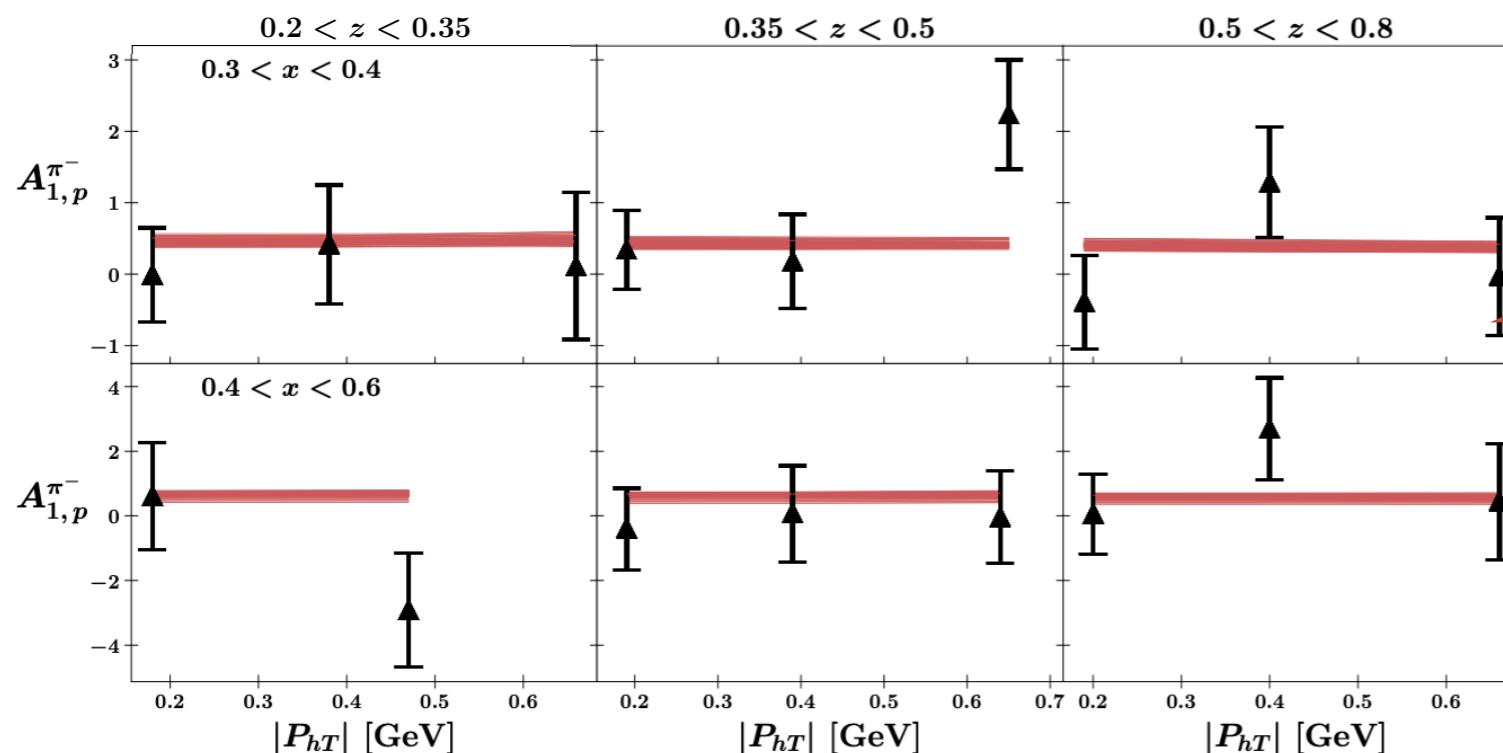
Parameters	N_{1g}	α_{1g}	σ_{1g}
NLL	0.70 ± 0.54	27.81 ± 27.70	0.42 ± 0.86
NNLL	0.87 ± 0.72	6.73 ± 6.58	3.04 ± 3.09

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- Experimental asymmetries ≥ 1 in several bins



What does it happen at large x ?

Need of **precise** data from JLab 22 GeV upgrade

Backup

MAP TMD fitting framework

<https://github.com/MapCollaboration/NangaParbat>



☰ README.md



Nanga Parbat is a fitting framework aimed at the determination of the non-perturbative component of TMD distributions.

Download

You can obtain NangaParbat directly from the github repository:

<https://github.com/MapCollaboration/NangaParbat>

For the last development branch you can clone the master code:

```
git clone git@github.com:MapCollaboration/NangaParbat.git
```

MAPTMD22 global fit: features

- Global analysis of Drell-Yan and SIDIS data sets: **2031** data points
- Perturbative accuracy: $N^3 LL^-$
- ***Normalization prefactor*** for SIDIS observables
- Number of fitted parameters: **21**
- Agreement with data: $\chi^2/N_{data} = 1.06$

Structure of a TMD

TMD in Fourier space

$$\hat{F}(x, b_T^2; \mu, \zeta) = \int \frac{d^2 \mathbf{k}_\perp}{(2\pi)^2} e^{i \mathbf{b}_T \cdot \mathbf{k}_\perp} F(x, k_\perp^2; \mu, \zeta)$$

Collinear extractions
b_{*}-prescription

$$\hat{f}_1^q(x, b_T^2; \mu, \zeta) = \sum_j C_{q/j}(x, b_*; \mu_{b_*}, \mu_{b_*}^2) \otimes f_1^j(x, \mu_{b_*})$$

: A

Perturbative TMD at the initial scale

$$\times \exp \left\{ K(b_*; \mu_{b_*}) \ln \frac{\sqrt{\zeta}}{\mu_{b_*}} + \int_{\mu_{b_*}}^\mu \frac{d\mu'}{\mu'} \left[\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta}}{\mu'} \right] \right\}$$

: B

Evolution to final scale (of the process)

$$\times f_{NP}(x, b_T^2) \exp \left\{ g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}} \right\}$$

: C

Non-perturbative part of the TMD

Parameterization

MAP22: Perturbative accuracy

Resummation of large logs

$$S_{\text{pert}}(\mu_b, \mu) = 1 + \sum_{k=0}^{\infty} R_{N^k \text{LL}}$$

$$S_{\text{pert}}(\mu_b, \mu) = 1 + \sum_{k=0}^{\infty} \sum_{n=1+[k/2]}^{\infty} \left(\frac{\alpha_S(\mu)}{4\pi} \right)^n \sum_{k=1}^{2n} L^{2n-k} R^{(n, 2n-k)} \quad L = \ln \left(\frac{\mu^2}{\mu_b^2} \right)$$

Accuracy	H and C	K and γ_F	γ_K	PDF/FF and α_S evol.
LL	0	-	1	-
NLL	0	1	2	LO
NLL'	1	1	2	NLO
NNLL	1	2	3	NLO
NNLL'	2	2	3	NNLO
N³LL'	2	3	4	NNLO/NLO
N ³ LL	2	3	4	NNLO
N ³ LL'	3	3	4	N ³ LO

MAP22: NP parametrization

$$f_{NP}(x, b_T^2) \exp \left\{ g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}} \right\}$$

Bacchetta, Gamberg, Goldstein, et al., PLB 659 (2008)

Bacchetta, Conti, Radici, PRD 78 (2008)

Pasquini, Cazzaniga, Boffi, PRD 78 (2008)

Matevosyan, Bentz, Cloet, Thomas, PRD 85 (2012)

Burkardt, Pasquini, EPJA (2016)

Grewal, Kang, Qiu, Signori, PRD 101 (2020)

$$f_{1NP}(x, b_T^2) \propto \text{F.T. of } \left(e^{-\frac{k_\perp^2}{g_{1A}}} + \lambda_B k_\perp^2 e^{-\frac{k_\perp^2}{g_{1B}}} + \lambda_C e^{-\frac{k_\perp^2}{g_{1C}}} \right)$$

$$g_1(x) = N_1 \frac{(1-x)^\alpha}{(1-\hat{x})^\alpha} \frac{x^\sigma}{\hat{x}^\sigma}$$

$$D_{1NP}(x, b_T^2) \propto \text{F.T. of } \left(e^{-\frac{P_\perp^2}{g_{3A}}} + \lambda_{FB} k_\perp^2 e^{-\frac{P_\perp^2}{g_{3B}}} \right)$$

$$g_3(z) = N_3 \frac{(z^\beta + \delta)(1-z)^\gamma}{(\hat{z}^\beta + \delta)(1-\hat{z})^\gamma}$$

$$g_K(b_T^2) = -g_2^2 \frac{b_T^2}{4}$$

11 parameters for TMD PDF
 + 1 for NP evolution + 9 for TMD FF
 = 21 free parameters

MAP22: Kinematic cuts

Drell-Yan

SIDIS

Fixed-target low-energy DY

HERMES data

RHIC data

COMPASS data

LHC and Tevatron data

$Q > 1.3 \text{ GeV}$

$9 \lesssim Q \lesssim 11 \text{ GeV}$ excluded (Υ resonance)

$0.2 < z < 0.7$

$q_T|_{\max} = 0.2Q$

$P_{hT}|_{\max} = \min[\min[0.2Q, 0.5zQ] + 0.3 \text{ GeV}, zQ]$

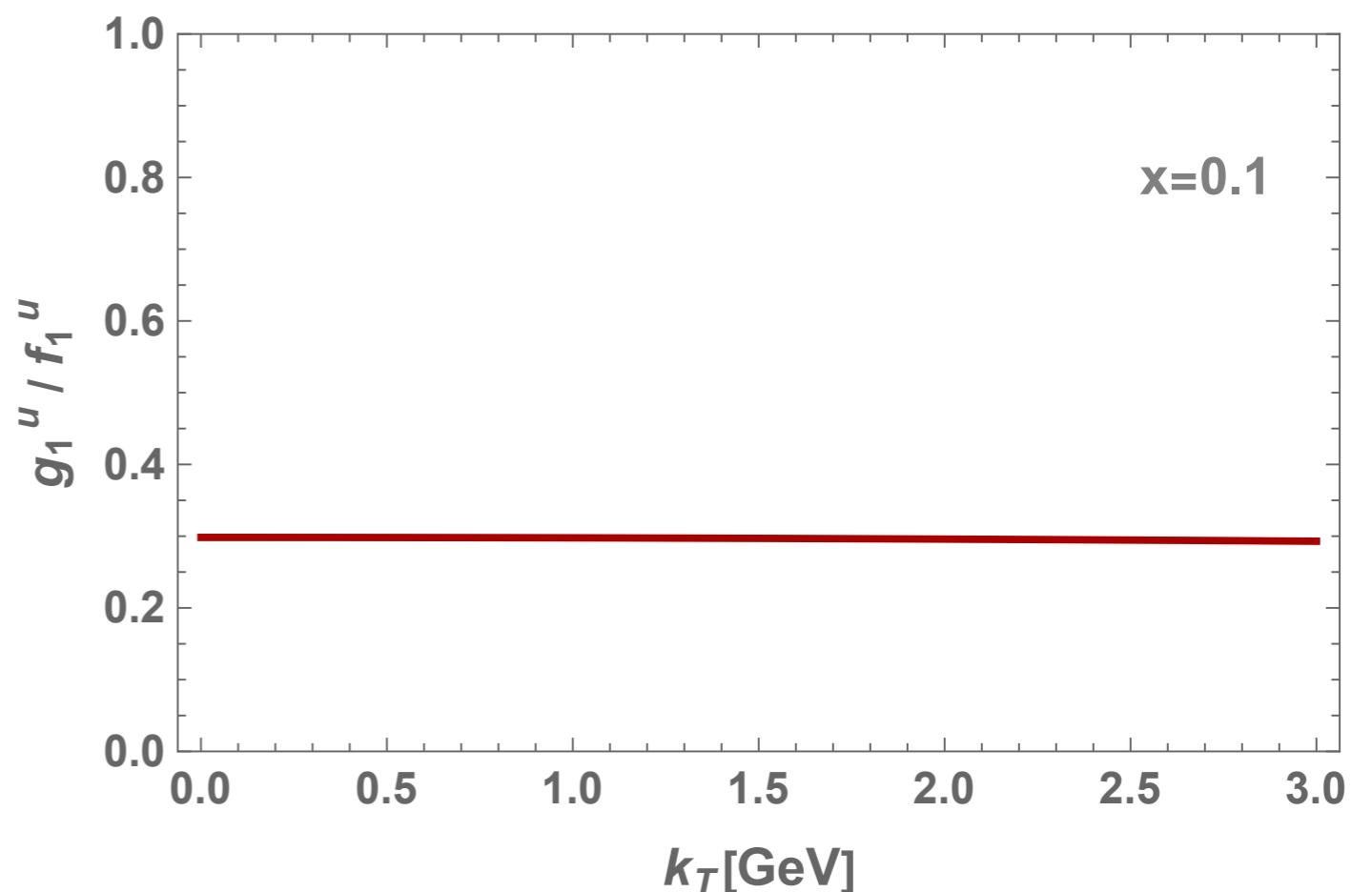
484 experimental points

1547 experimental points

Role of Gaussian

$$g_{NP}(x, k_\perp^2, Q_0) = f_{NP}^{MAP22}(x, k_\perp^2, Q_0) \frac{e^{-\frac{k_\perp^2}{\omega_1(x)}}}{k_{norm}(x)}$$

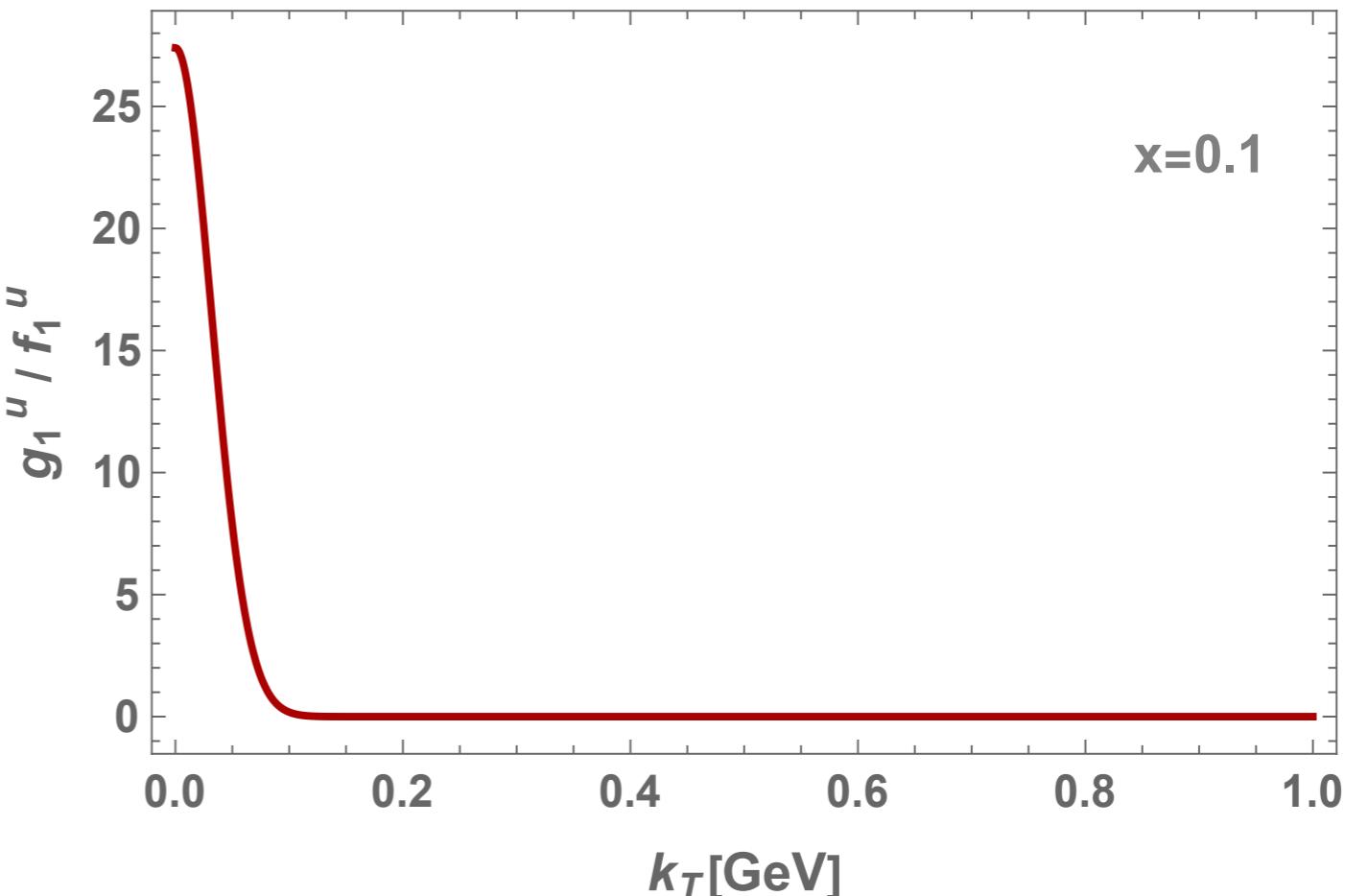
$$\omega_1(x) \rightarrow \infty \implies g_1 \simeq f_1$$



Role of Gaussian

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$\omega_1(x) \rightarrow 0 \implies$ Positivity broken
at small k_T



Positivity constraint

- ♦ $f_{pos.}(x)$ guarantees the positivity bound

at TMD level

$$10^{-4} \leq x \leq 0.7$$

$$f_{pos.}(x) \approx c + h^2 e^{-\frac{(x-\mu)^2}{\sigma^2}}$$

(interpolation function)

- ♦ $N_{1g}, \alpha_{1g}, \sigma_{1g}$ are the **free parameters** of the fit, $\hat{x} = 0.1$

$$\omega_1(x) = f_{pos.}(x) + N_{1g} \frac{(1-x)^{\alpha_{1g}^2} x^{\sigma_{1g}}}{(1-\hat{x})^{\alpha_{1g}^2} \hat{x}^{\sigma_{1g}}}$$

Propagation of errors

$f_1(x) \rightarrow$ MMHT2014 set, $D_1(z) \rightarrow$ DSS14, DSS17 sets

$g_1(x) \rightarrow$ NNPDFpol1.1: 100 MC members



100 replicas of A_1 data points to be fitted

i-th replica of $g_1(x)$ and the extracted g_1 TMD
associated with the same replica of unpolarized TMDs



Uncertainty of extracted collinear PDF propagated onto TMD's uncertainty