

Search for missing Sigma-hyperon states

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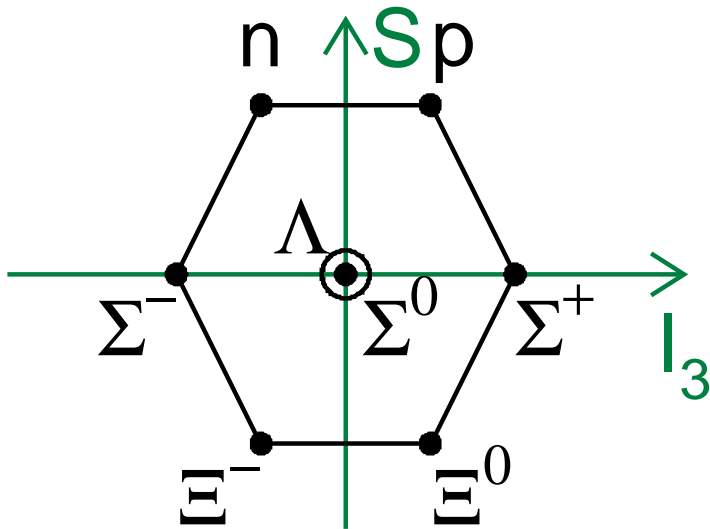


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$$3 \otimes 3 \otimes 3 = 10_S \oplus 8_M \oplus 8_M \oplus 1_A$$

Octet



Decuplet

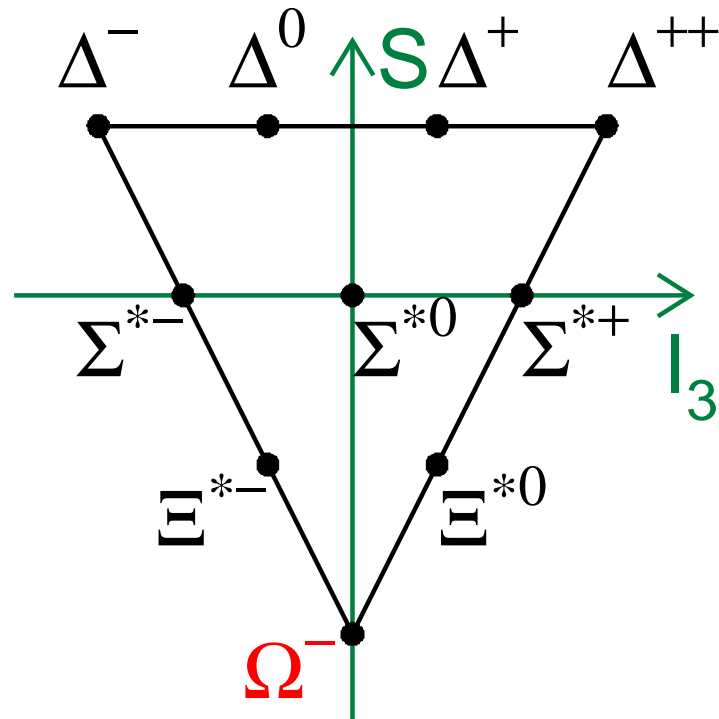


Table 1: Λ -hyperons used in the first fit of the data.

		J^P	Status	Mass	Width
singlet	$\Lambda(1405)$	$1/2^-$	****	$1405_{-1.0}^{+1.3}$	50.5 ± 2.0
$N(1535)$	$\Lambda(1670)$	$1/2^-$	****	$1660 - 1680$	$25 - 50$
$N(1650)$	$\Lambda(1800)$	$1/2^-$	***	$1720 - 1850$	$200 - 400$
singlet	$\Lambda(1520)$	$3/2^-$	****	1519.5 ± 1.0	15.6 ± 1.0
$N(1520)$	$\Lambda(1690)$	$3/2^-$	****	$1685 - 1695$	$50 - 70$
$N(1675)$	$\Lambda(1830)$	$5/2^-$	****	$1810 - 1830$	$60 - 110$
$N(2190)$	$\Lambda(2100)$	$7/2^-$	****	$2090 - 2110$	$100 - 250$
$N(1440)$	$\Lambda(1600)$	$1/2^+$	***	$1560 - 1700$	$50 - 250$
$N(1710)$	$\Lambda(1810)$	$1/2^+$	***	$1750 - 1850$	$50 - 250$
$N(1700)$	$\Lambda(1890)$	$3/2^+$	****	$1850 - 1910$	$60 - 200$
$N(1680)$	$\Lambda(1820)$	$5/2^+$	****	$1815 - 1825$	$70 - 90$
$N(2060)$	$\Lambda(2110)$	$5/2^+$	***	$2090 - 2140$	$150 - 250$

Table 2: Σ -Hyperons used in the first fit of the data.

		J^P	Status	Mass	Width
$N(1440)$	$\Sigma(1660)$	$1/2^+$	***	1630 – 1690	40 – 200
$\Delta(1230)$	$\Sigma(1385)$	$3/2^+$	****	1382.80 ± 0.35	36.0 ± 0.7
$N(1680), \Delta(1905)$	$\Sigma(1915)$	$5/2^+$	****	1900 – 1935	80 – 160
$N(1990), \Delta(1950)$	$\Sigma(2030)$	$7/2^+$	****	2025 – 2040	150 – 200
$N(1520)$	$\Sigma(1670)$	$3/2^-$	****	1665 – 1685	40 – 80
$N(1535), \Delta(1620), N(1650)$	$\Sigma(1750)$	$1/2^-$	***	1730 – 1800	60 – 160
$N(1675)$	$\Sigma(1775)$	$5/2^-$	****	1770 – 1780	105 – 135
$N(1700), \Delta(1700)$	$\Sigma(1940)$	$3/2^-$	***	1900 – 1950	150 – 300

Many Σ states are missing.

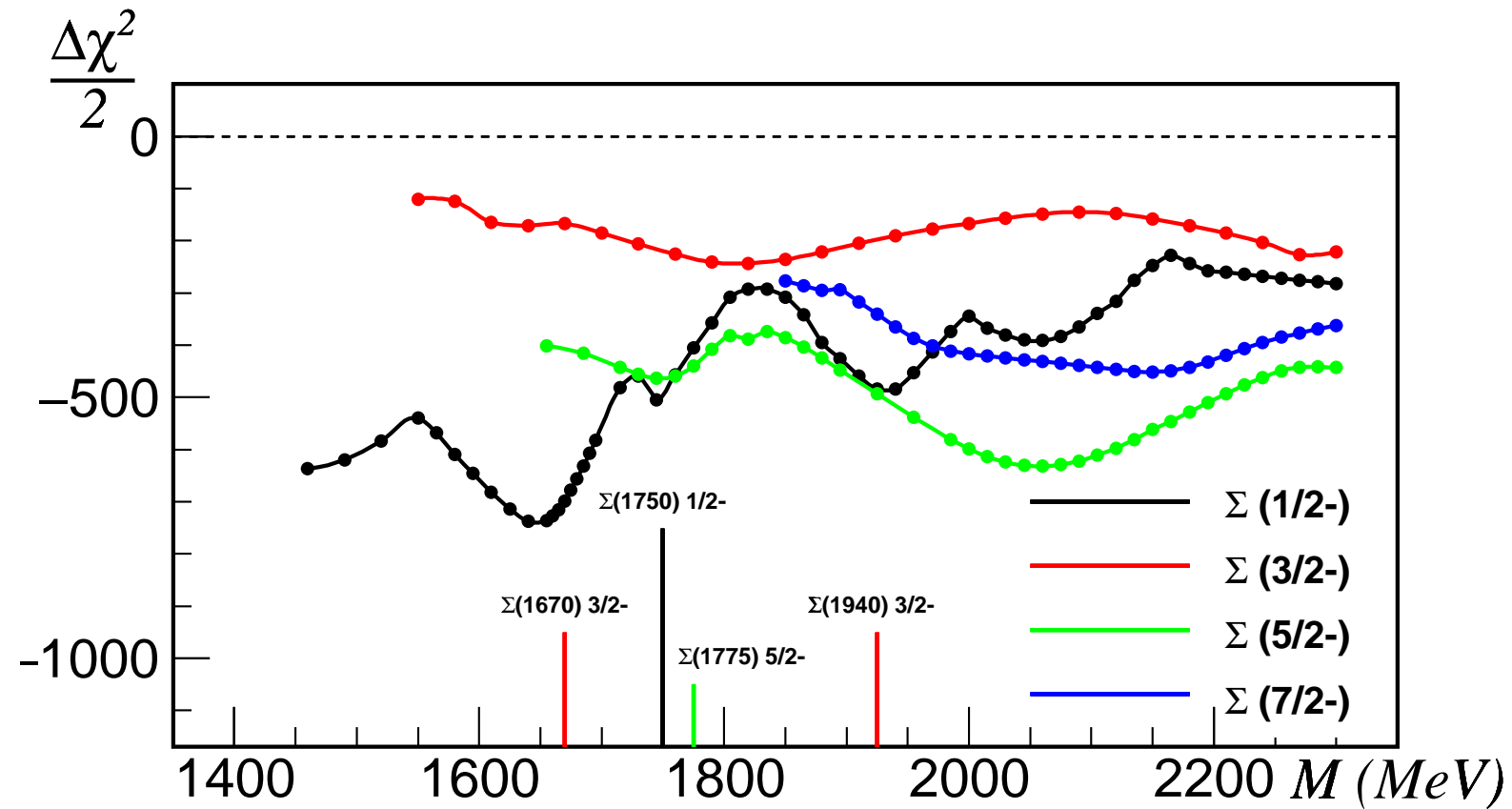
Kaon beam motivation

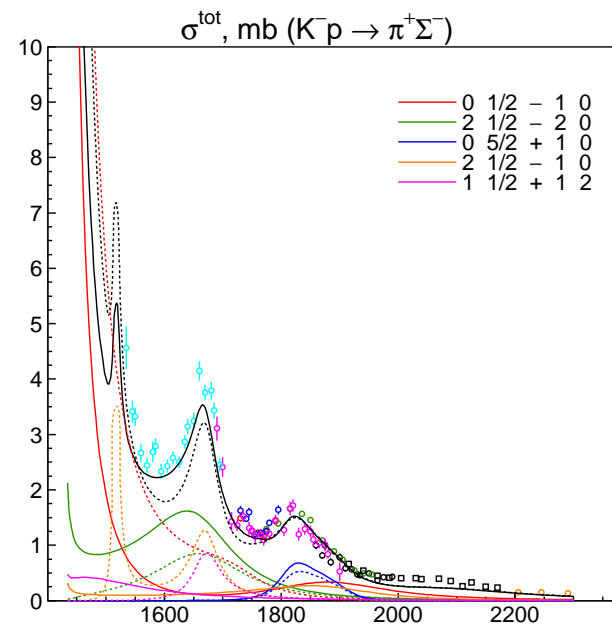
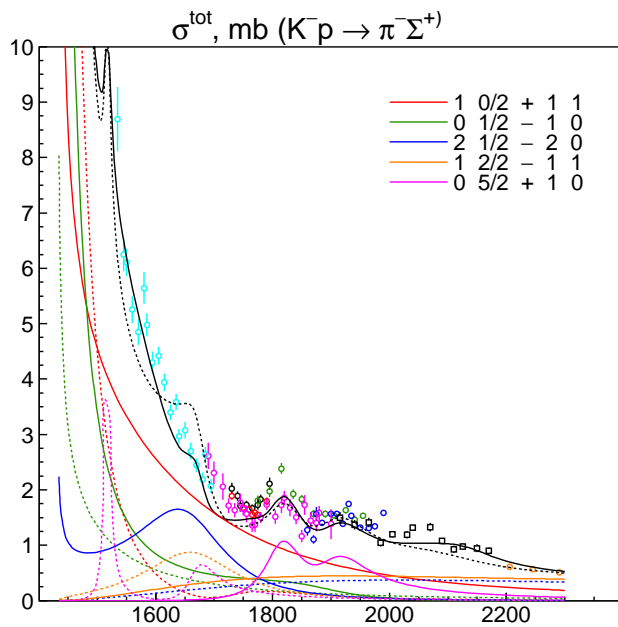
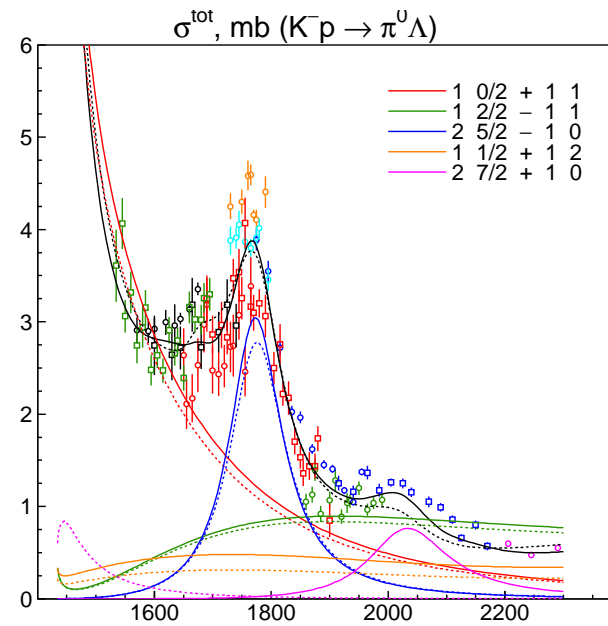
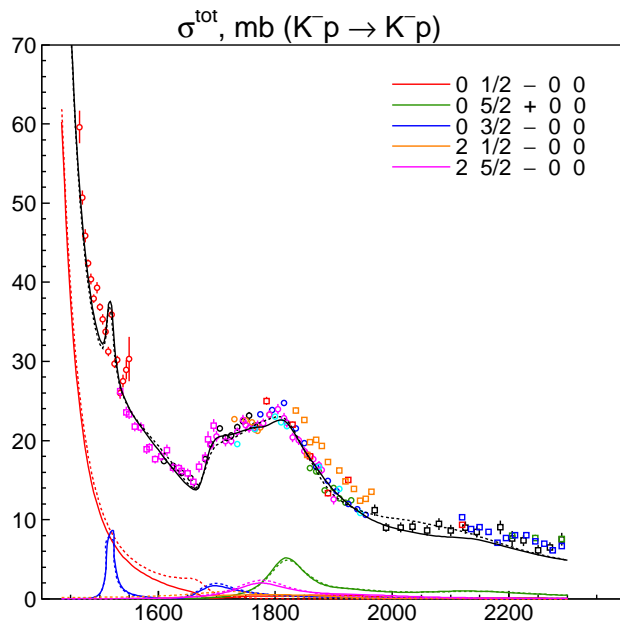
There is a hope to observe the baryon multiplets and therefore to confirm the states observed in the Nucleon and Delta sector.

Table 3: List of reactions used in the partial wave analysis.

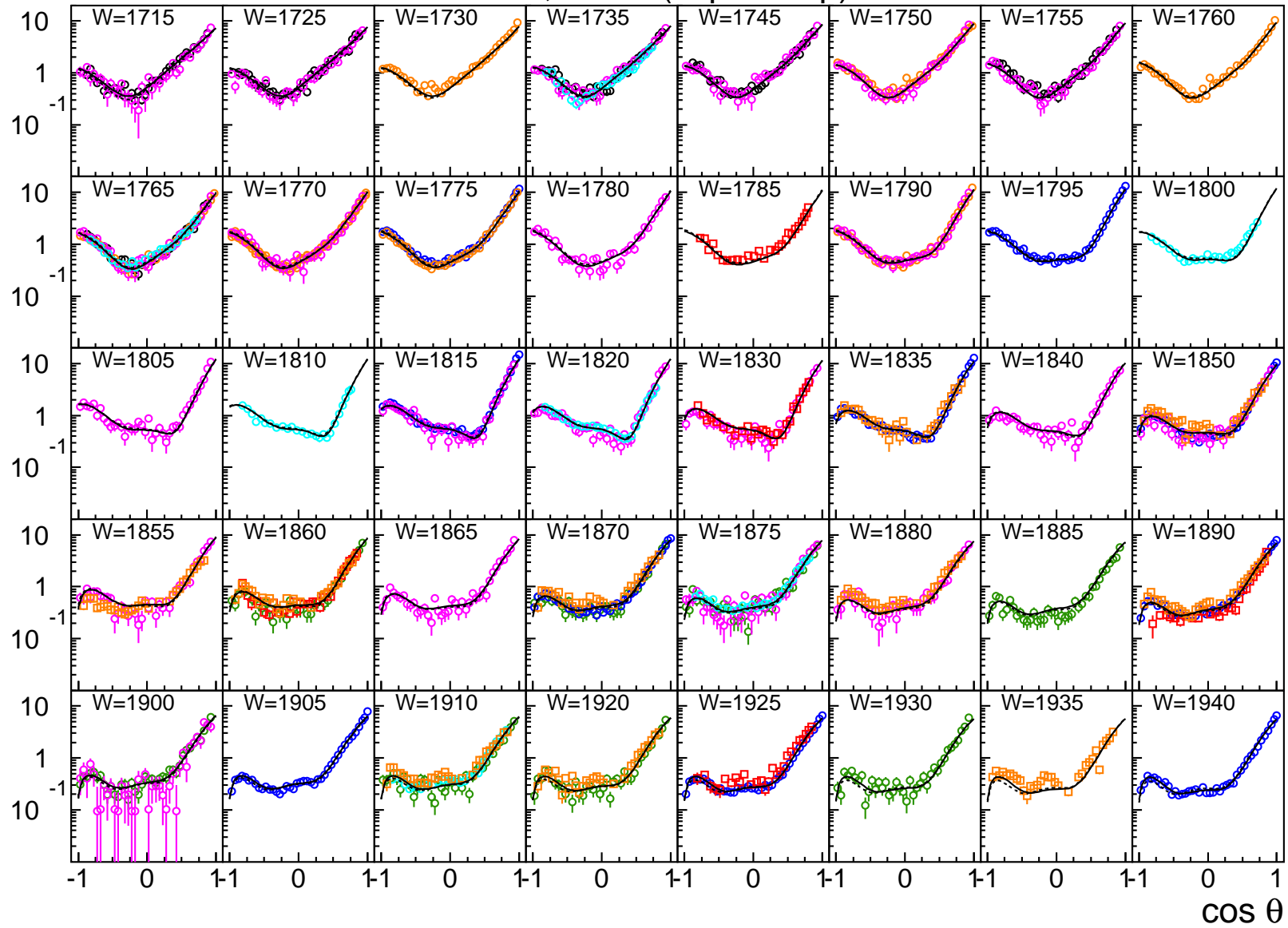
$K^- p \rightarrow K^0 n$	$K^- p \rightarrow K^- p$	$K^- p \rightarrow \omega \Lambda$
$K^- p \rightarrow \pi^0 \Lambda$	$K^- p \rightarrow \eta \Lambda$	$K^- p \rightarrow \pi^+ \Sigma^-$
$K^- p \rightarrow \pi^0 \Sigma^0$	$K^- p \rightarrow \pi^- \Sigma^+$	$K^- p \rightarrow \pi^0 \pi^0 \Lambda$
$K^- p \rightarrow K^+ \Xi^-$	$K^- p \rightarrow K^0 \Xi^0$	$K^- p \rightarrow \pi^0 \pi^0 \Sigma^0$

Mass scan of additional states

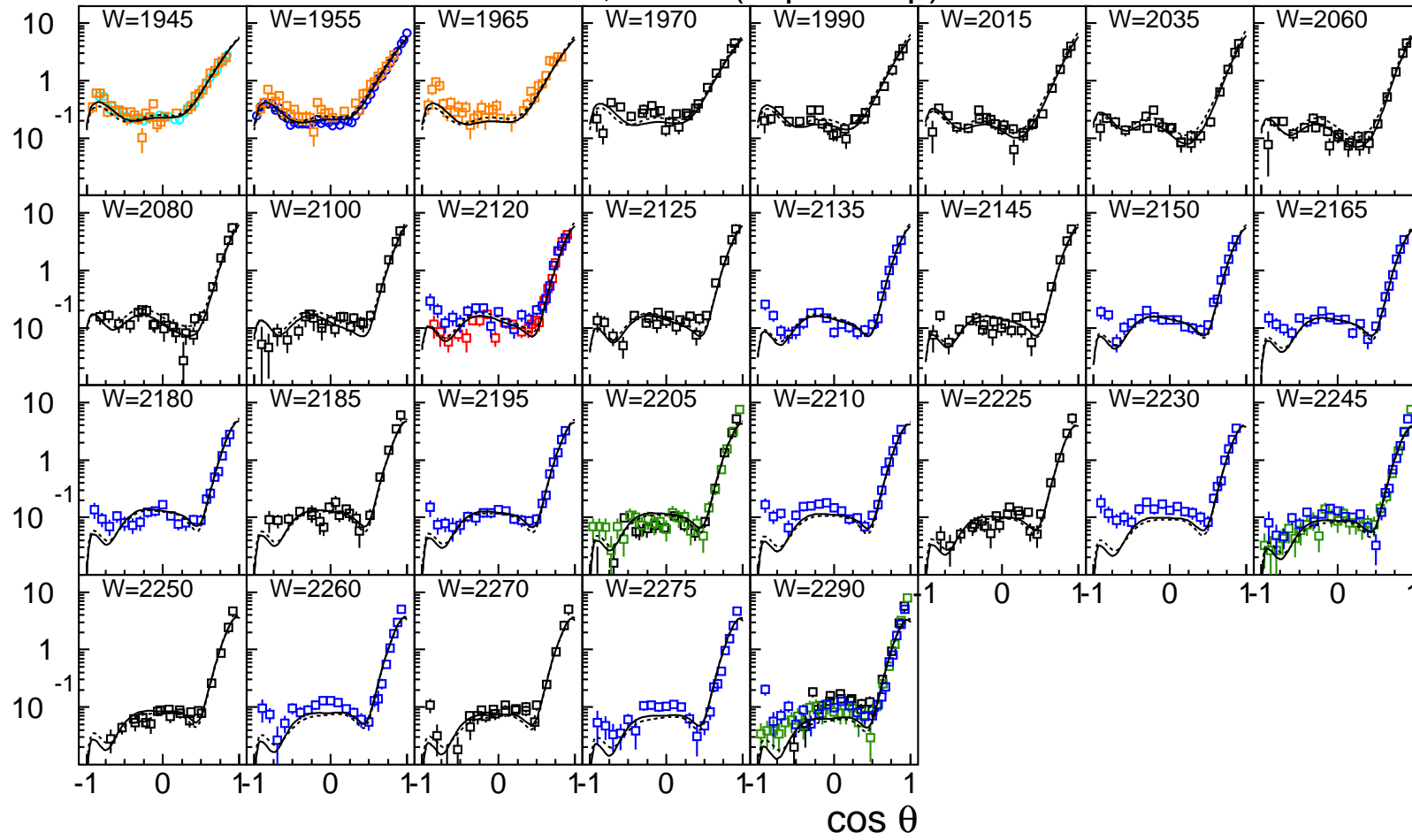




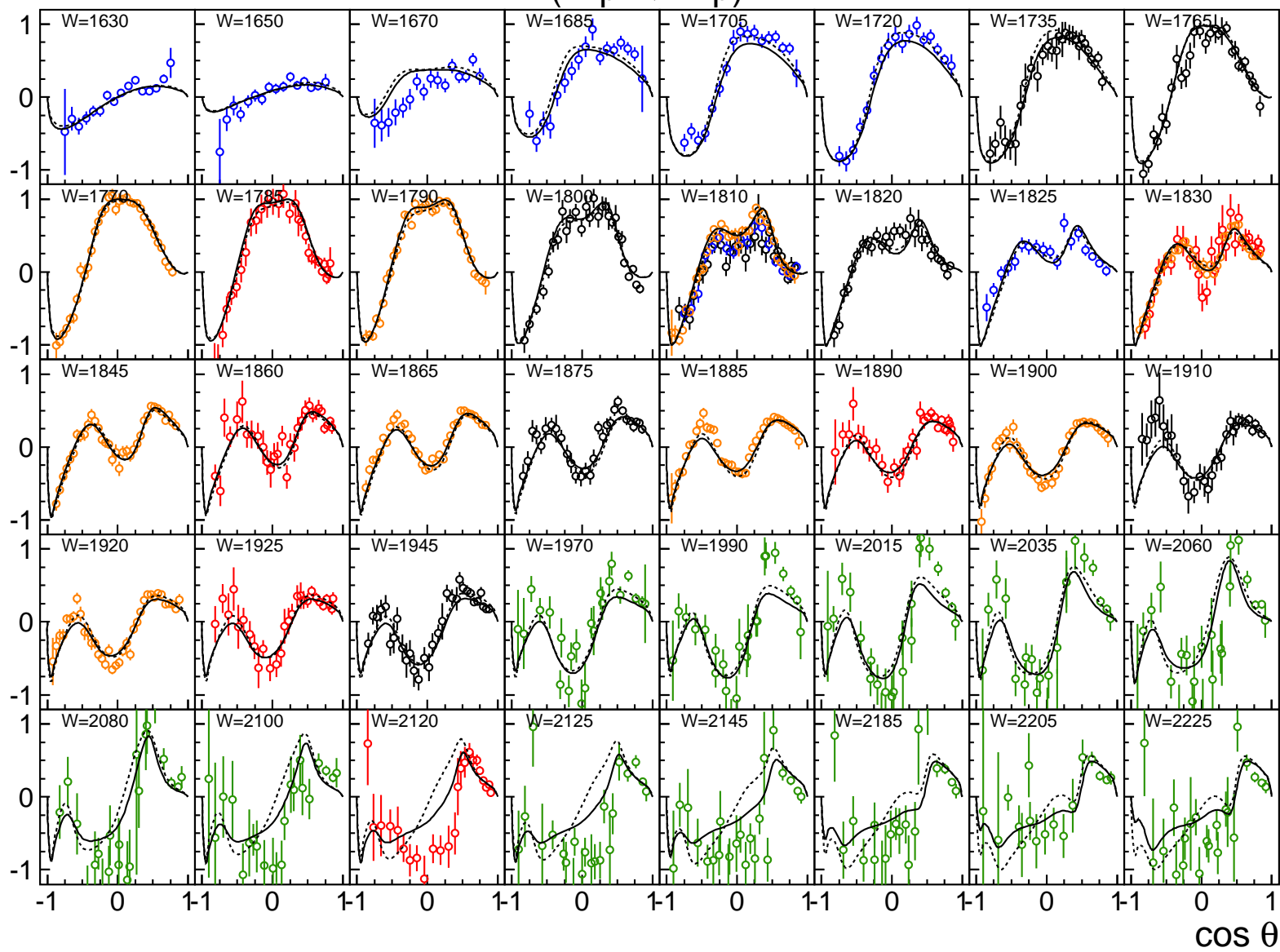
$d\sigma/d\Omega$, mb/sr ($K^-p \rightarrow K^-p$)



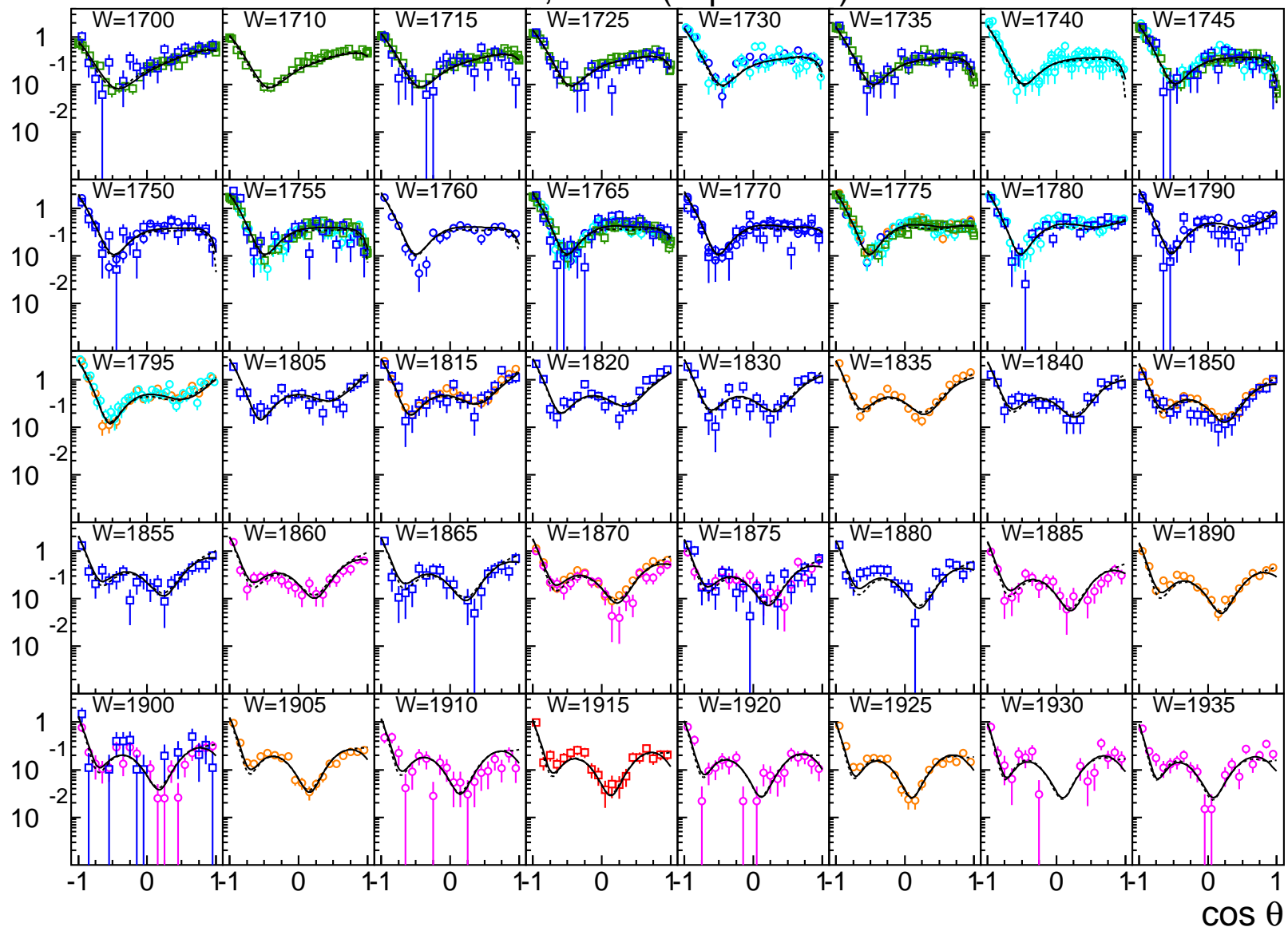
$d\sigma/d\Omega$, mb/sr ($K^- p \rightarrow K^- p$)



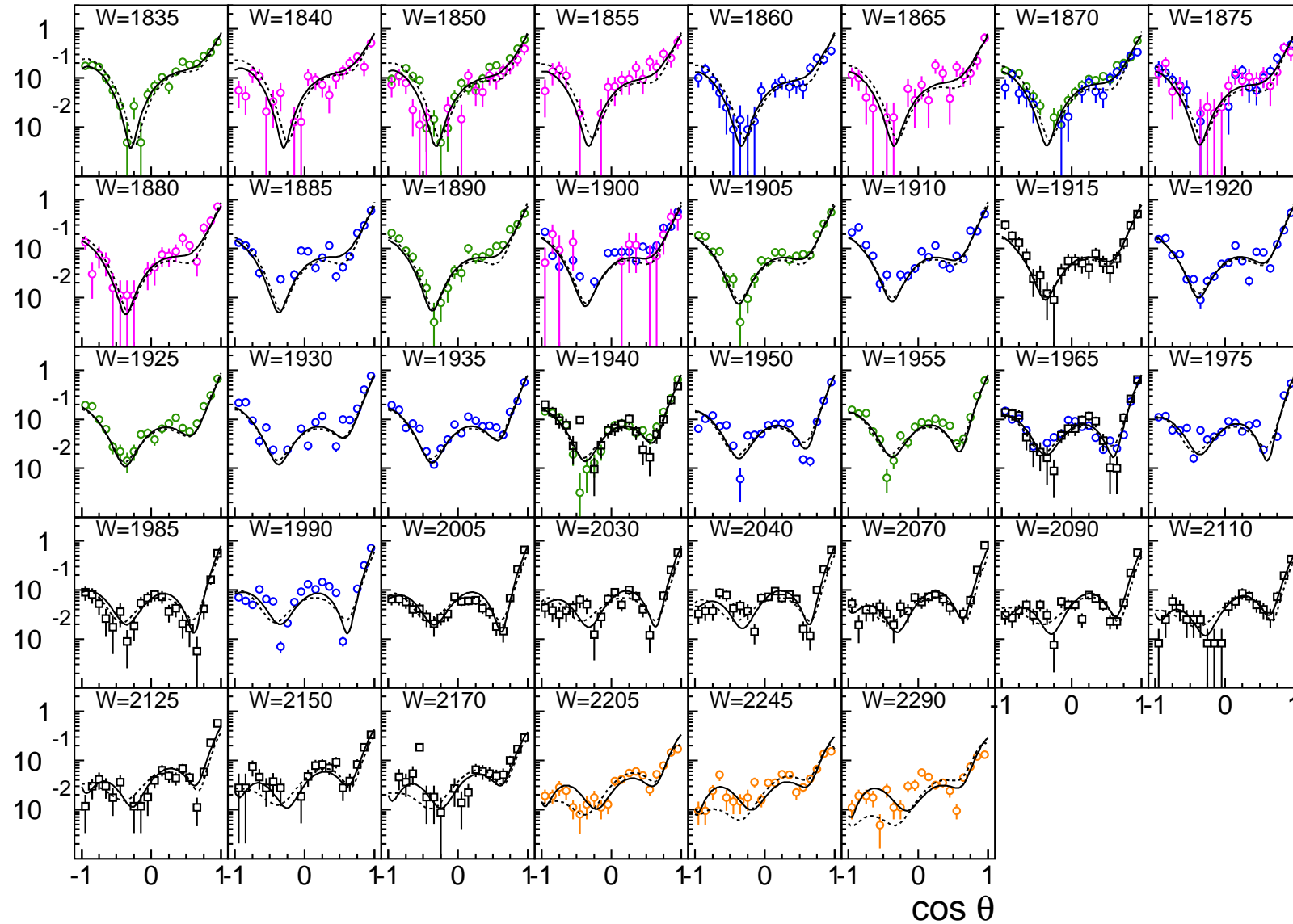
$P(K^-p \rightarrow K^-p)$

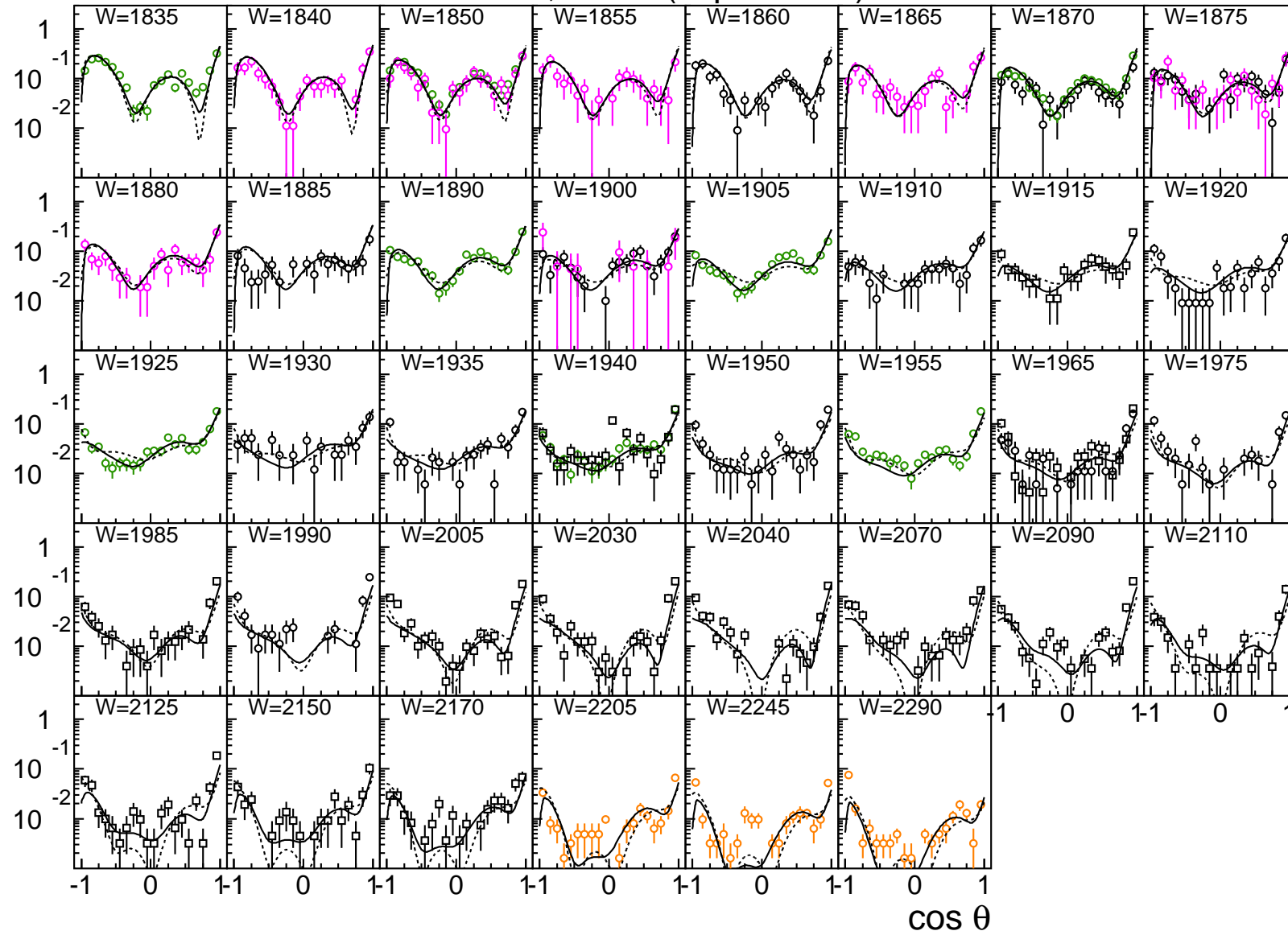


$d\sigma/d\Omega$, mb/sr ($K^-p \rightarrow K^0n$)



$d\sigma/d\Omega$, mb/sr ($K^-p \rightarrow \pi^- \Sigma^+$)



$$d\sigma/d\Omega, \text{ mb/sr } (K^- p \rightarrow \pi^+ \Sigma^-)$$


$d\sigma/d\Omega$, mb/sr ($K^-p \rightarrow \pi^0\Sigma^0$)

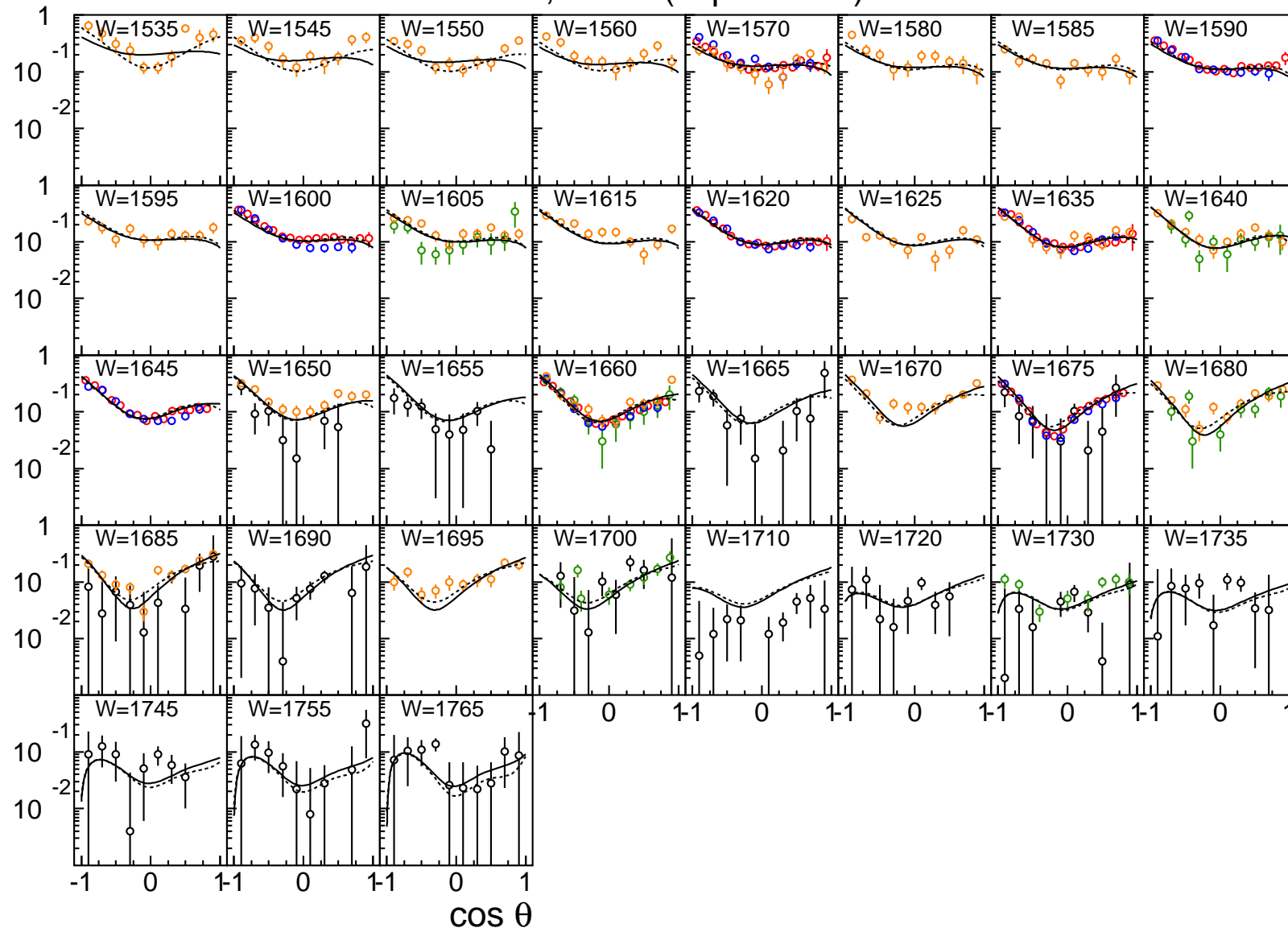
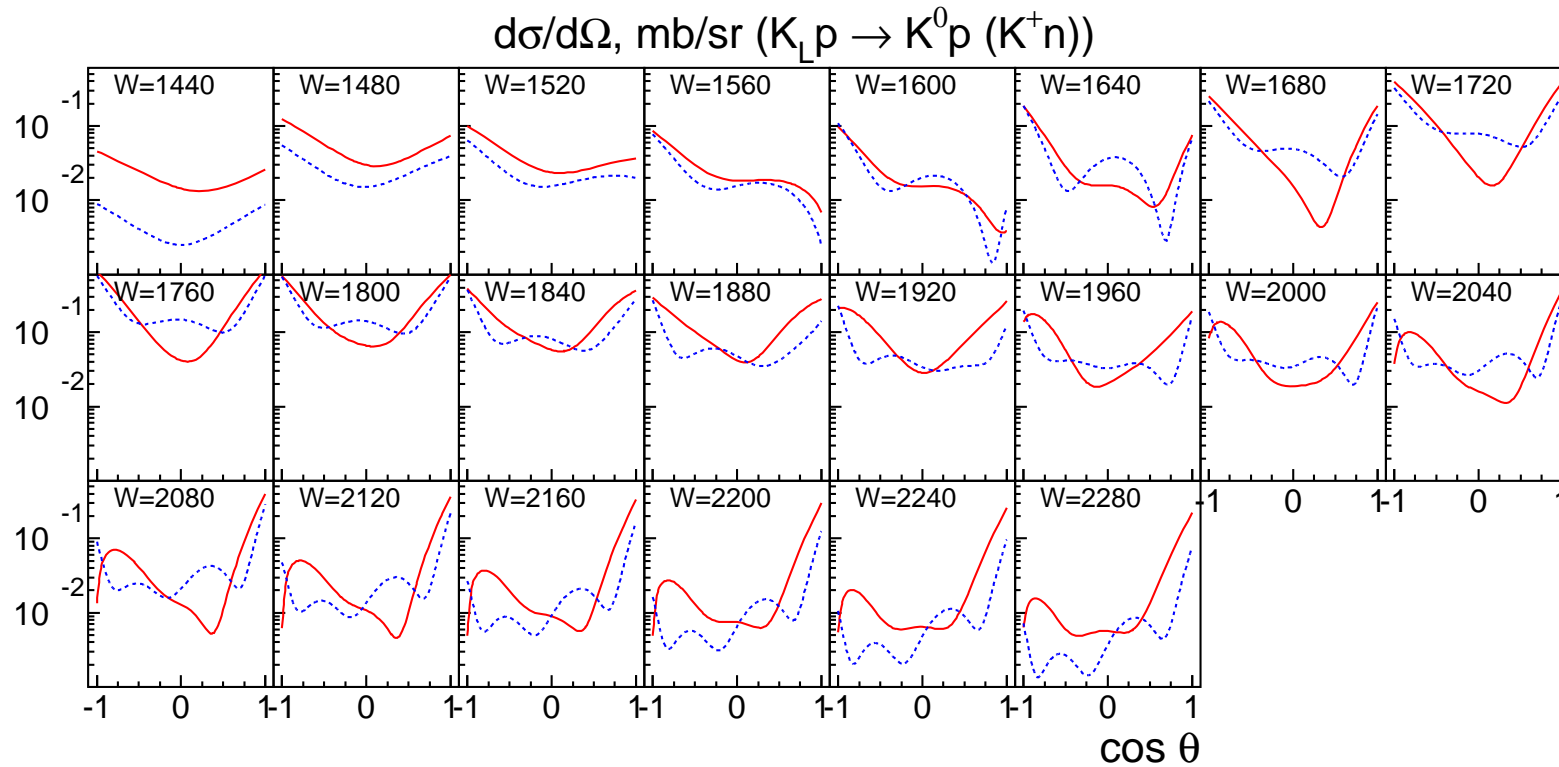


Table 4: Σ -Hyperons Observed states

J^P		Known state	New state	Mass
$1/2^+$	$N(1440)$	$\Sigma(1660)$		
$3/2^+$	$\Delta(1230)$	$\Sigma(1385)$		
$5/2^+$	$N(1680), \Delta(1905)$	$\Sigma(1915)$????	
$7/2^+$	$N(1990), \Delta(1950)$	$\Sigma(2030)$????	
$3/2^-$	$N(1520)$	$\Sigma(1670)$		
$1/2^-$	$N(1535), \Delta(1620), N(1650)$	$\Sigma(1750)$	$\Sigma(1620)$	1680 ± 8
			$\Sigma(1900)$	1936 ± 10
$5/2^-$	$N(1675)$	$\Sigma(1775)$		
$3/2^-$	$N(1700), \Delta(1700)$	$\Sigma(1940)$	$\Sigma(1860)$	1856 ± 10
$1/2^-$	$N(1895)$		$\Sigma(2120)$	2158 ± 25

Prediction for $\frac{d\sigma}{d\Omega} (K_L p \rightarrow K^0 p (K^+ n))$



Let us consider the decay of the isospin 0 and isospin 1 states into $K^- p$ and $K^0 n$

$$|A(K^- p)|^2 = \left(A_1 \frac{1}{\sqrt{2}} + A_0 \frac{1}{\sqrt{2}} \right)^2 = \frac{1}{2} (|A_1|^2 + |A_0|^2 + 2\text{Re}(A_1 A_0^*))$$

$$|A(K^0 n)|^2 = \left(A_1 \frac{1}{\sqrt{2}} - A_0 \frac{1}{\sqrt{2}} \right)^2 = \frac{1}{2} (|A_1|^2 + |A_0|^2 - 2\text{Re}(A_1 A_0^*))$$

$$A_{KN} = \omega^* [G(s, t) + H(s, t)i(\vec{\sigma}\vec{n})] \omega' \quad \vec{n}_j = \varepsilon_{\mu\nu j} \frac{q_\mu k_\nu}{|\vec{k}||\vec{q}|}.$$

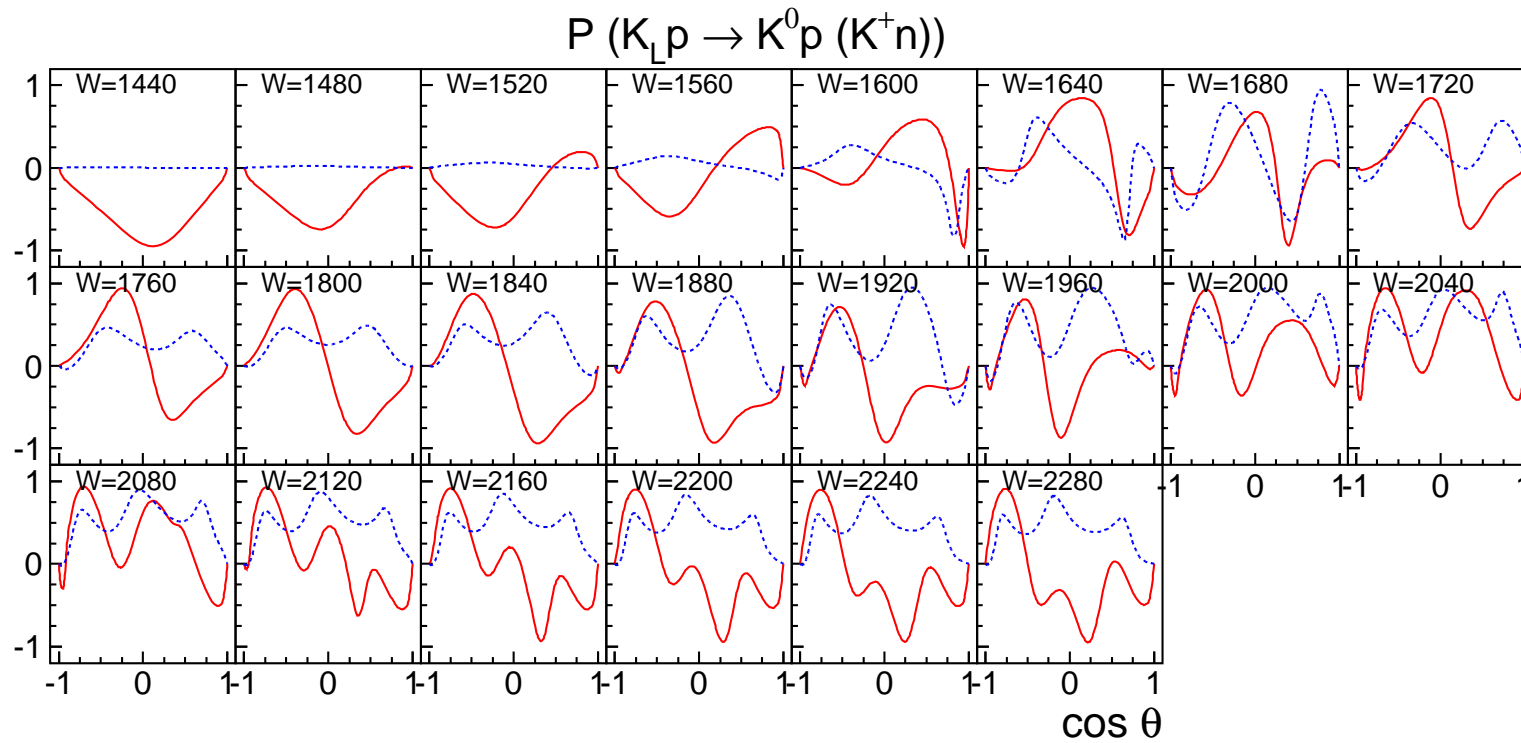
Differential cross section in c.m.s. of the reaction

$$|A|^2 = \frac{1}{2} \text{Tr} [A_{\pi N}^* A_{\pi N}] = |G(s, t)|^2 + |H(s, t)|^2 (1 - z^2)$$

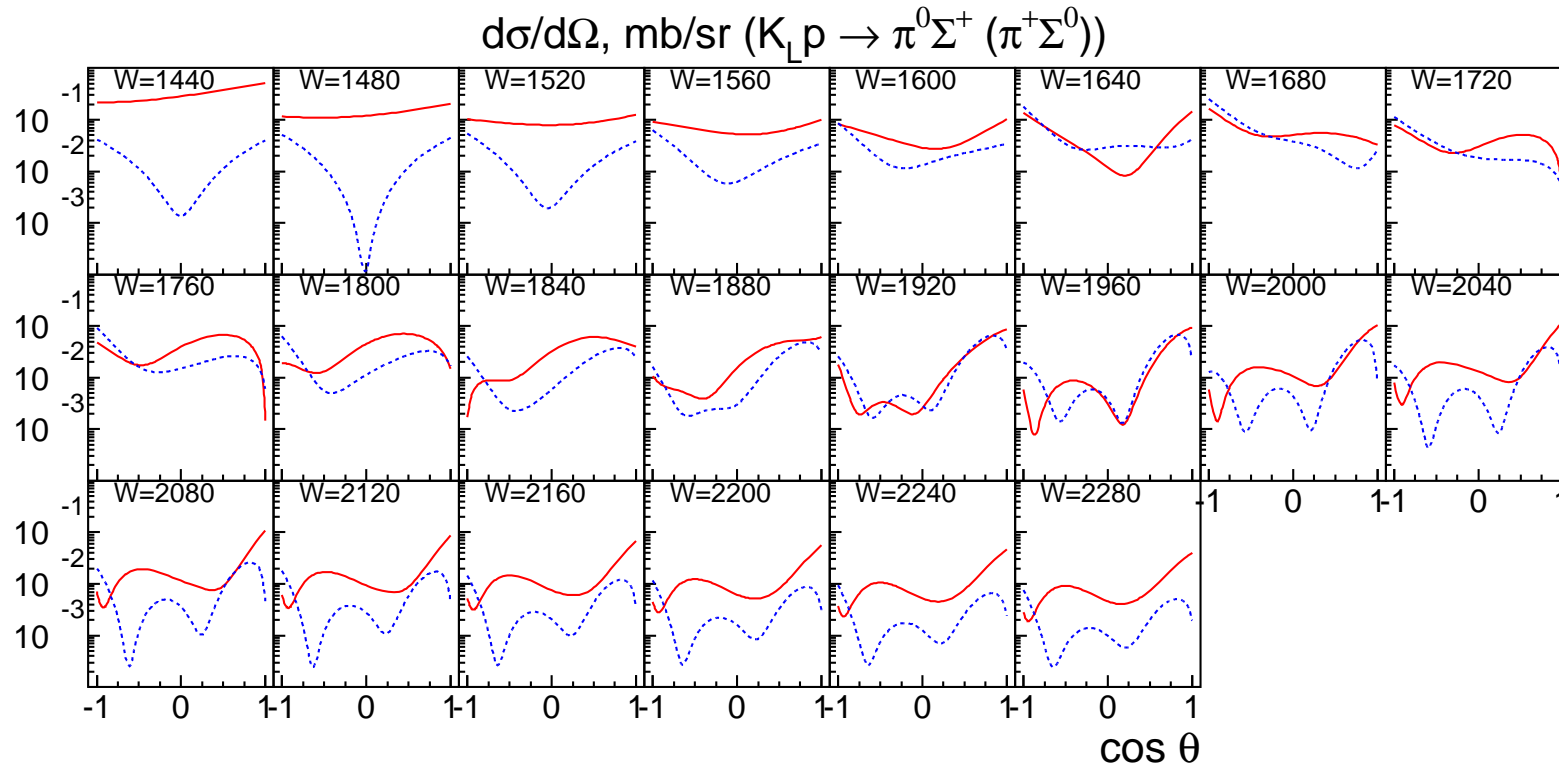
the recoil asymmetry:

$$P = \frac{\text{Tr} [A_{\pi N}^* \sigma_2 A_{\pi N}]}{2|A|^2 \cos \phi} = \sin \Theta \frac{2\text{Im} (H^*(s, t)G(s, t))}{|A|^2}.$$

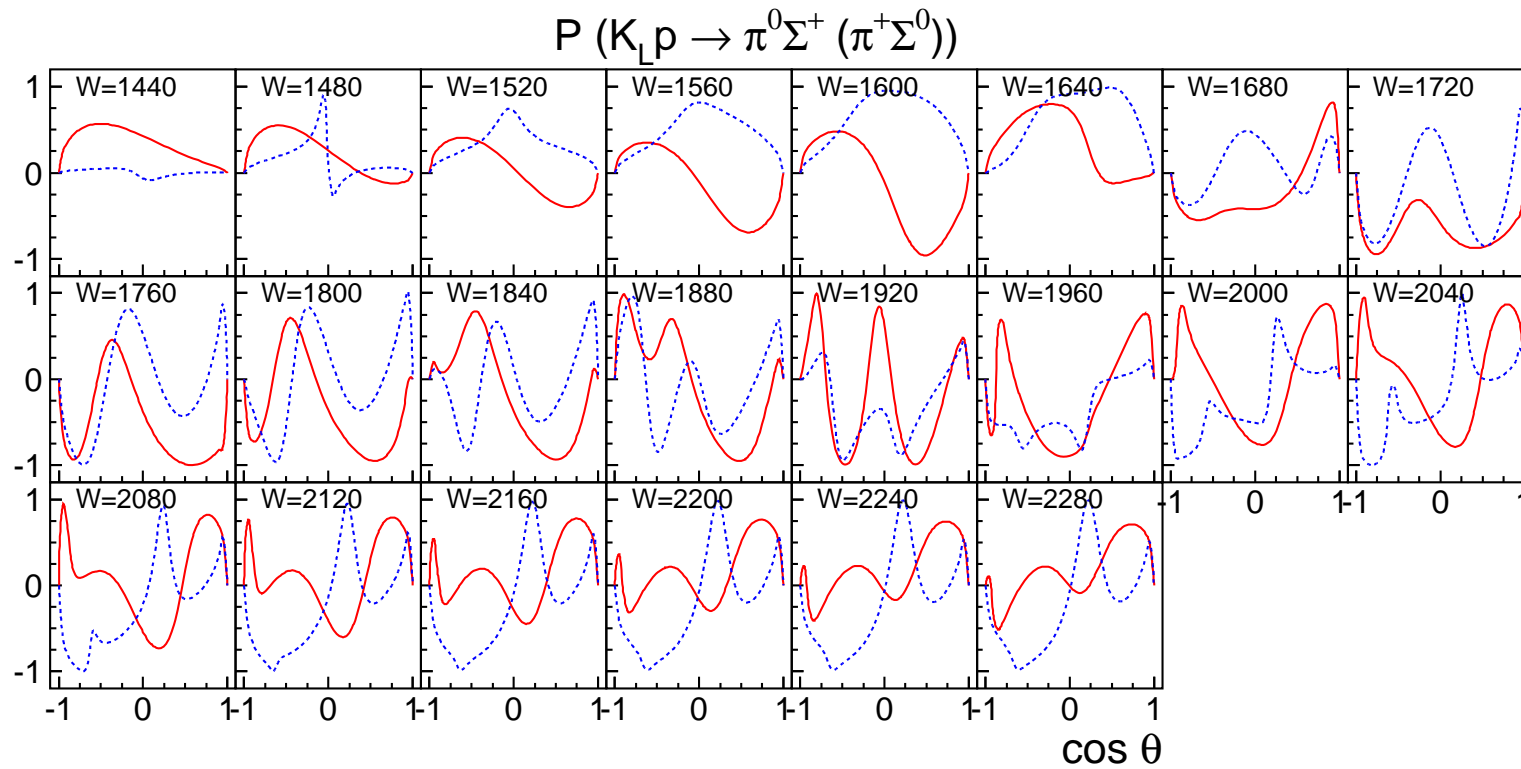
Prediction for the recoil asymmetry $K_L p \rightarrow K^0 p(K^+ n)$



Prediction for $\frac{d\sigma}{d\Omega} (K_L p \rightarrow \pi^0 \Sigma^+ (\pi^+ \Sigma^0))$



Prediction for the recoil asymmetry $K_L p \rightarrow \pi^0 \Sigma^+ (\pi^+ \Sigma^0)$



SUMMARY

- The planned $K_L p$ experiment provides a unique possibility to study the spectrum and properties of Σ hyperons.
- The data on the $K_L p$ collision are necessary to perform the full amplitude decomposition of the observables measured in the $K^- p$ collision reactions: therefore the new data are important for the determination of the spectrum and properties of Λ hyperons.