#### A Dalitz Plot Analysis of  $\omega \to 3\pi$  Decay

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<span id="page-0-0"></span>June 4, 2015

## **Overview**

Goals

- G12 Data
- **•** Fit Function in Brief
- Analysis
- **•** Next Steps

Main Goal

- Fit a modified Khuri-Treiman (KT) Model for the  $\omega \to 3\pi$  decay.
- Fit the decay amplitude via event-based likelihood fits using AmpTools framework.
- Compare fit parameters with known results from other models

## CLAS G12 Data



- G12 data covers incoming photon energy range  $E_{\gamma}$ : [1150-5400]MeV
- Have G12 x-section  $E_{\gamma}$ : [1150-3800]MeV, extending to 5400MeV, need G12 SDME's still
- We use G11 x-section, SDME's data; range  $E_{\gamma}$ : [1107.4-3828.9]MeV
- We consider bins in range  $E_{\gamma}$ :[1150-3800]MeV

## G12, G11 Cross-Section Comparison



 $E\gamma$ :[1500-2010]MeV, Zulkaida Akbar

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## G12, G11 Cross-Section Comparison



 $E\gamma$ :[2150-3150]MeV, Zulkaida Akbar

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## G12, G11 Cross-Section Comparison



 $E\gamma$ :[3150-3800]MeV, Zulkiada Akbar

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JLAB-THY-14-1960

#### Dispersive Analysis of  $\omega/\phi \rightarrow 3\pi$ ,  $\pi \gamma^*$

I.V. Danilkin,<sup>1,</sup> \* C. Fernández-Ramírez,<sup>1</sup> P. Guo,<sup>2,3</sup> V. Mathieu,<sup>2,3</sup> D. Schott,<sup>4</sup> M. Shi,<sup>1,5</sup> and A. P. Szczepaniak<sup>1, 2, 3</sup> (Joint Physics Analysis Center)

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(Dated: September 30, 2014)

The decays  $\omega/\phi \to 3\pi$  are considered in the dispersive framework that is based on the isobar decomposition and sub-energy unitarity. The inelastic contributions are parametrized by the power series in a suitably chosen conformal variable that properly account for the analytic properties of the amplitude. The Dalitz plot distributions and integrated decay widths are presented. Our results indicate that the final state interactions may be sizable. As a further application of the formalism we also compute the electromagnetic transition form factors of  $\omega/\phi \to \pi^0 \gamma^*$ .

#### arXiv:1409.7708v1 [hep-ph] (2014)

arXiv:1409.7708v1 [hep-ph] 26 Sep 2014

resonances [1]. With high precision data already avail-

and expected from Jefferson Lab  $\mathbb{Z}$ , in the near future it is the near future it in the near future it is

#### Illustration



FIG. 1: Isobar decomposition.

#### Isobar Model



FIG. 2: Crossed channel rescattering effects.

#### consequence of elastic unitarity requirement of model

#### Model

$$
F(s) = \Omega(s) \left( \frac{1}{\pi} \int_{s_{\pi}}^{s_i} ds' \frac{\rho(s') t^*(s')}{\Omega^*(s')} \frac{\hat{F}(s')}{s'-s} + \Sigma(s) \right). \tag{31}
$$

This is an alternative to the standard way which employs The Decay Amplitude the becay Amplitude

$$
\Sigma(s) = \sum_{i=0}^{\infty} a_i \,\omega^i(s) \tag{29}
$$

inelastic contribution  $p_{\text{reduced}}$  since it is restricted to the elastic-region,  $p_{\text{reduced}}$  $a_0 =$  igorParameter)  $(a_0 = \text{lgorParameter})$ inelastic contribution  $(a_0 = \text{lgorParameter})$ <br>Dalitz Analysis of  $\omega \to 3\pi$  June 4, 2015 10 / 48

## Dalitz Analysis



dzie dzie

$$
x = \frac{\sqrt{3}}{Q}(T_1 - T_2) = \frac{\sqrt{3}(t - u)}{2M(M - 3m_{\pi})},
$$
  

$$
y = \frac{3T_3}{Q} - 1 = \frac{3(s_c - s)}{2M(M - 3m_{\pi})}.
$$
 (38)

#### is the kinetic energy of the i-th pion in the three-

 $y = \sqrt{z} \sin \theta$  $x = \sqrt{z} \cos \theta$  $y = \sqrt{z} \sin \theta$  $y - \sqrt{2} \sin \theta$ 

<sup>2</sup>

$$
|F_{par}(z,\vartheta)|^2 = |N|^2 (1 + 2\alpha z + 2\beta z^{3/2} \sin(3\vartheta) + 2\gamma z^2 + 2\delta z^{5/2} \sin(3\vartheta) + \mathcal{O}(z^3))
$$
(40)

#### Dalitz Plot Amplitude Expansion Dalitz Plot Amplitude Expansion

This paper (full) 80 27 8 - 0.24  $\overline{\phantom{a}}$  and  $\overline{\phantom{a}}$  of  $\overline{\phantom{a}}$  of  $\overline{\phantom{a}}$  and  $\overline{\phantom{a}}$  and  $\overline{\phantom{a}}$  and  $\overline{\phantom{a}}$ 

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# Dalitz Plot Parameter Comparison

$$
|F_{par}(z,\vartheta)|^2 = |N|^2 (1 + 2 \alpha z + 2 \beta z^{3/2} \sin(3\vartheta) + 2 \gamma z^2 + 2 \delta z^{5/2} \sin(3\vartheta) + \mathcal{O}(z^3))
$$
(40)

TABLE I: Dalitz Plot parameters and  $\sqrt{\overline{\chi}^2}$  of the polynomial parametrization (40) for  $\omega \to 3\pi$ . In addition to our results we also show the selected results from Niecknig et al.  $[37]$  (dispersive study with incorporated crossed-channel effects) and<br>Torschluson et al. [10] (Lagrangian based study with the pion pion rescattering effects) Terschlusen et al. [19] (Lagrangian based study with the pion-pion rescattering effects).

	$\alpha \times 10^3$	$\beta \times 10^3$	$\gamma \times 10^3$	$\delta \times 10^3$	$\sqrt{\bar{X}^2} \times 10^3$
This paper $(\hat{F} = 0)$	136	н.	÷	$\overline{\phantom{a}}$	$3.5\,$
This paper (full)	94	$\overline{\phantom{a}}$	٠	$\overline{\phantom{a}}$	3.2
Niecknig et al. [37]	8496			٠	0.91.1
Terschlusen et al. [19]	202	$\overline{\phantom{a}}$	$\overline{\phantom{a}}$	۰	$6.6\,$
This paper $(\hat{F} = 0)$	125	30		÷	0.74
This paper (full)	84	28	٠	٠	0.35
Niecknig et al. [37]	7484	2428	٠	٠	0.0520.078
Terschlusen et al. [19]	190	54	$\overline{\phantom{0}}$	۰	2.1
This paper $(\hat{F} = 0)$	113	27	24		0.1
This paper (full)	80	27	8		0.24
Niecknig et al. [37]	7381	2428	36	٠	0.0380.047
Terschlusen et al. [19]	172	43	50	٠	$0.4\,$
This paper $(\hat{F} = 0)$	114	24	20	6	0.005
This paper (full)	83	22		14	0.079
Niecknig et al. [37]	7483	2124	02	78	0.0120.011
Terschlusen et al. [19]	174	35	43	$20\,$	0.1



 $E\gamma$ :[1500-1600]MeV and t:[0.100-0.200]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.300-0.400]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.500-0.600]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.700-0.800]GeV<sup>2</sup>



 $E\gamma$ :[1500-1600]MeV and t:[0.900-1.000]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.100-0.200]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.300-0.400]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.500-0.600]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.700-0.800]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.900-1.000]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.100-0.200]GeV<sup>2</sup>



 $E\gamma$ :[1500-1600]MeV and t:[0.300-0.400]GeV<sup>2</sup>



 $E\gamma$ :[1500-1600]MeV and t:[0.500-0.600]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.700-0.800]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.900-1.000]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.100-0.200]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.300-0.400]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.500-0.600]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.700-0.800]GeV<sup>2</sup>

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 $E\gamma$ :[1500-1600]MeV and t:[0.900-1.000]GeV<sup>2</sup>

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## Analysis: Decay Amplitude Parameter (IgorParameter)



 $E\gamma$ :[1500-1600]MeV and t:[0.400-1.100]GeV<sup>2</sup>

## Analysis: Decay Amplitude Parameter (IgorParameter)



 $E\gamma$ :[2000-2100]MeV and t:[0.400-1.300]GeV<sup>2</sup>

## Analysis: Decay Amplitude Parameter (IgorParameter)



 $E\gamma$ :[2500-2600]MeV and t:[0.400-3.100]GeV<sup>2</sup>

# Analysis: Decay Amplitude Acceptance (IgorParameter)



 $E\gamma$ :[1500-1600]MeV and t:[0.400-1.100]GeV<sup>2</sup>

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# Analysis: Decay Amplitude Acceptance (IgorParameter)



 $E\gamma$ :[2000-2100]MeV and t:[0.400-1.300]GeV<sup>2</sup>

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# Analysis: Decay Amplitude Acceptance (IgorParameter)



 $E\gamma$ :[2500-2600]MeV and t:[0.400-3.100]GeV<sup>2</sup>

## Next Steps

**o** discussed

## Analysis: Dalitz Plot Parameter  $(\alpha)$



 $E\gamma$ :[2500-3000]MeV and t:[0.400-3.400]GeV<sup>2</sup>

## Analysis: Dalitz Plot Parameters  $(\alpha,\beta)$



 $E\gamma$ :[2500-3000]MeV and t:[0.400-3.400]GeV<sup>2</sup>

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## Analysis: Dalitz Plot Parameters  $(\alpha,\beta,\gamma)$



 $E\gamma$ :[2500-3000]MeV and t:[0.400-3.400]GeV<sup>2</sup>

## Analysis: Dalitz Plot Parameters  $(\alpha, \beta, \gamma, \delta)$



 $E\gamma$ :[2500-3000]MeV and t:[0.400-3.400]GeV<sup>2</sup>

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## Analysis: Integrated Populations and Acceptance



 $E\gamma$ :[2500-3000]MeV and t:[0.400-3.400]GeV<sup>2</sup>

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#### The Decay Amplitude The inelastic contribution is described by an analytical by an analy function on the s-plane cut along the real axis above

$$
F(s) = \Omega(s) \left( \frac{1}{\pi} \int_{s_{\pi}}^{s_i} ds' \frac{\rho(s') \, t^*(s')}{\Omega^*(s')} \frac{\hat{F}(s')}{s'-s} + \Sigma(s) \right). \tag{31}
$$

This is an alternative to the standard way which employs KT Amplitude  $\mathsf{k}$  $T_{\text{min}}$ 

$$
\Sigma(s) = \sum_{i=0}^{\infty} a_i \,\omega^i(s) \qquad (29) \qquad \omega(s) = \frac{\sqrt{s_i - s_E} - \sqrt{s_i - s}}{\sqrt{s_i - s_E} + \sqrt{s_i - s}} \qquad (30)
$$

#### $\overline{1}$ inelastic contribution

c contribution and inelastic dispersion relation on dispersion relation

$$
a_0 = \text{ "lgorParameter"}
$$

count the contributions from the more distant left-hand successfully applied in other descriptions of two-to-two amplitudes e.g. in [49–51] it was used to take into ac-

# Previous Parameter Values and Errors,  $E\gamma$ :[2500-3500]MeV



Note: Broken paddle, Eγ:[3000-3100]MeV

## Previous Comments by Mike P., Adam S., Igor D. and Carlos S.

- Fitted values are unstable. Have 8M data and 1.5M MC accepted, Need at lest the same amount of of acc-MC and data.
- Could you bin in t? Your rho parameters are t dependent and that trend will reflect in your fits and in the Igor's parameter that you will finally want to extract.
- Adam suggested that you may have too many parameters in the fit, as rho00 is also part of the overall normalization.
- 4) It will be good to start with a fit using  $F = 1$  (w/o Igor correction), and extract the rho parameters in that way to compare with  $g11$ results, and then include the new Igor parameter.
- <span id="page-47-0"></span>5) Igor's parameter should be independent of production, i.e. on Ebeam and t.