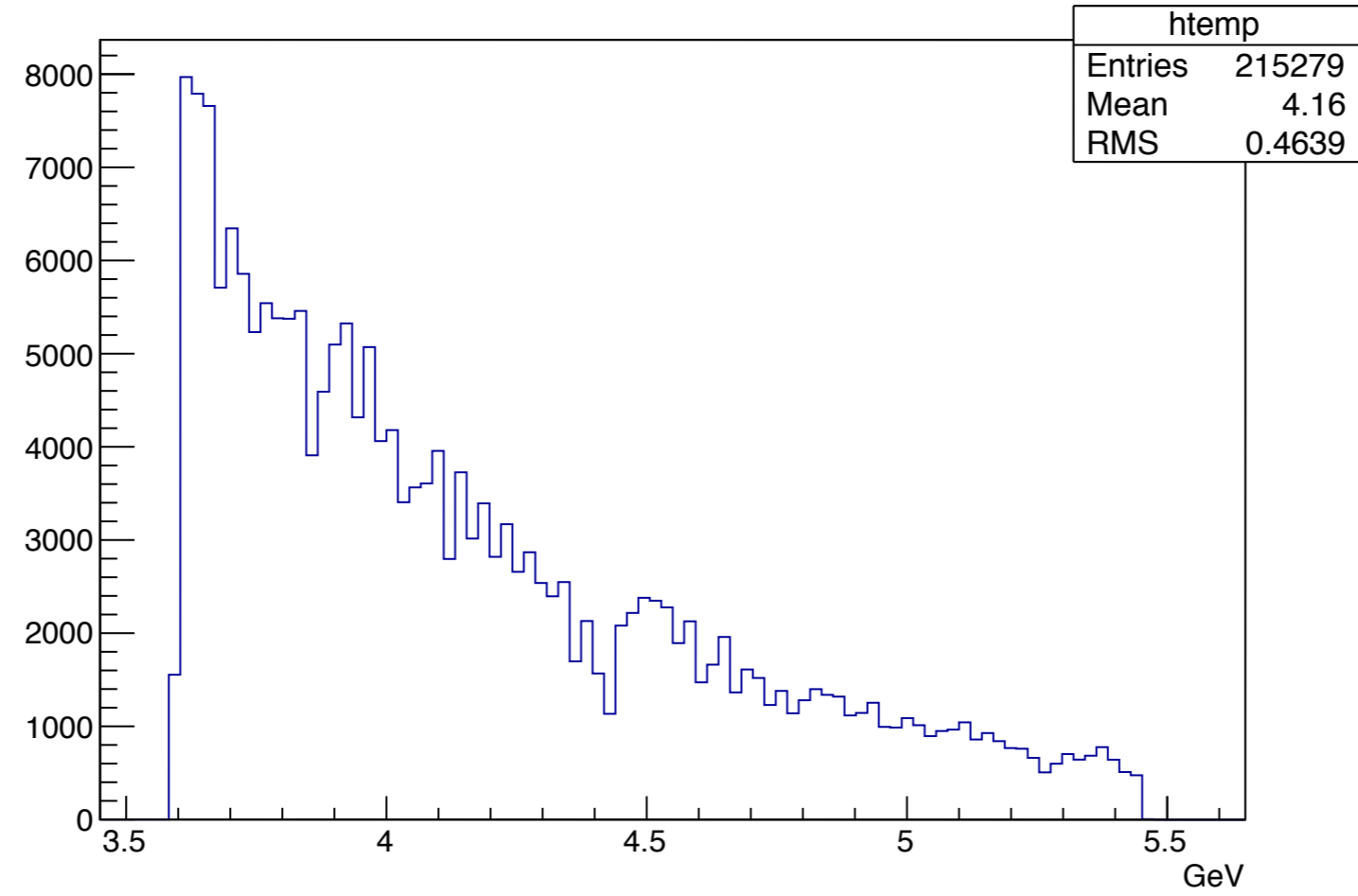
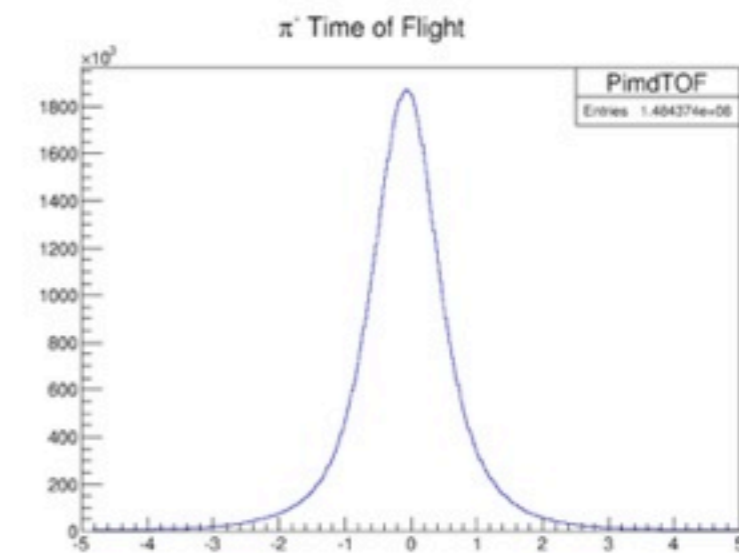
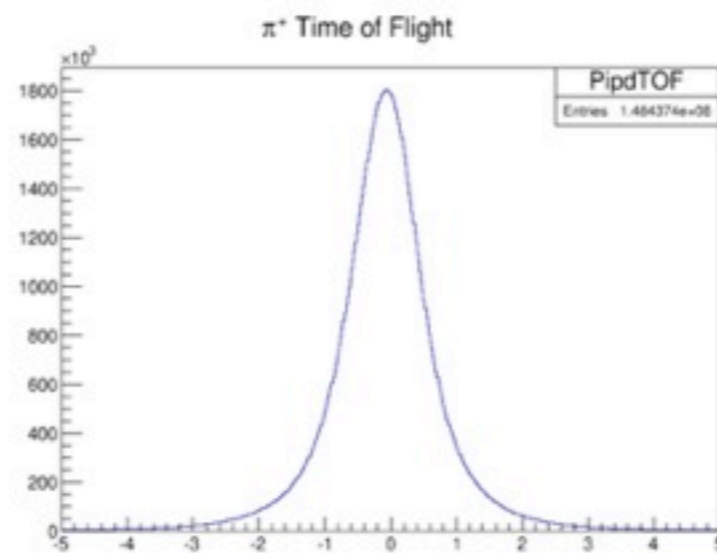
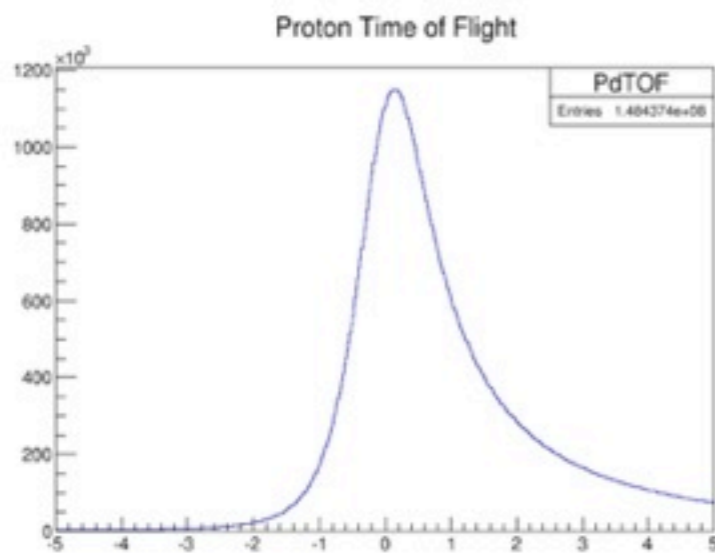
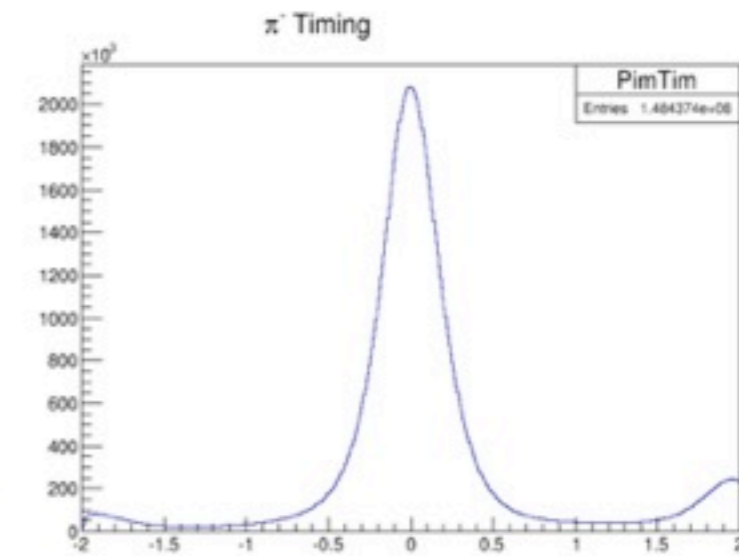
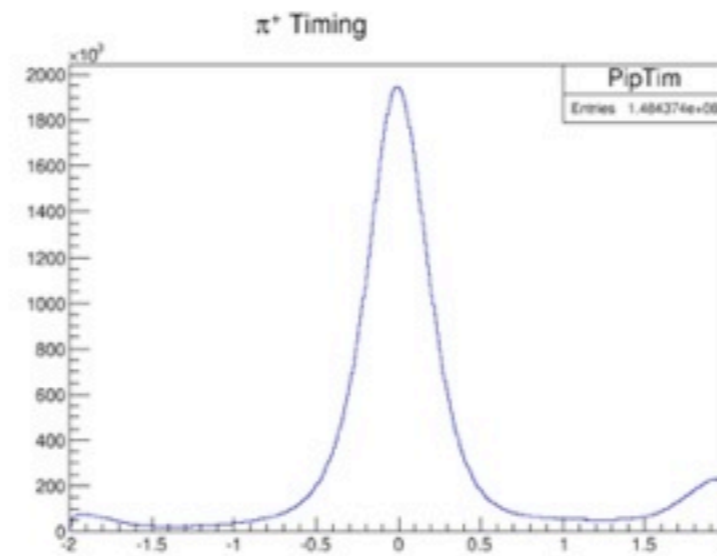
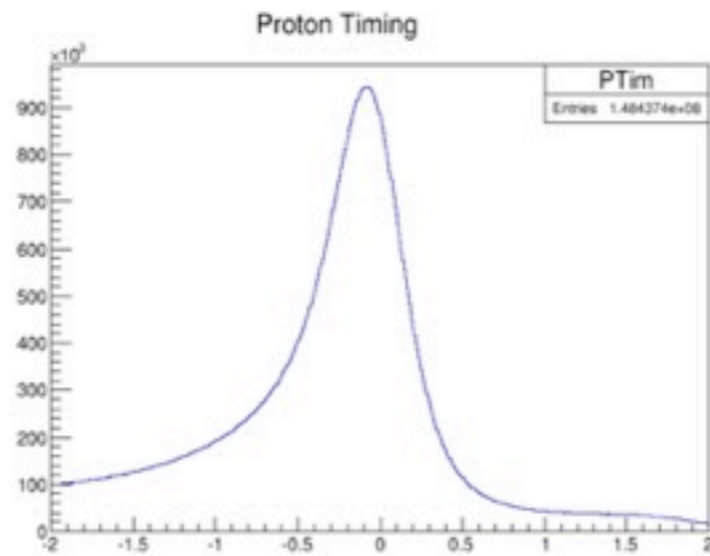


Omega Beam Energies



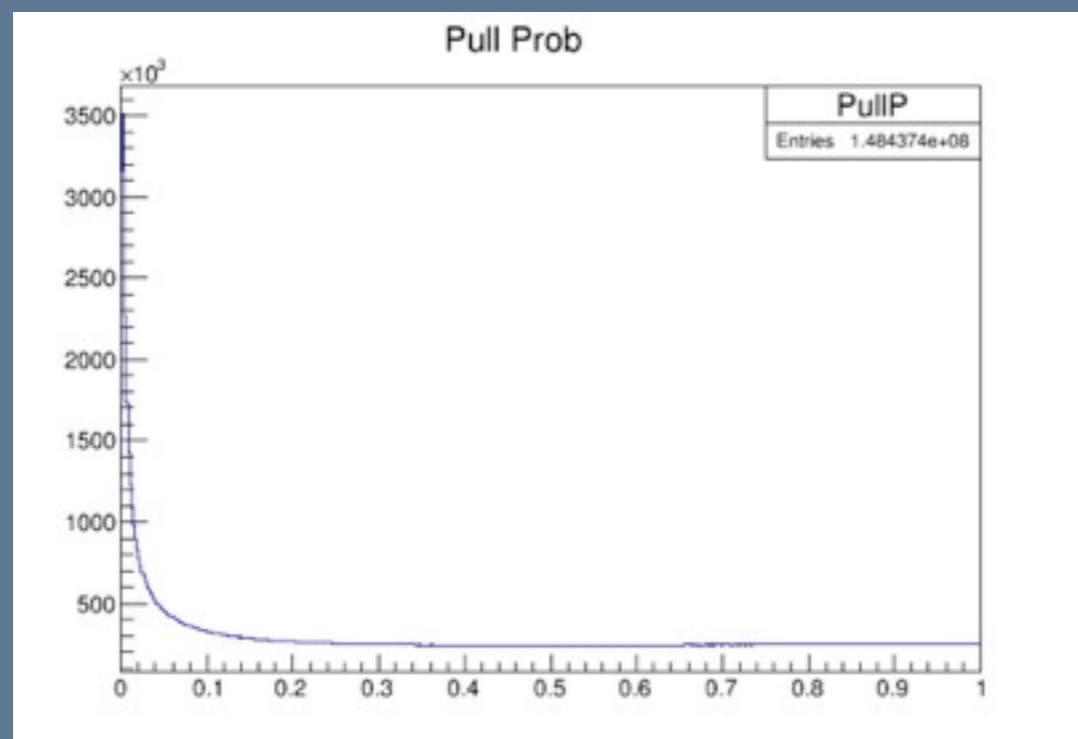
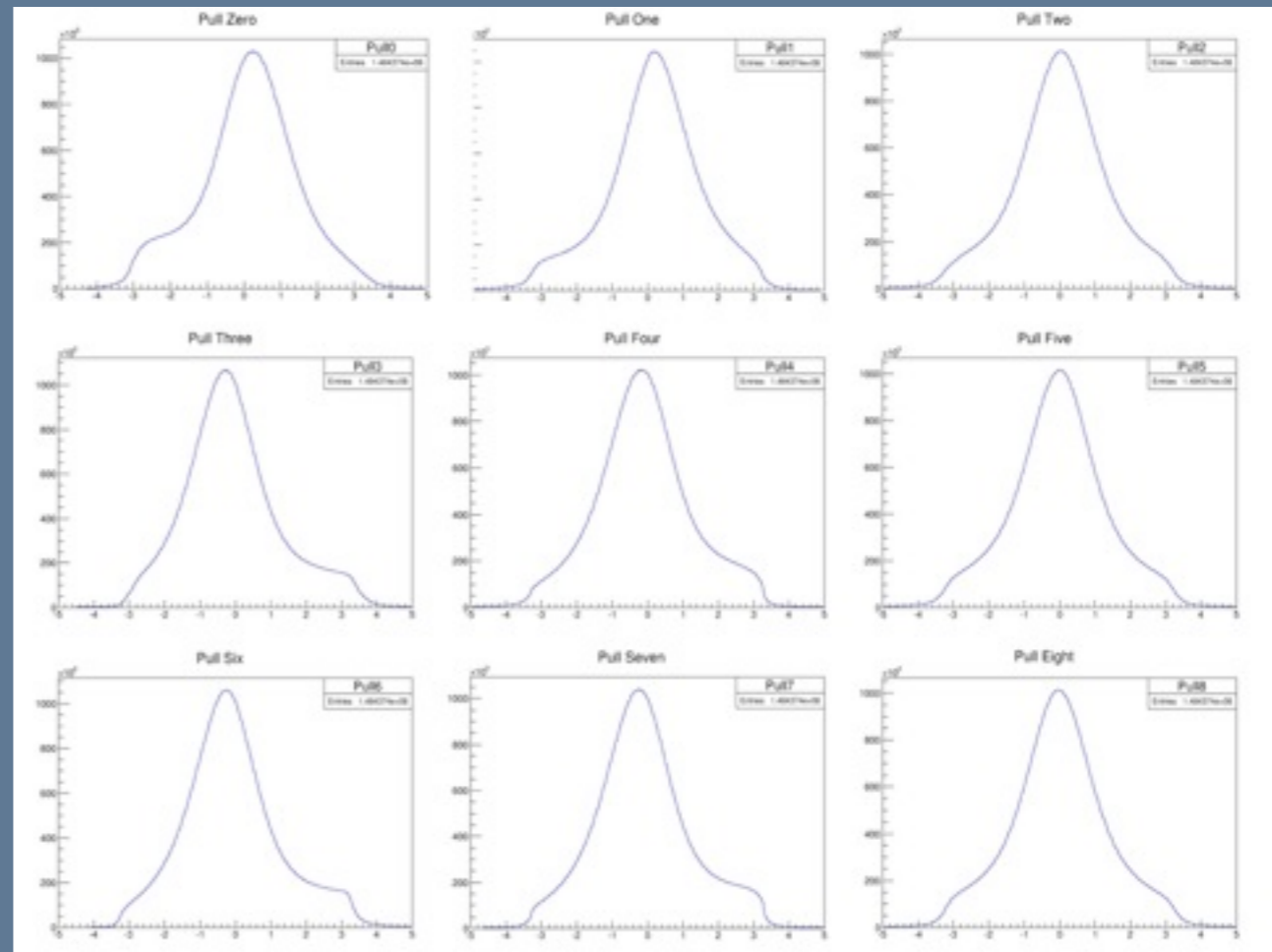
We also performed cuts on the particles timing (ToF), which is the difference between the time of flight measured (ToF-SC) and the calculated through the pathlength assuming a given ID mass. We used

$$(dTOF \pi^+)^2 + (dTOF \pi^-)^2 < 4, (dTOF p)^2 + (dTOF \pi^-)^2 < 4$$



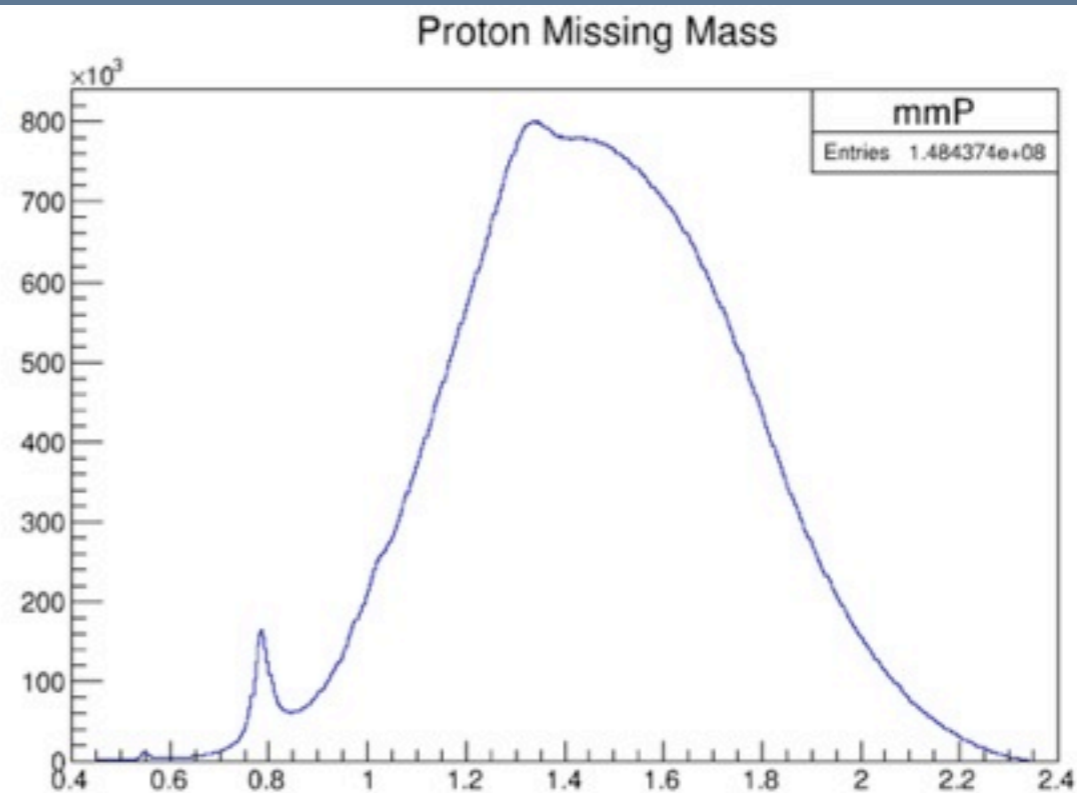
Kinematical Fit to π^0

using standard g12 KF -> see g12 general procedures note

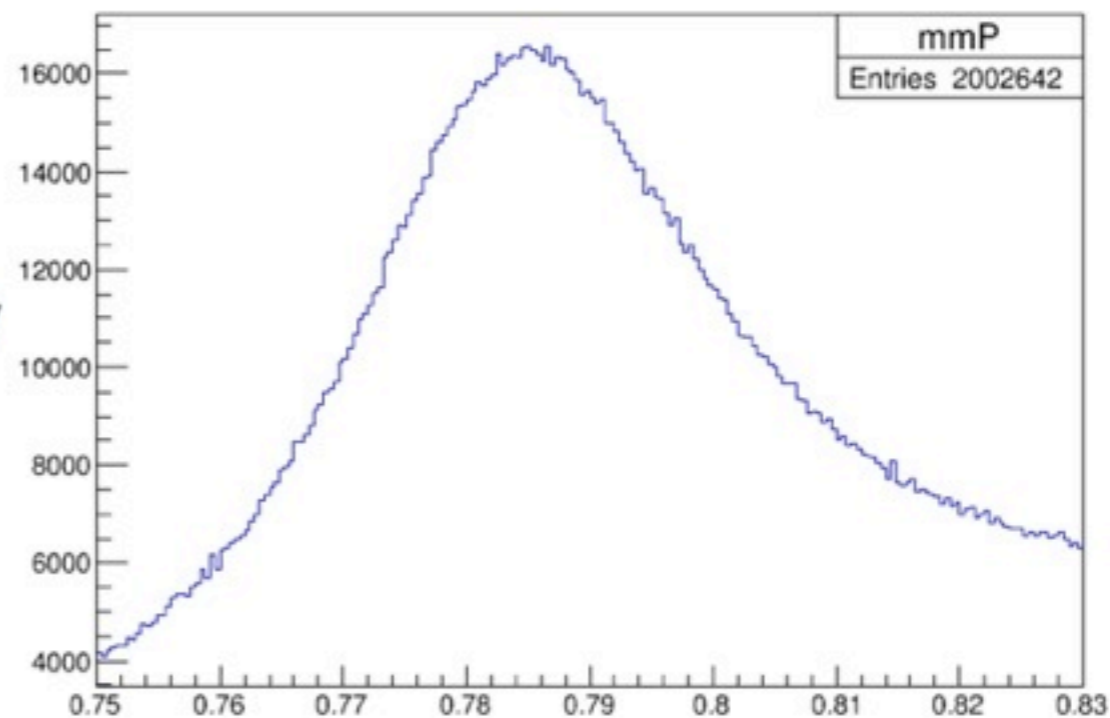




Full mass region \rightarrow

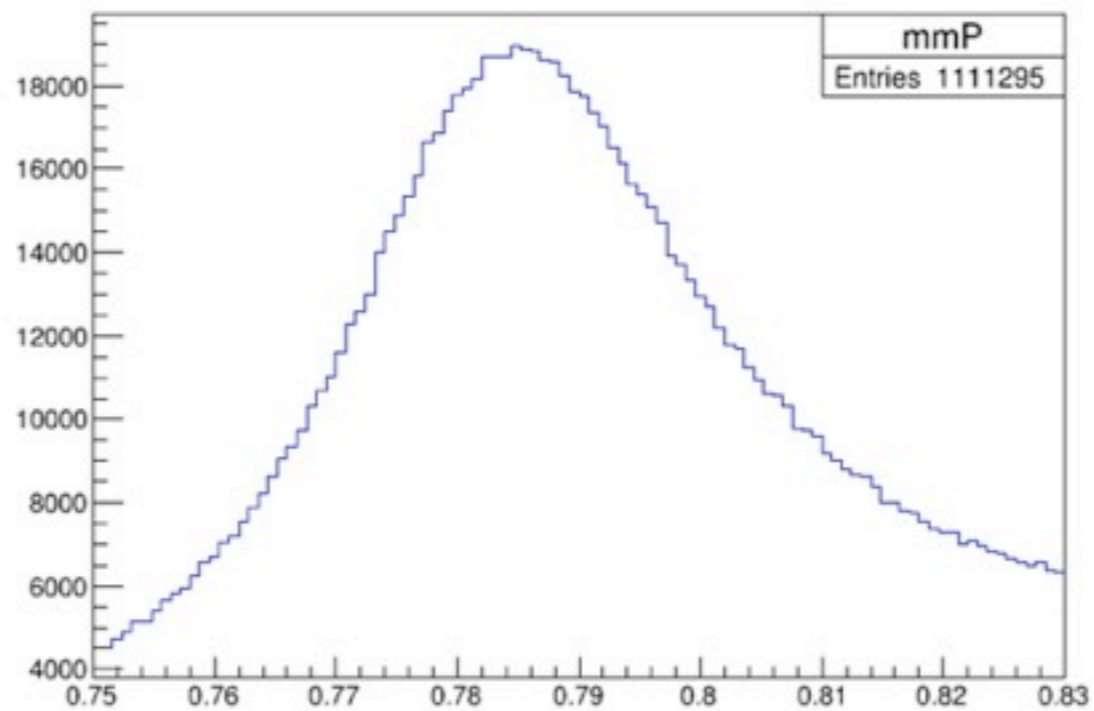


Proton Missing Mass (energy cut)

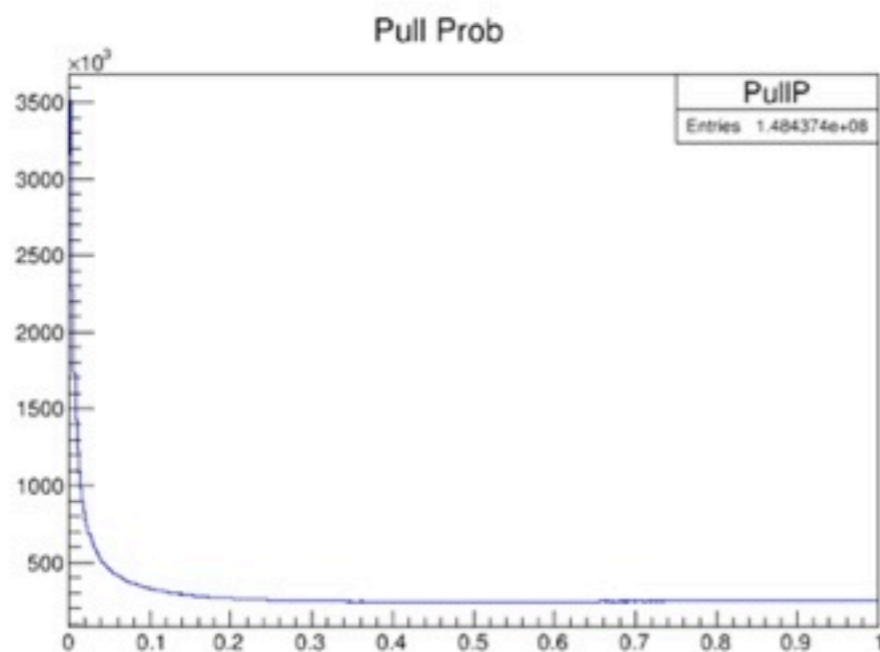


We will focus on the omega mass zone:
 $750 \text{ MeV} < \text{proton missing mass} < 830 \text{ MeV}$
and perform the mentioned energy cut:
 $3.6 \text{ GeV} < \text{photon energy}$

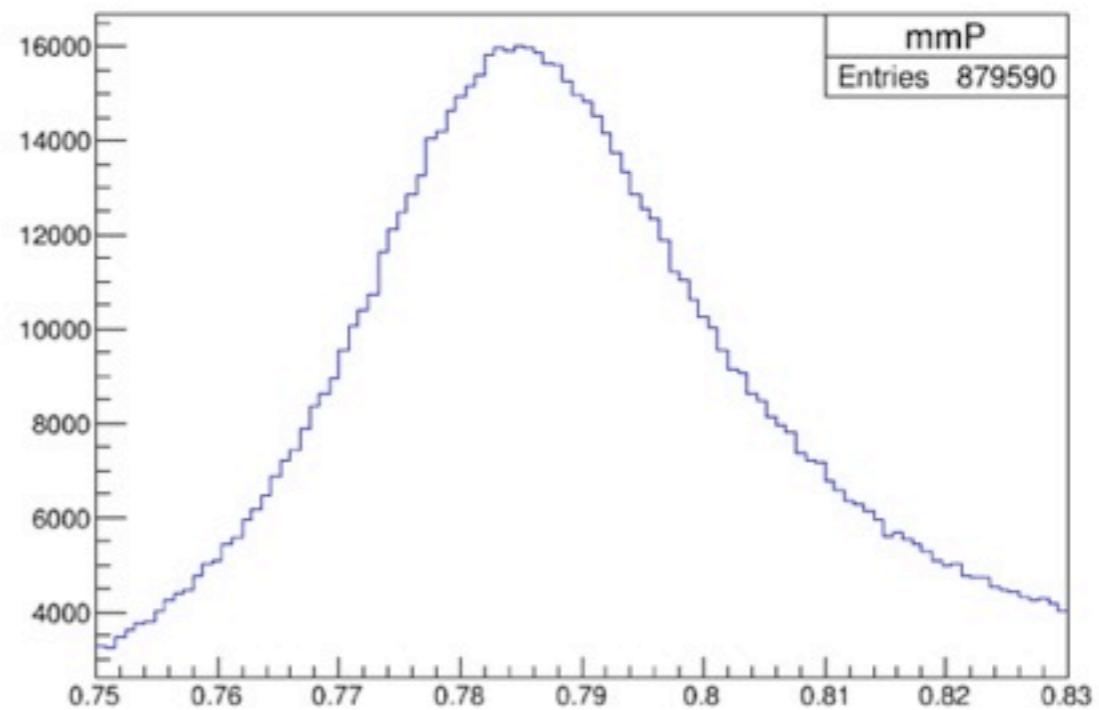
Proton Missing Mass (energy and time cuts)



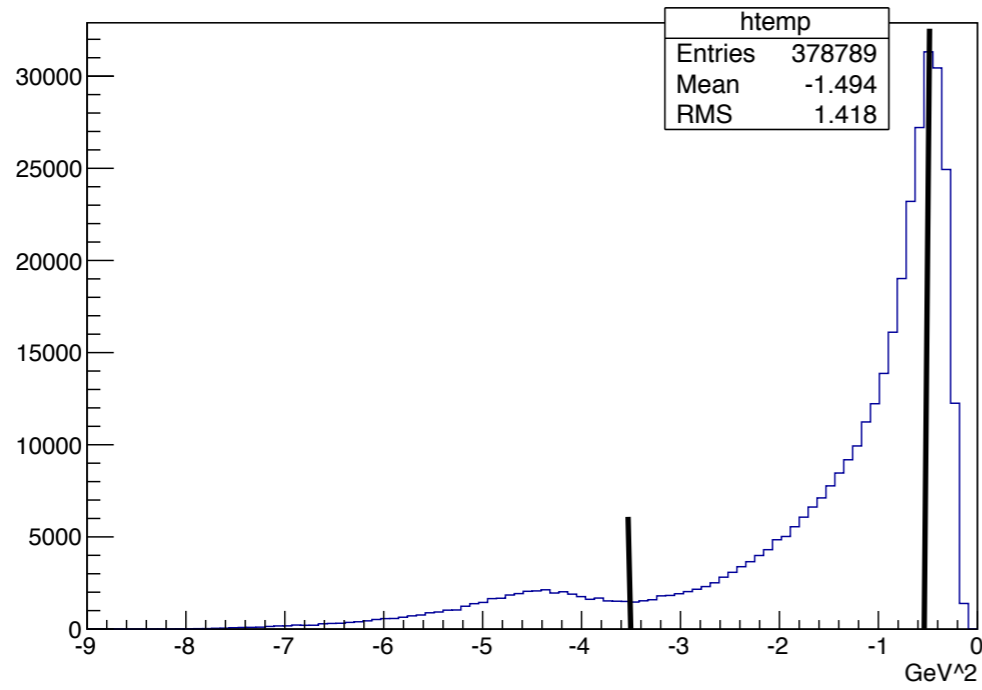
Then we impose Pull Prob > 0.15.



Proton Missing Mass (energy, time and pull prob cuts)

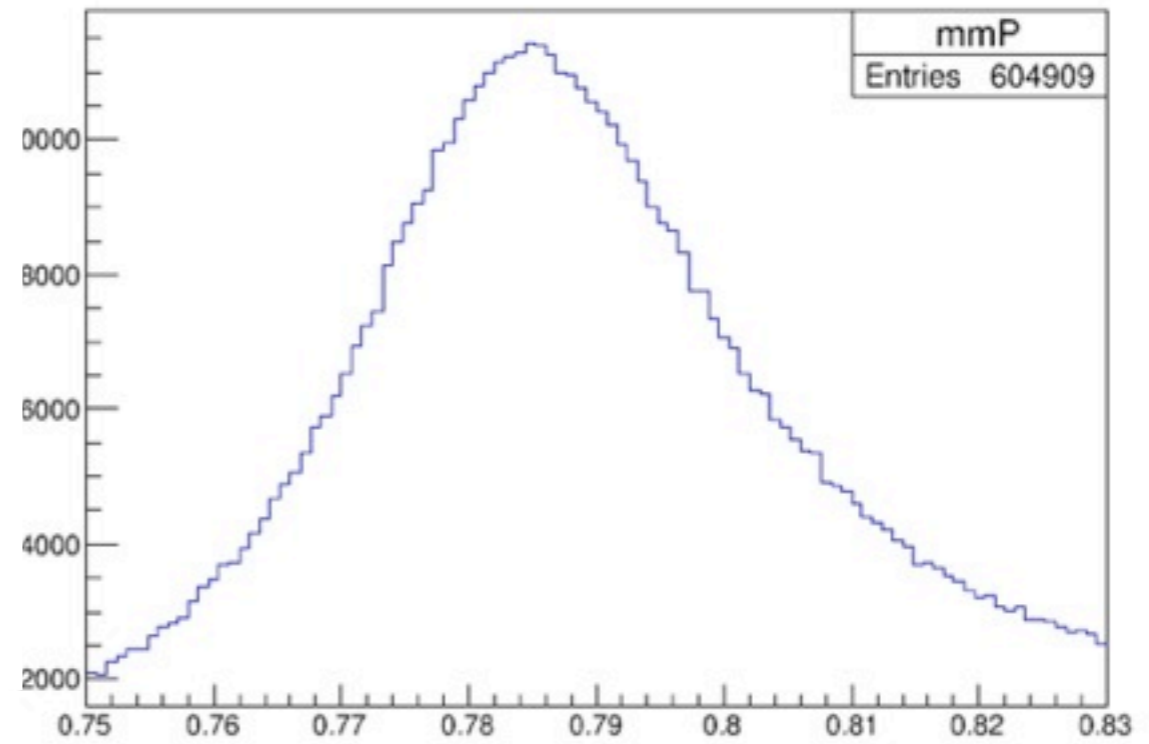


Mandelstam_tprime

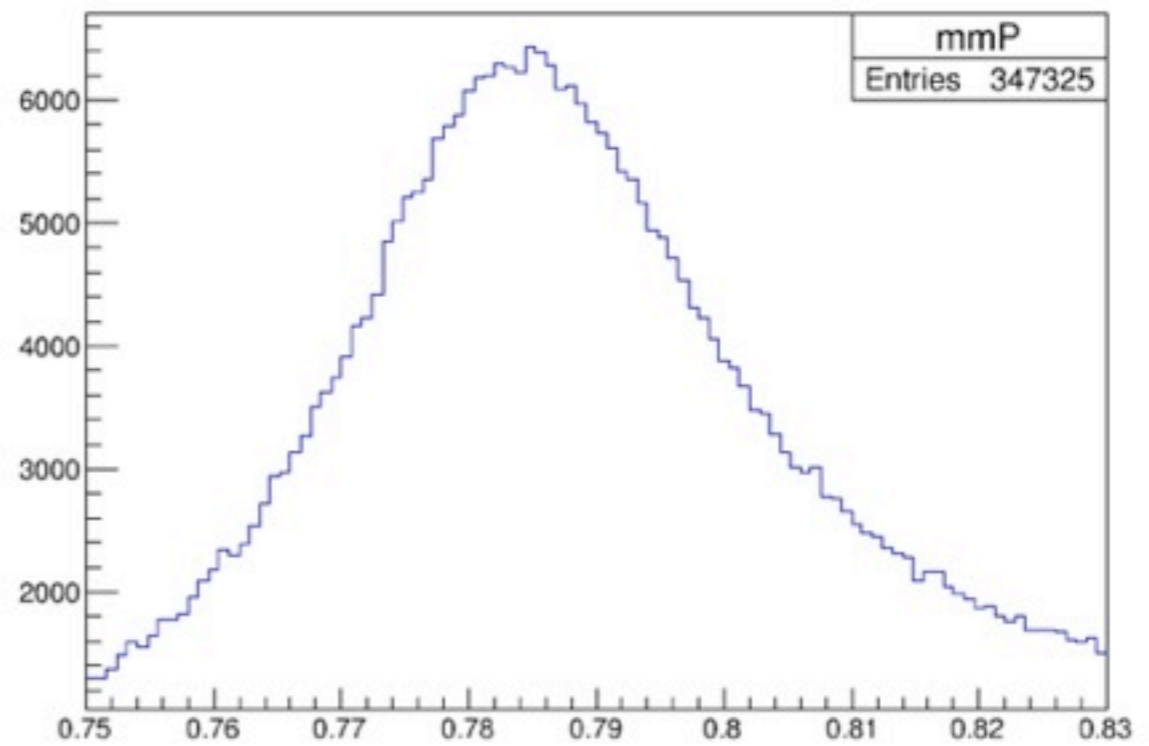


Also, a cut on the Mandelstam variable t' was performed: $3.5 < t' < -0.6$

Proton Missing Mass (energy, time and all pull cuts)



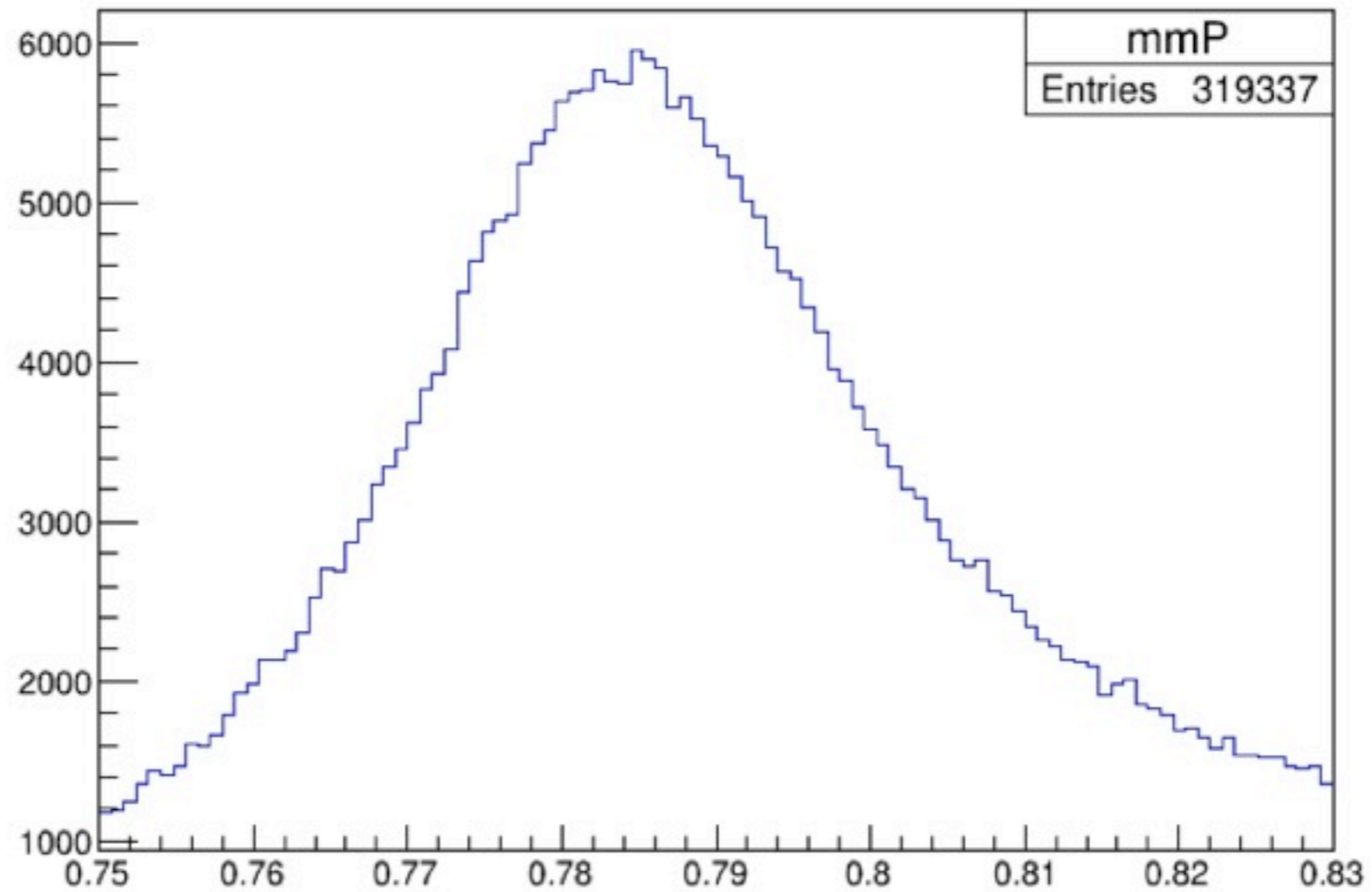
Proton Missing Mass (energy, time, all pull and t cuts)

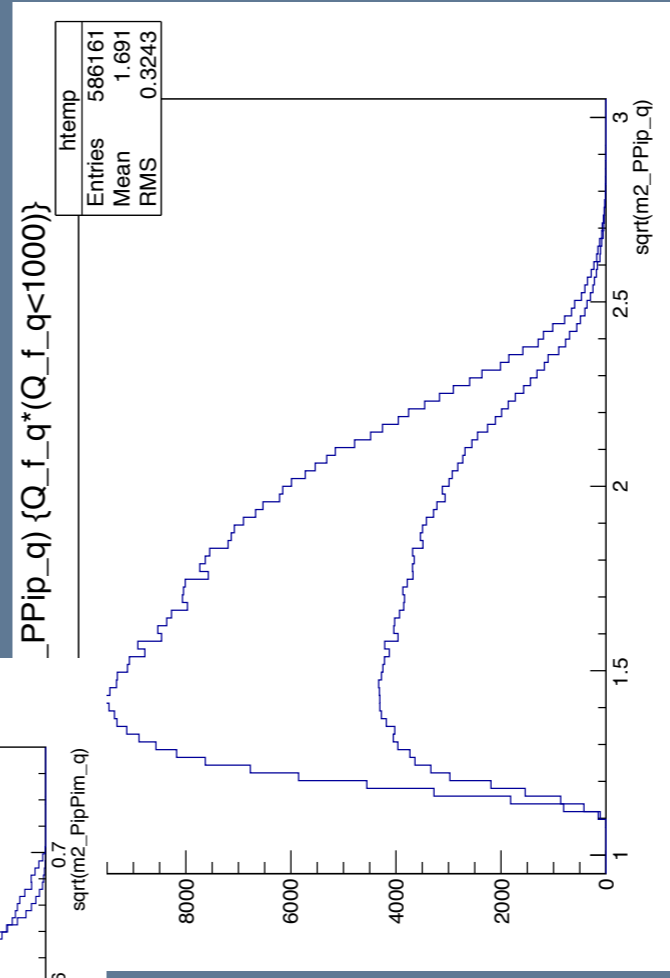
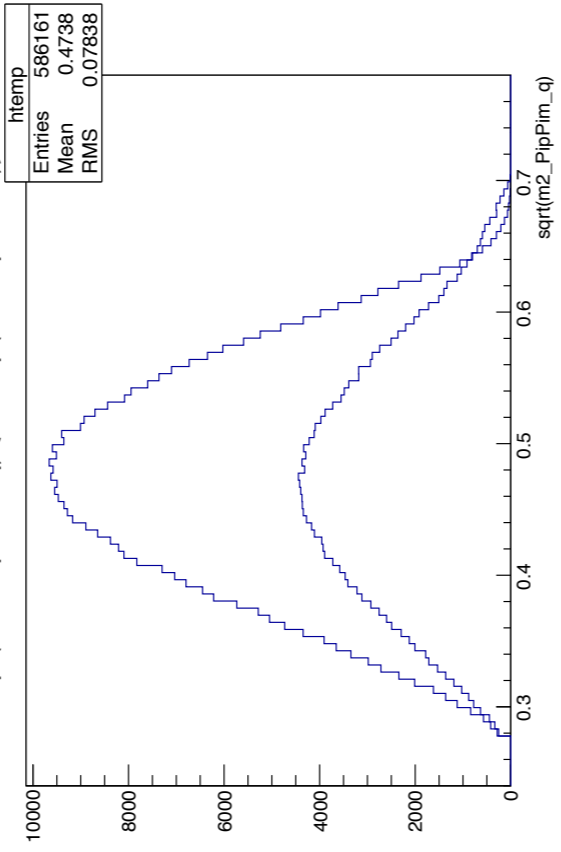
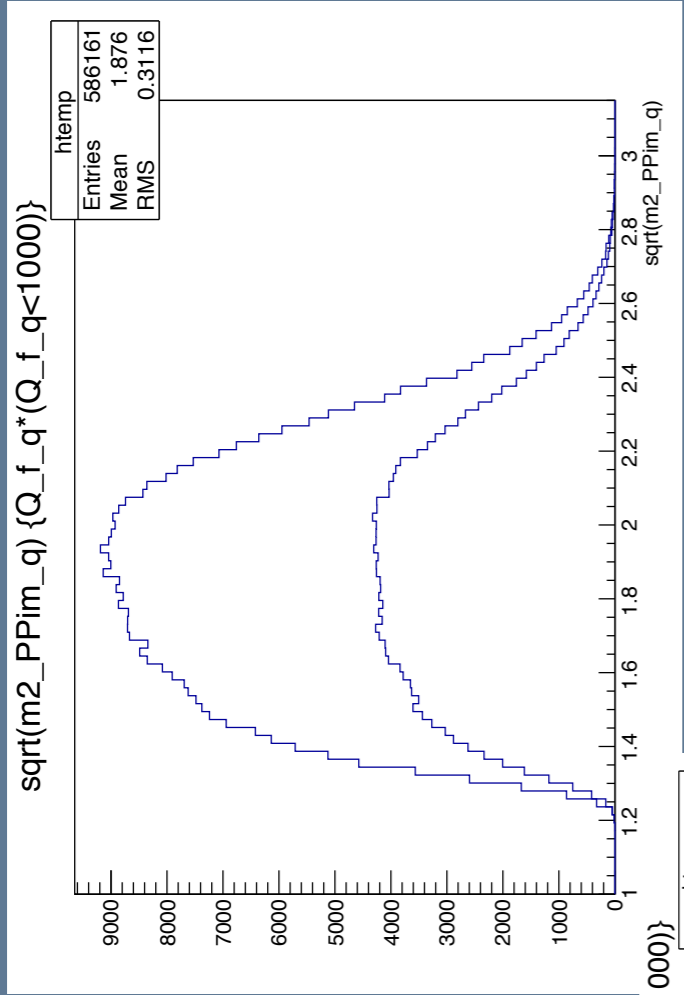


Finally, the fiducial and TOF knockout cuts are considered.

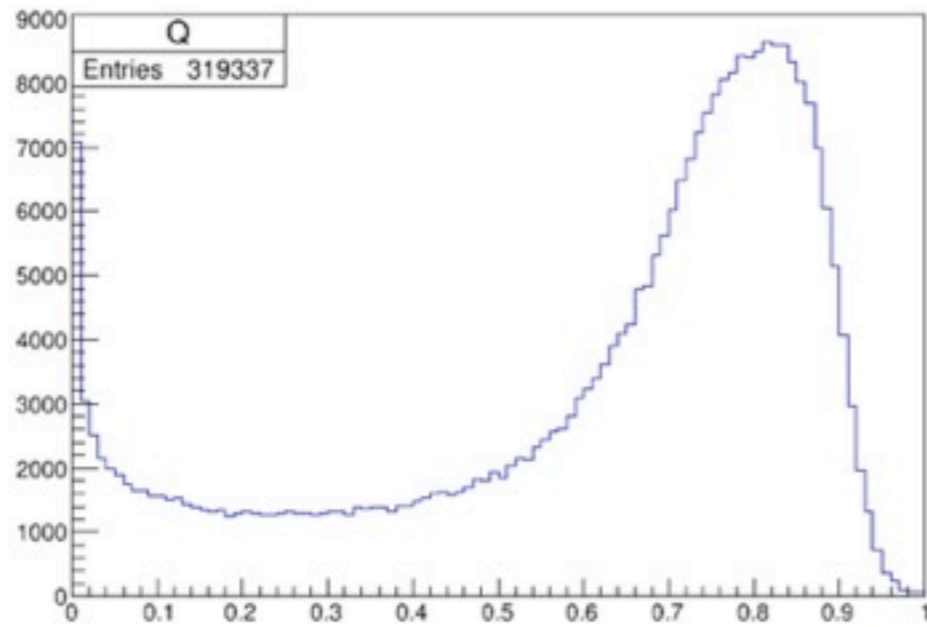
(standard g12 - see general note)

Proton Missing Mass (all cuts)





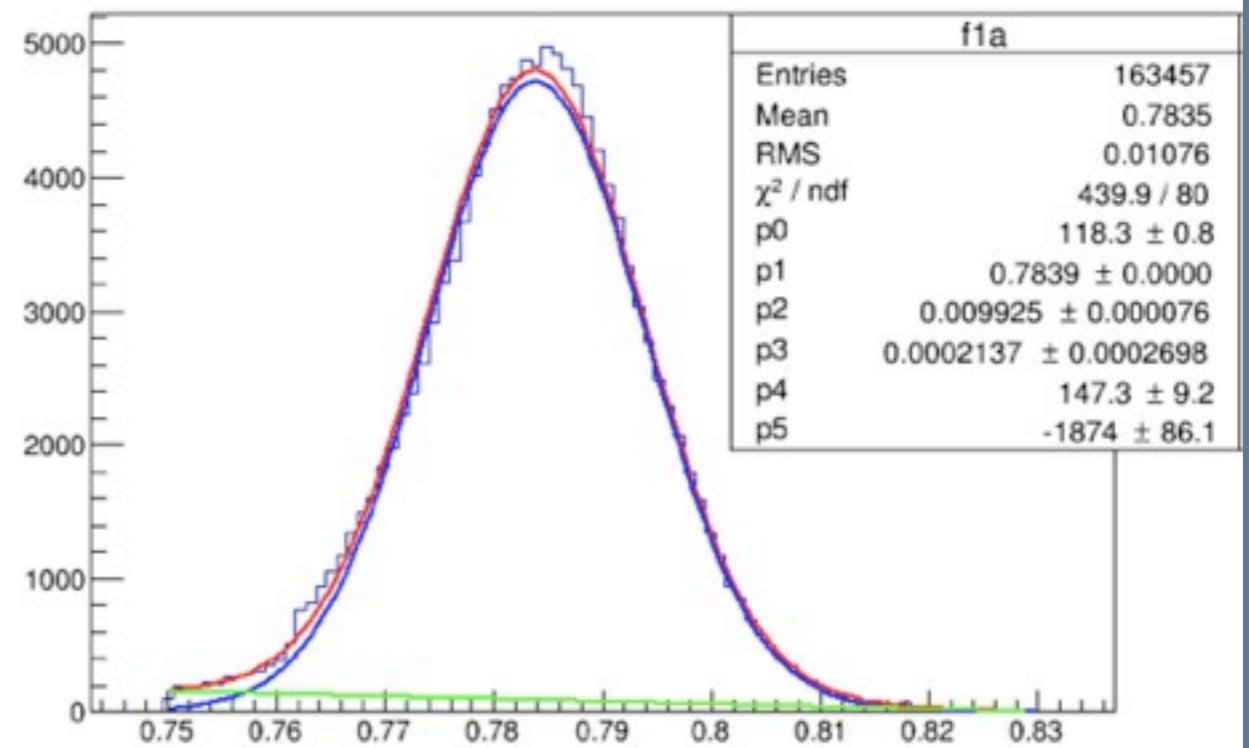
Q Factor



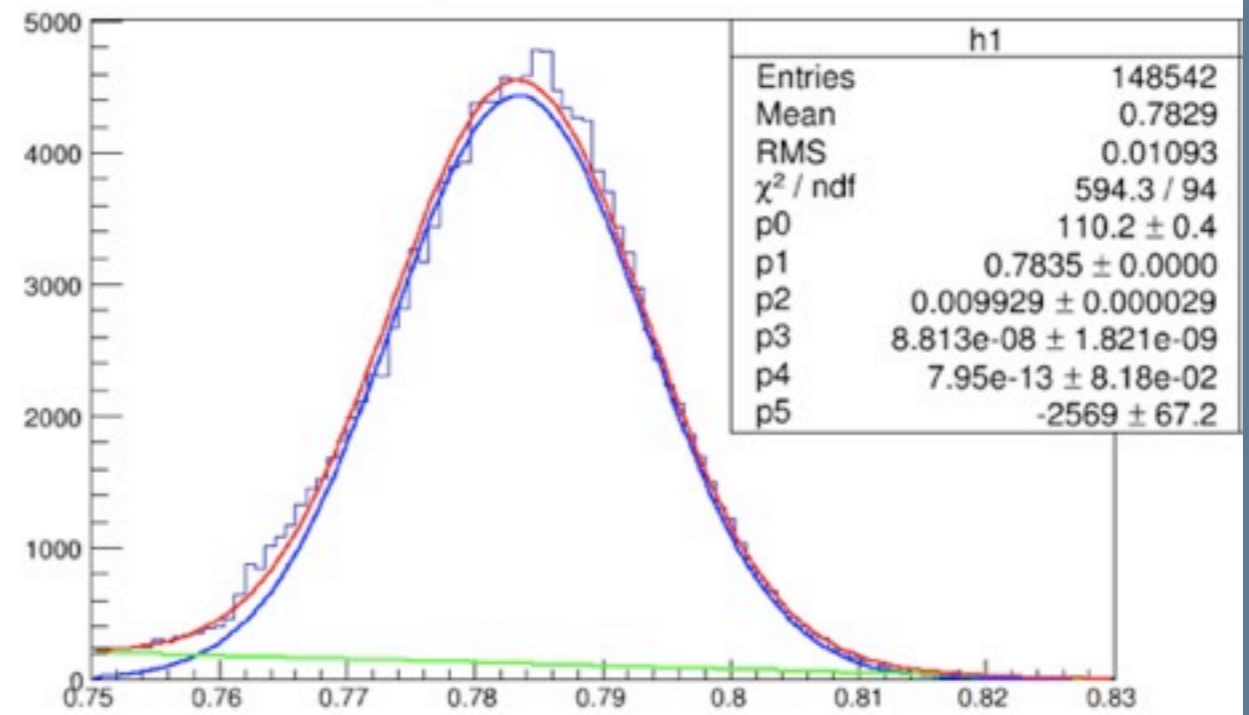
$$Q_f > 0.7$$

(Qf as defined in Williams Thesis)

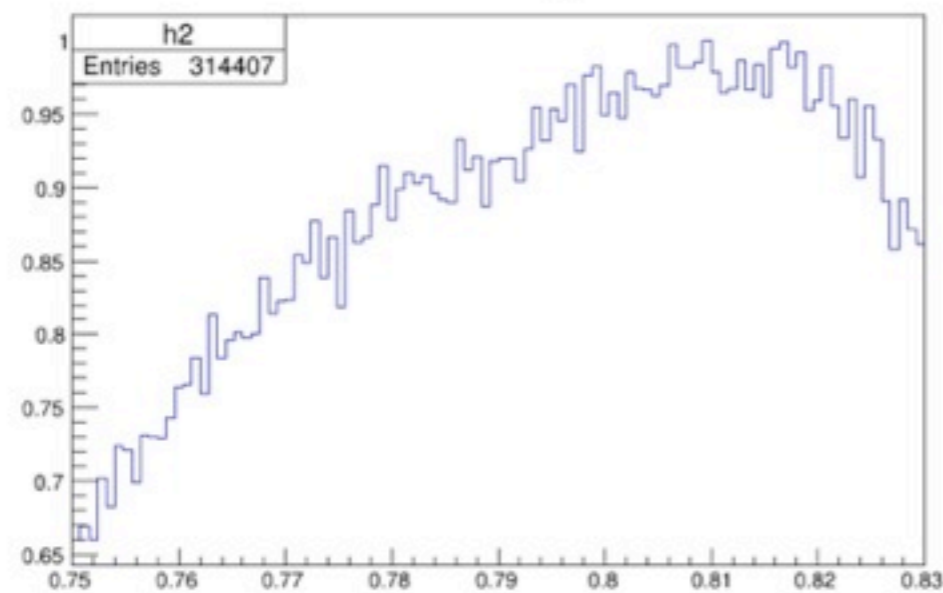
Proton Missing Mass



Acceptance Corrected Data



Normalized Acceptance



Dispersive Analysis of $\omega/\phi \rightarrow 3\pi, \pi\gamma^*$

I.V. Danilkin,^{1,*} C. Fernández-Ramírez,¹ P. Guo,^{2,3} V. Mathieu,^{2,3} D. Schott,⁴ M. Shi,^{1,5} and A. P. Szczepaniak^{1,2,3}
 (Joint Physics Analysis Center)

Dispersive analysis of $\omega \rightarrow 3\pi$ and $\phi \rightarrow 3\pi$ decays

F. Niecknig, B. Kubis, and S. P. Schneider, Eur.Phys.J. C72, 2014 (2012).

Franz Niecknig, Bastian Kubis, Sebastian P. Schneider

Helmholtz-Institut für Strahlen- und Kernphysik (Theorie) and Bethe Center for Theoretical Physics, Universität Bonn, D-53115 Bonn, Germany

$$\frac{d^2\Gamma}{ds dt} = \frac{1}{(2\pi)^3} \frac{1}{32M^3} \frac{1}{3} P(s, t) |F(s, t, u)|^2$$

$$P(s,t) = stu - m_\pi^2 (M^2 - m_\pi^2)^2.$$

$$F(s) = \Omega(s) \left(\frac{1}{\pi} \int_{s_\pi}^{s_i} ds' \frac{\rho(s') t^*(s') \hat{F}(s')}{\Omega^*(s') s' - s} + \Sigma(s) \right). \tag{31}$$

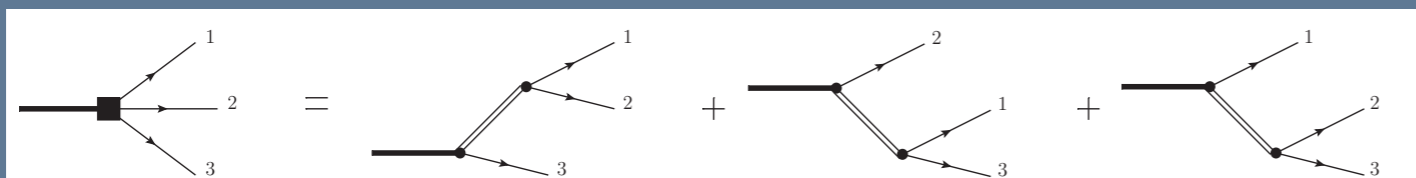


FIG. 1: Isobar decomposition.

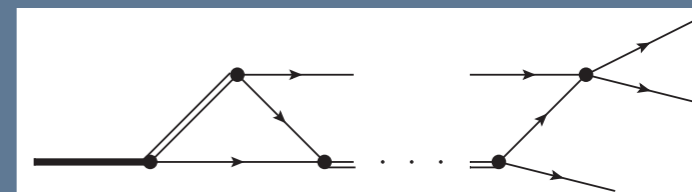


FIG. 2: Crossed channel rescattering effects.

$$\Sigma(s) = \sum_{i=0}^{\infty} a_i \omega^i(s)$$

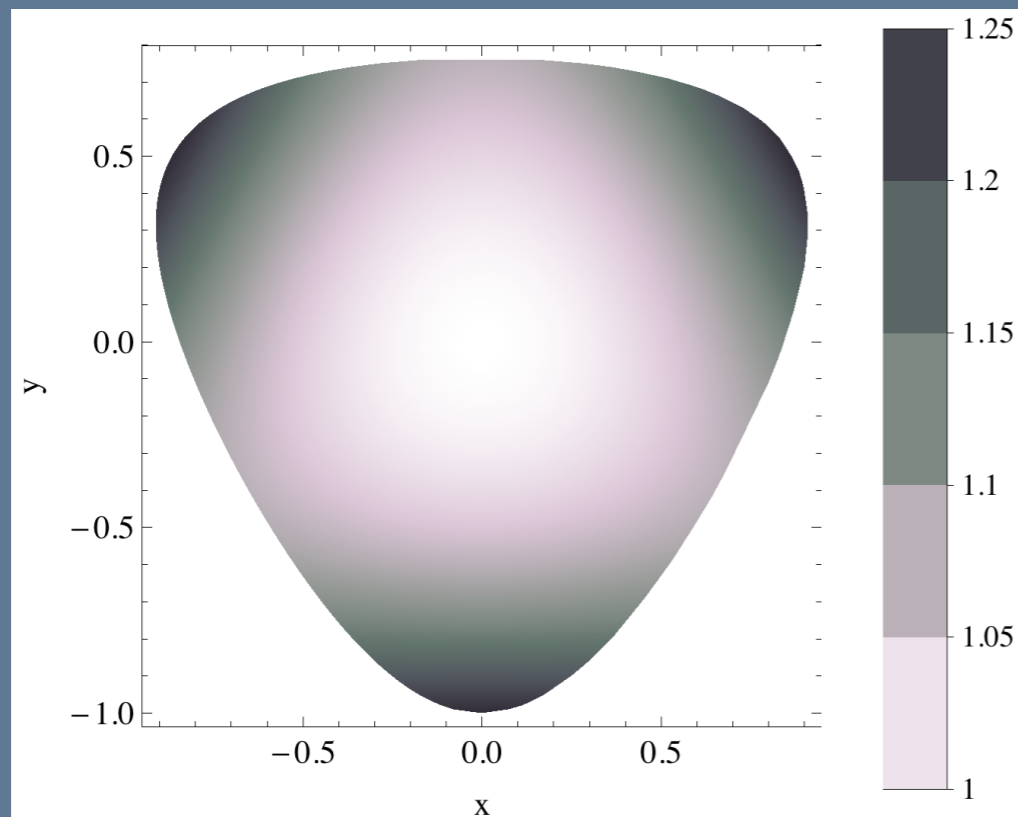
The variable

$$\omega(s) = \frac{\sqrt{s_i - s_E} - \sqrt{s_i - s}}{\sqrt{s_i - s_E} + \sqrt{s_i - s}}$$

$$F(s,t,u) = F_0(s,t) + a_0 F_1(s,t)$$

(code available in the web at: <http://cgl.soic.indiana.edu/jpac/>)

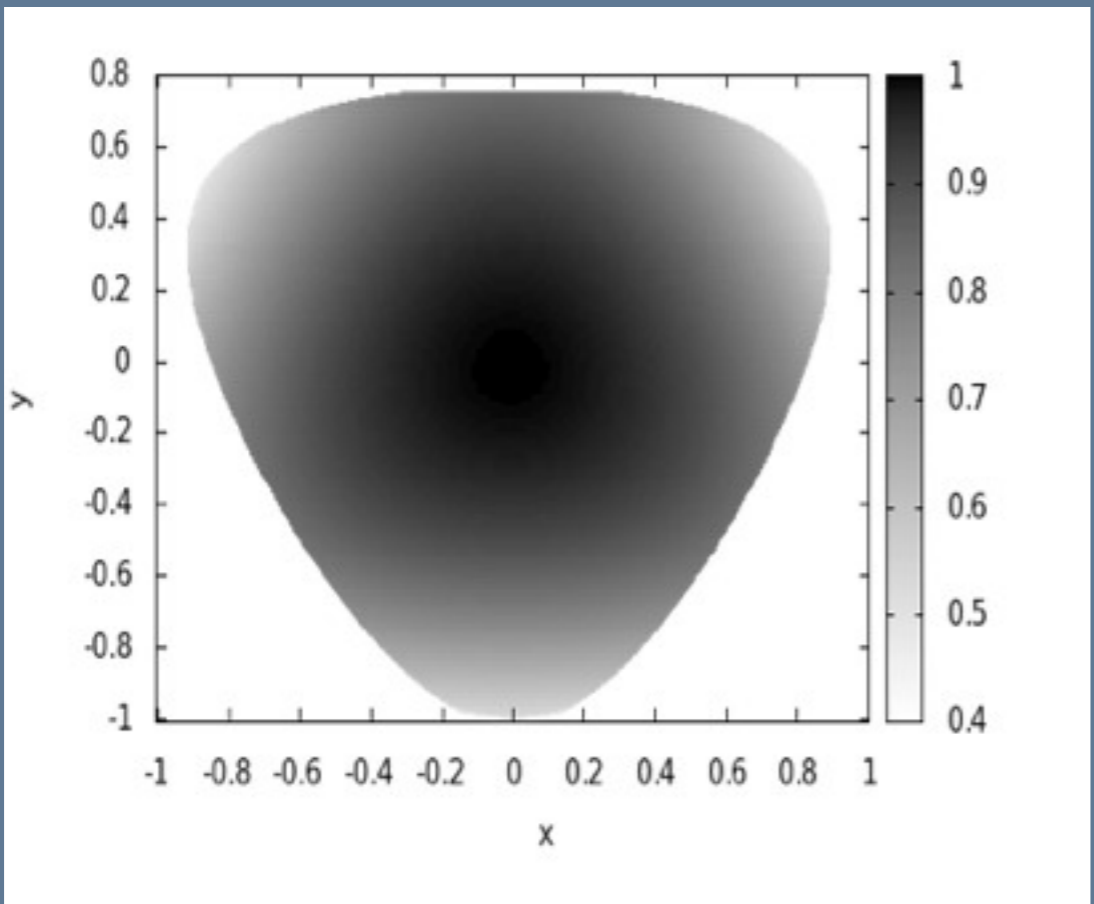
Dalitz Analysis in x vs y



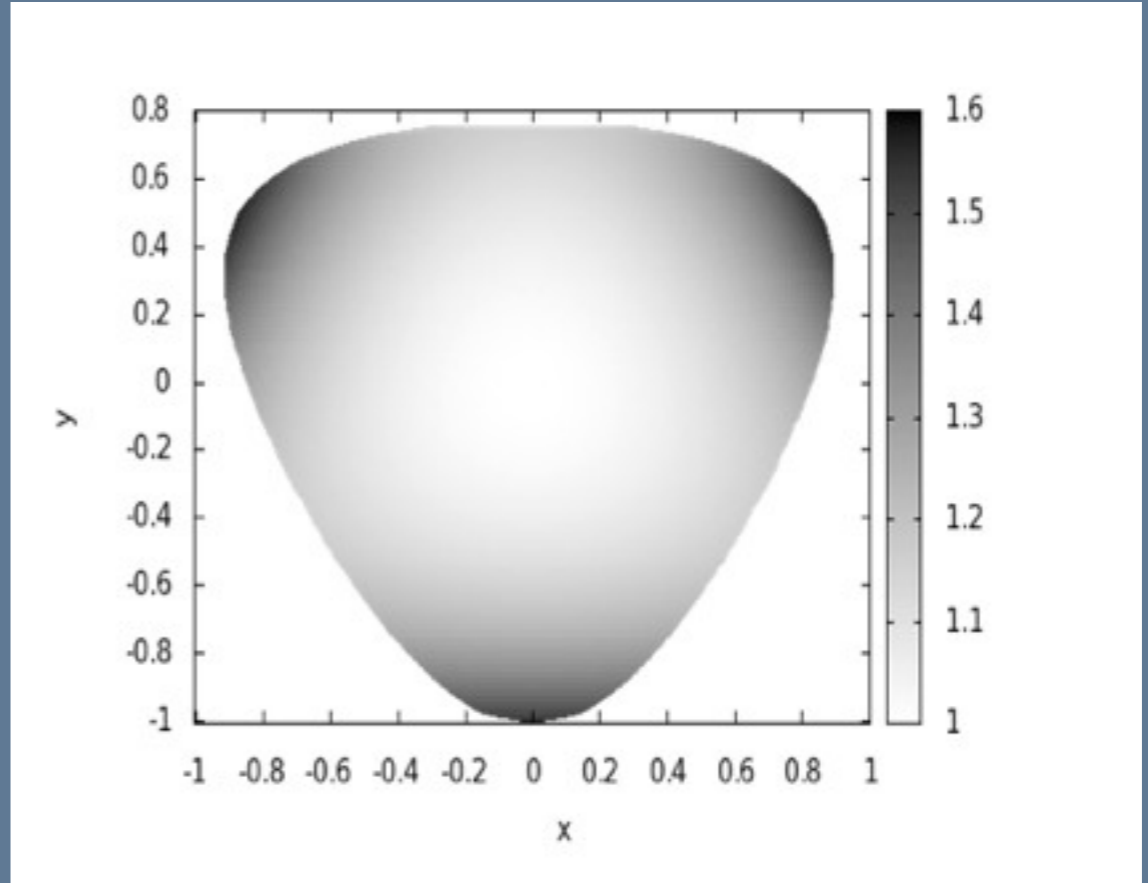
$$a_0 = 0$$

$$x = \frac{\sqrt{3}}{Q} (T_1 - T_2) = \frac{\sqrt{3}(t - u)}{2M(M - 3m_\pi)},$$

$$y = \frac{3T_3}{Q} - 1 = \frac{3(s_c - s)}{2M(M - 3m_\pi)}.$$

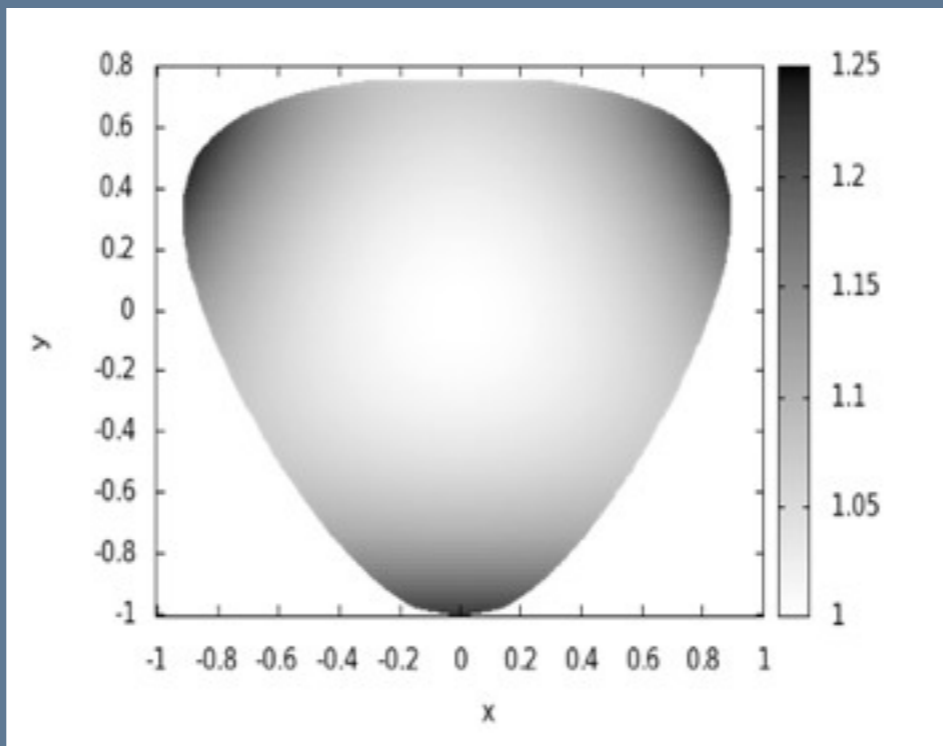


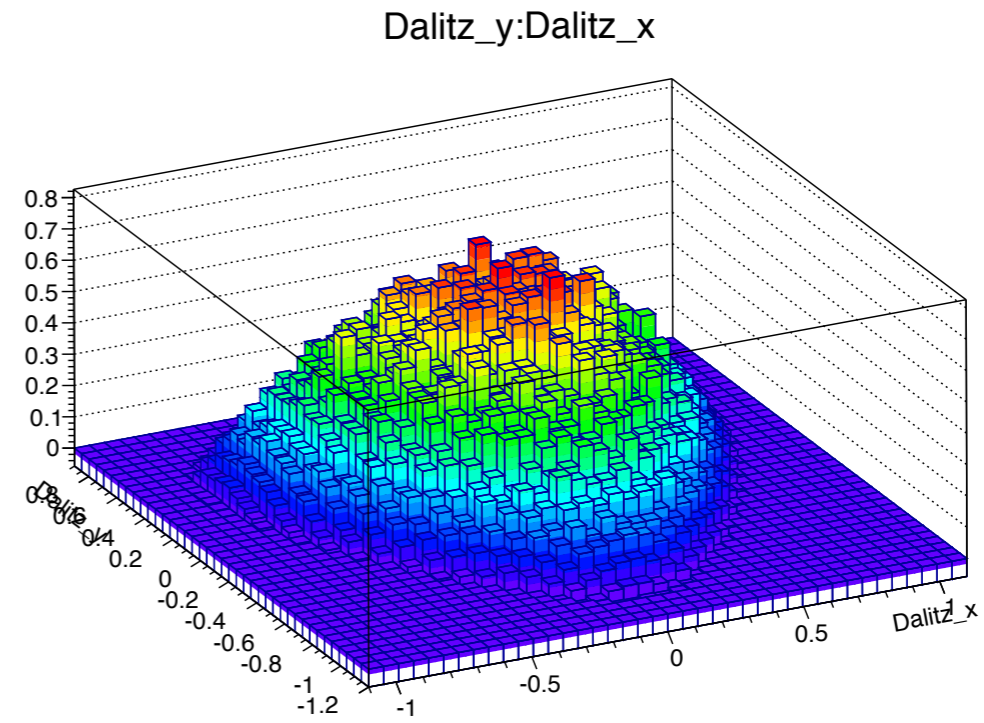
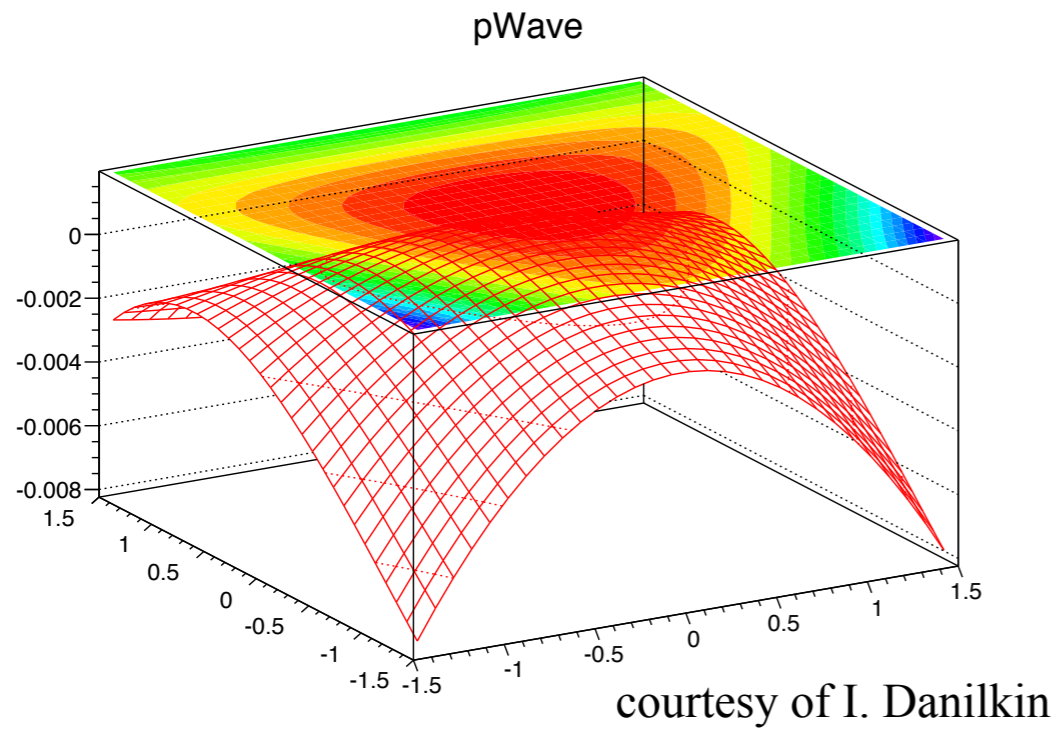
$a_o = -10$



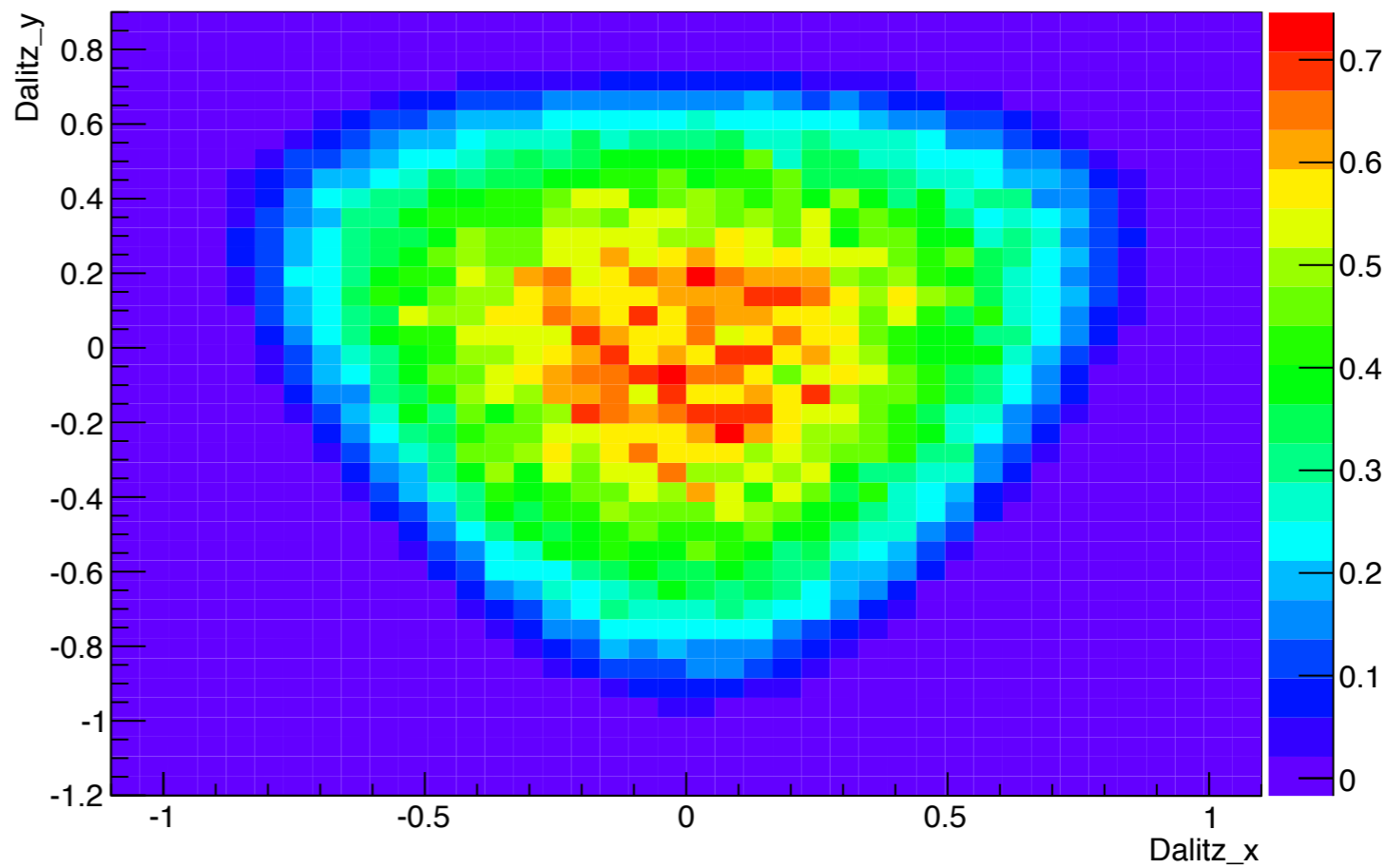
$a_o = 10$

$a_o = 0$



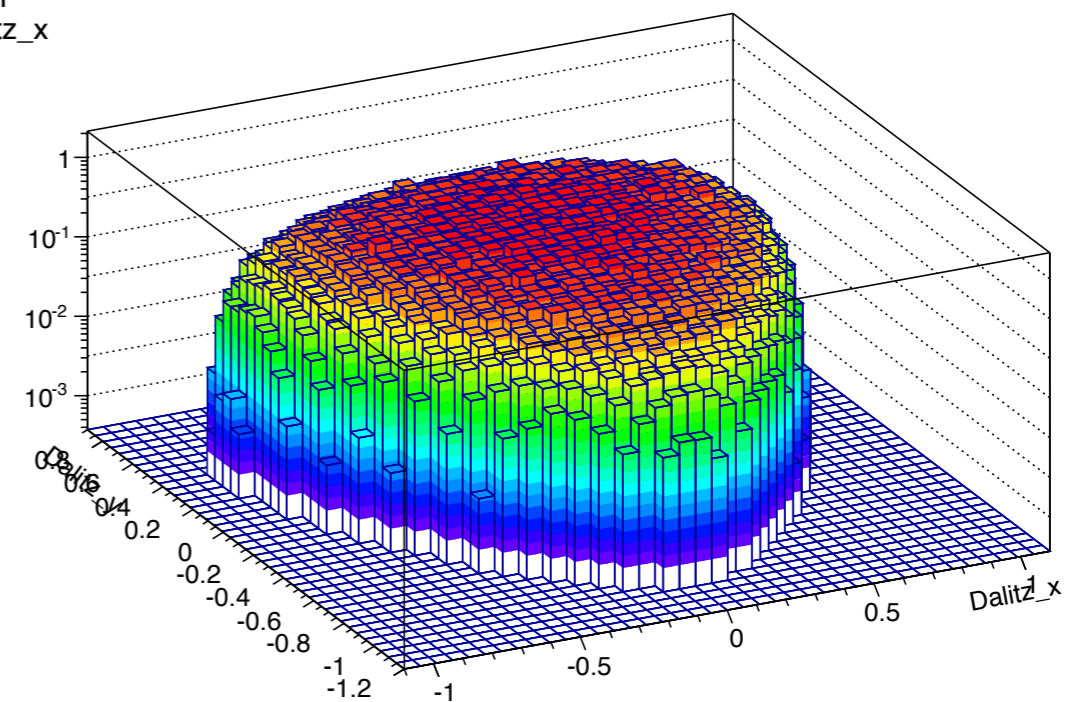
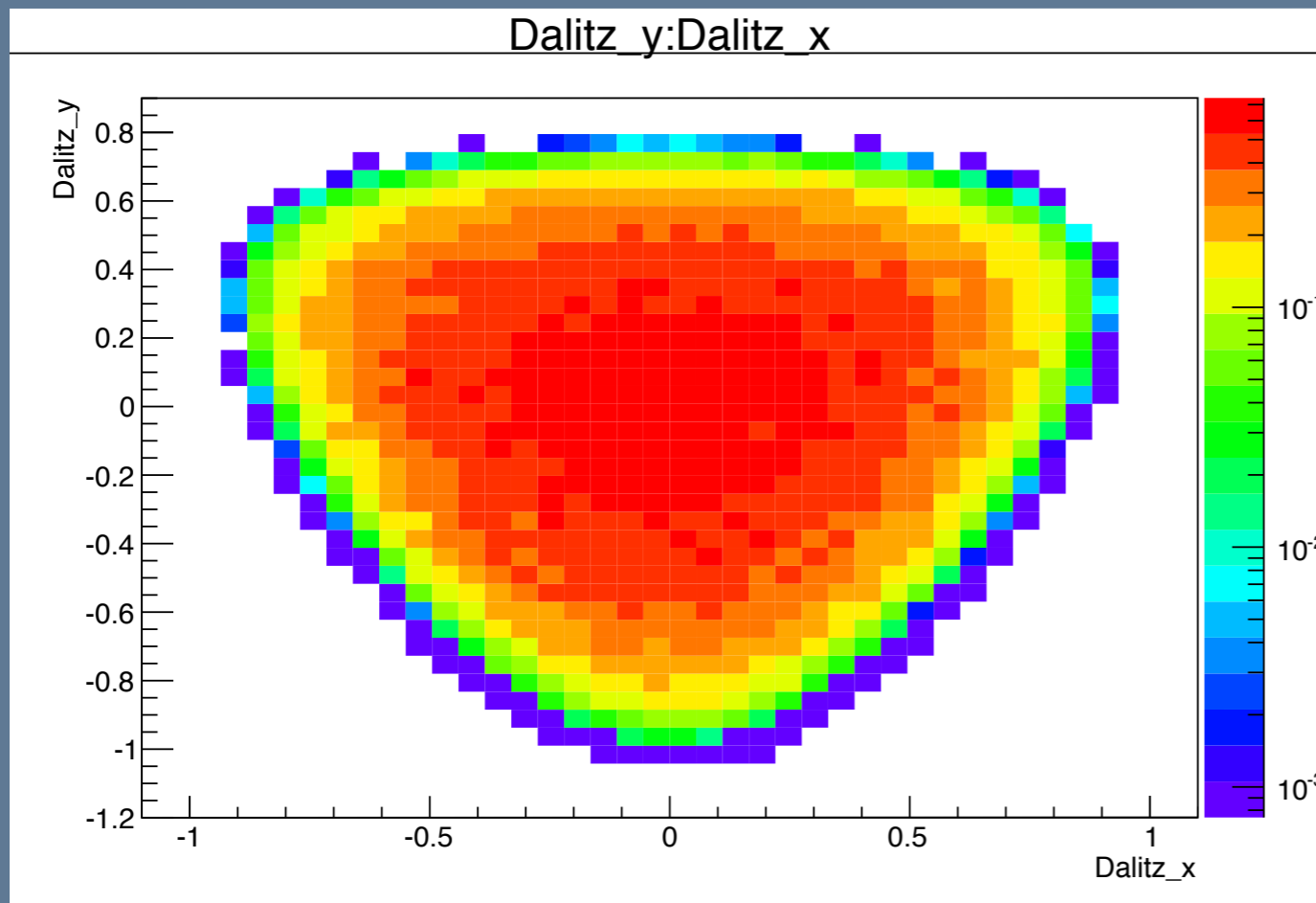


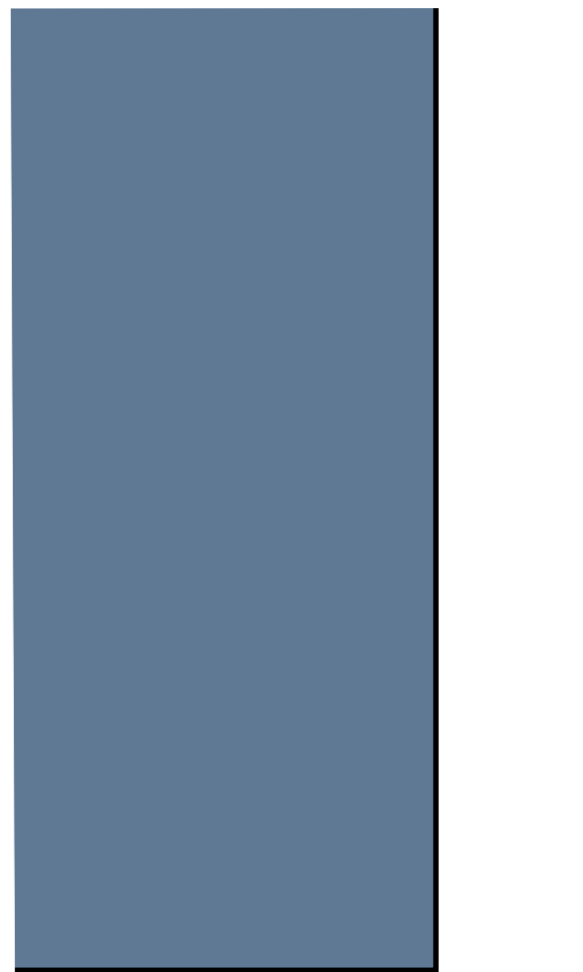
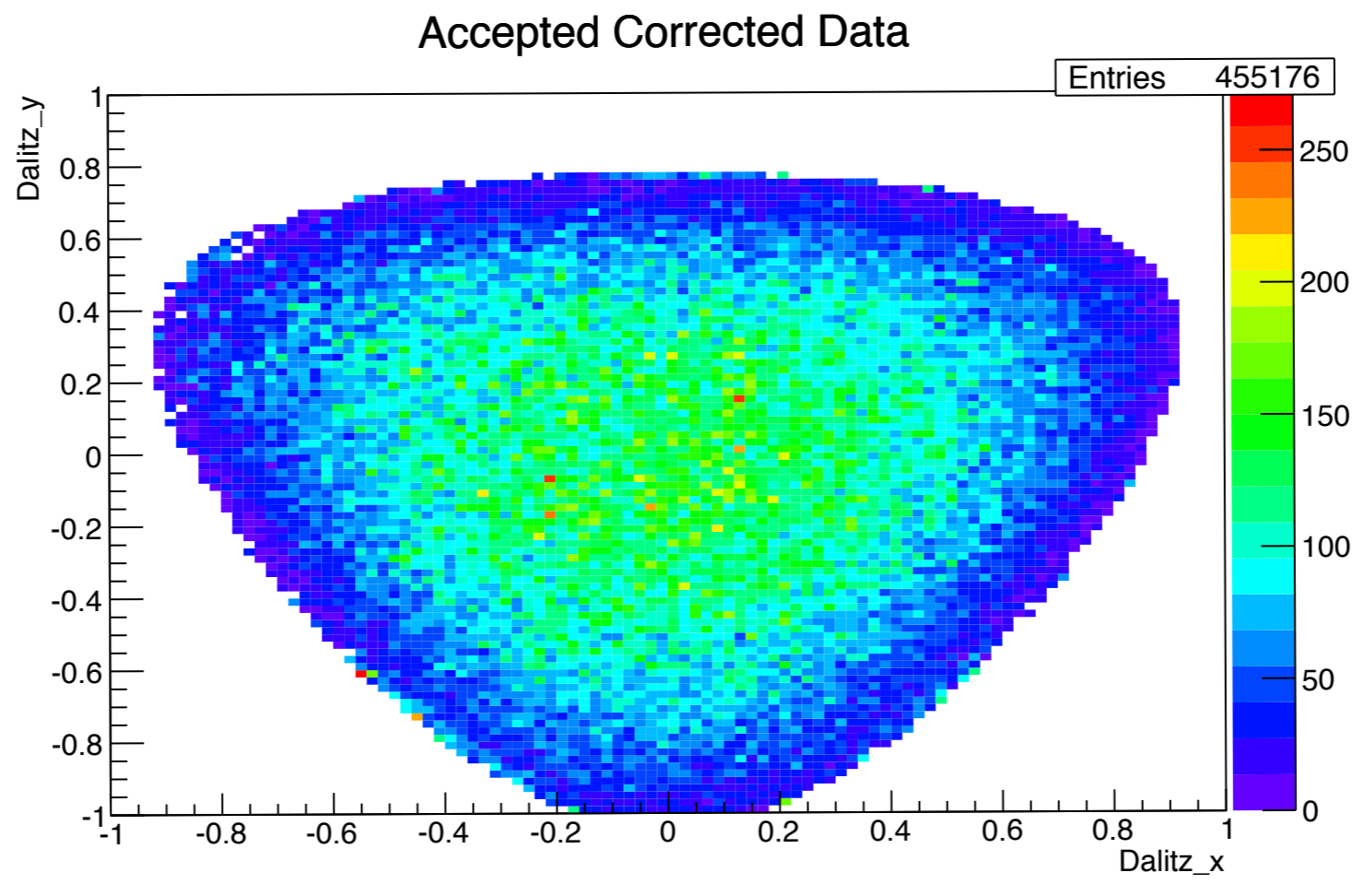
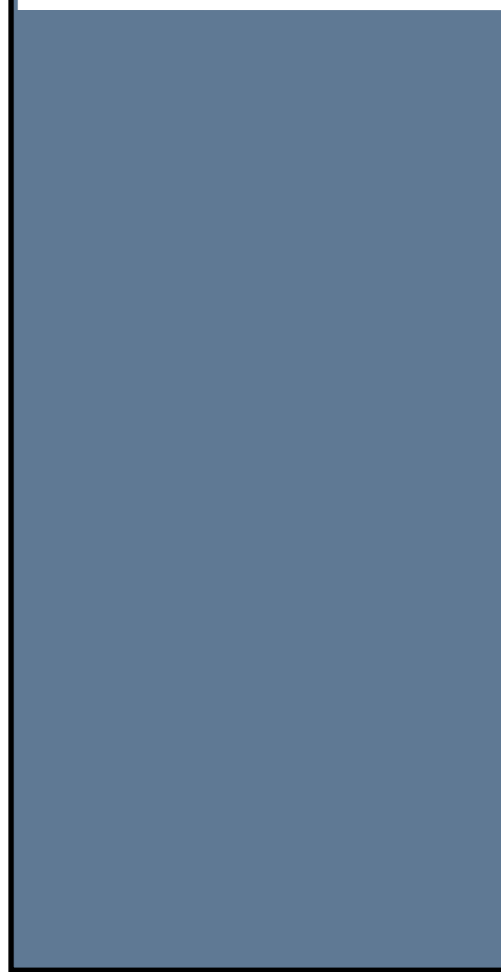
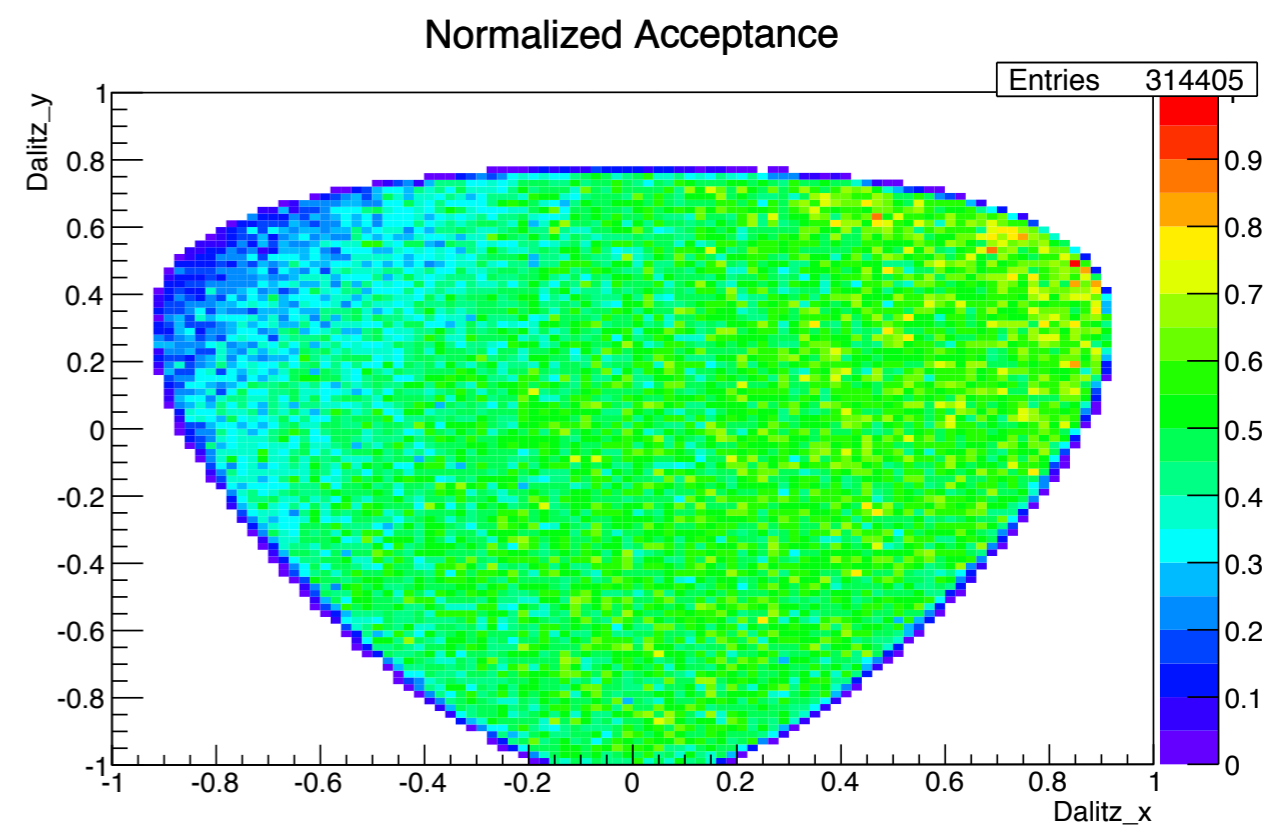
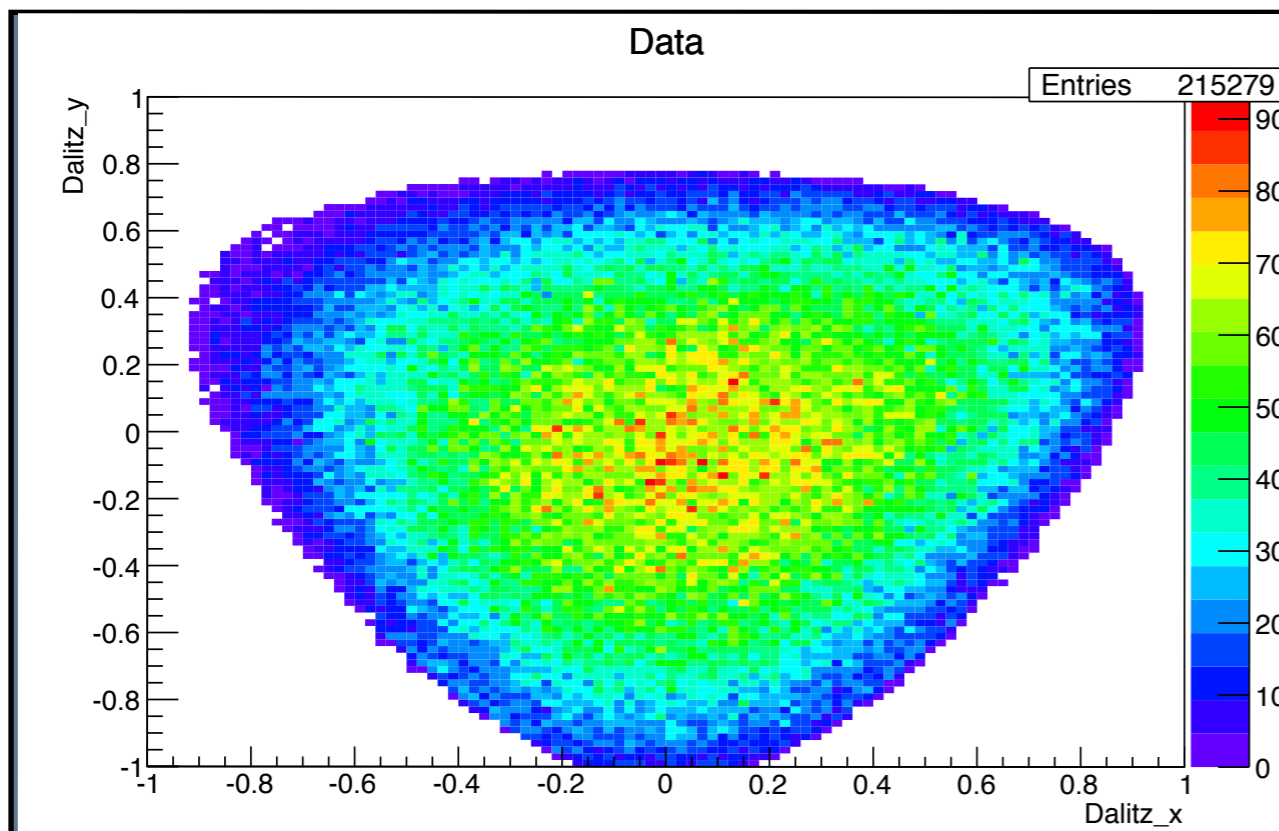
Dalitz_y:Dalitz_x



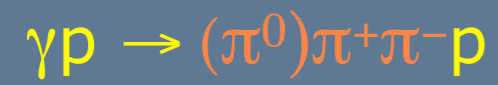
$P(s,t)$ factor

P-factor on log-z scale

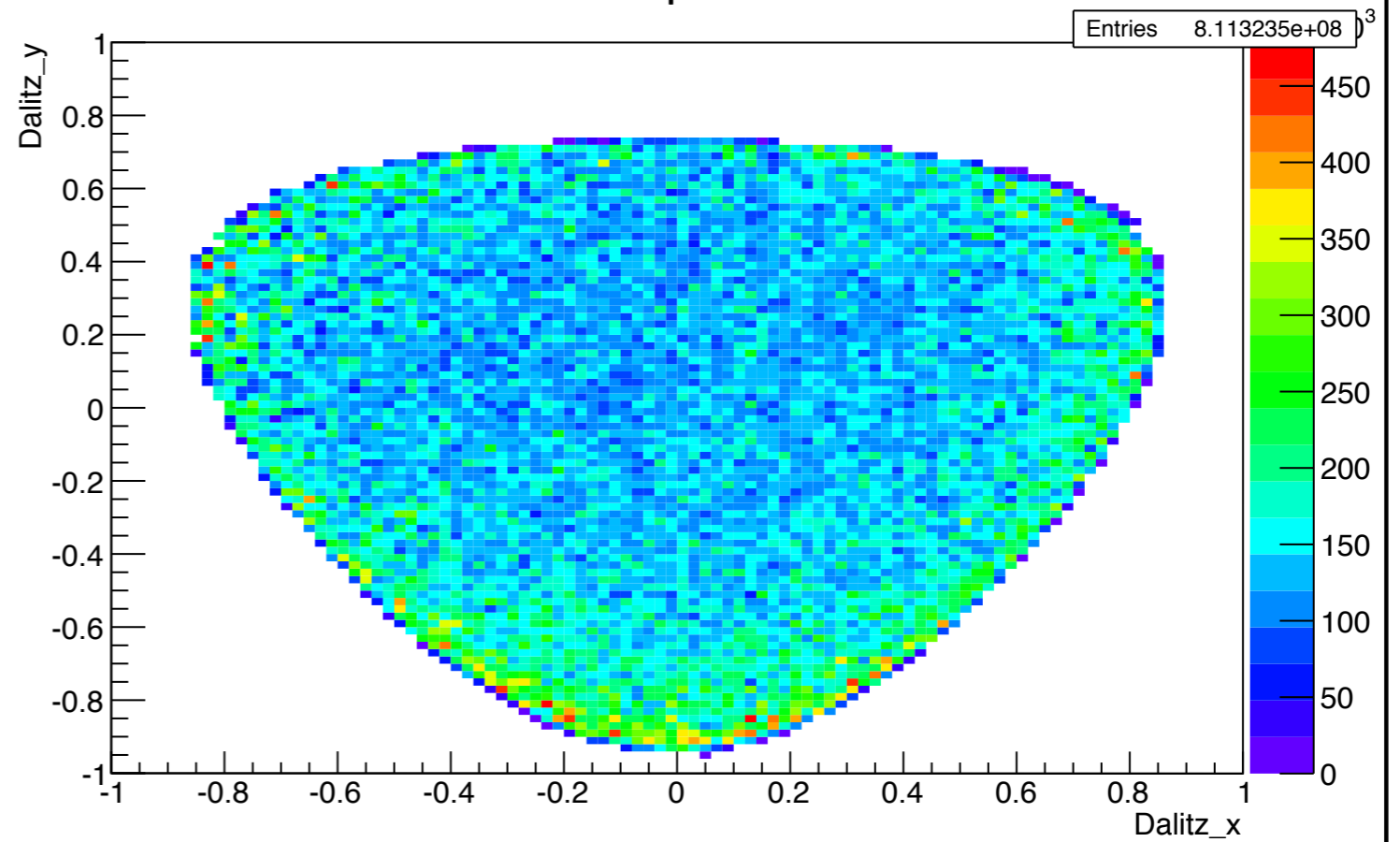




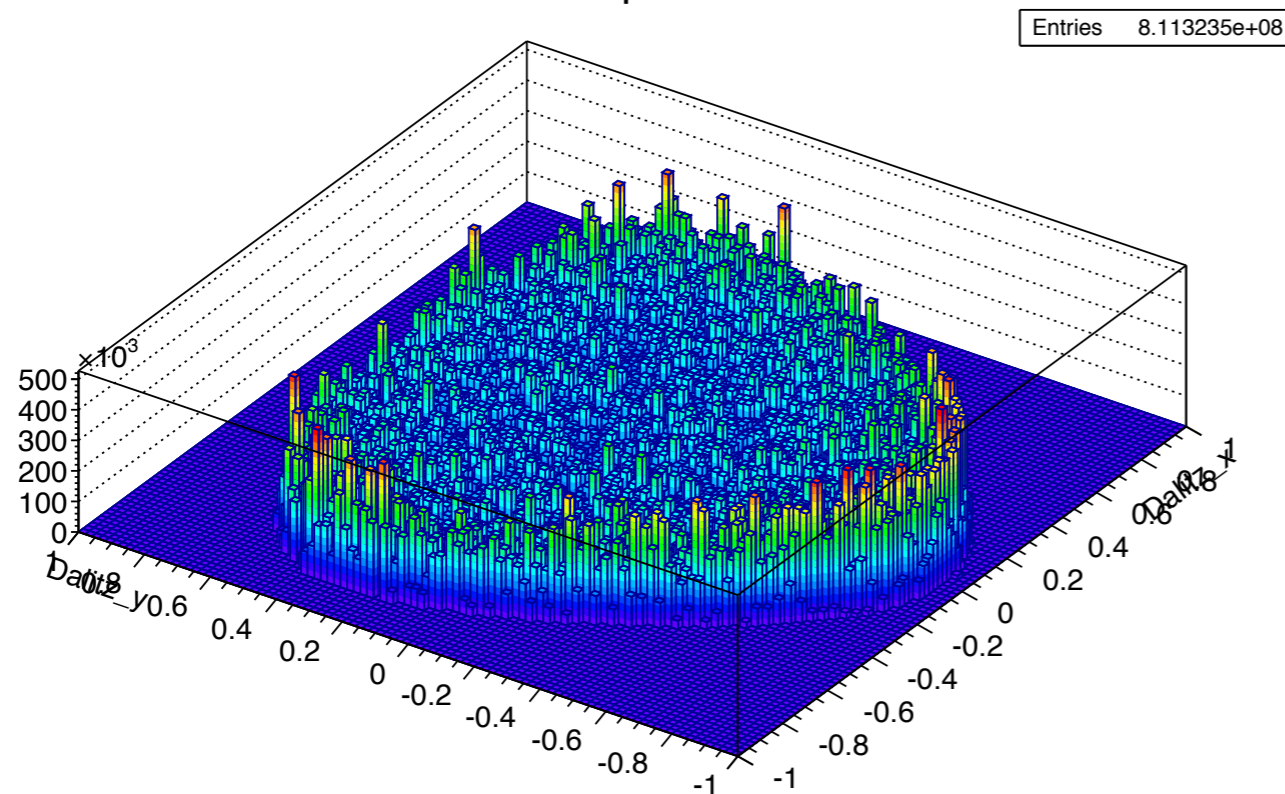
g12 Data



P-Normalized Accepted Corrected Data



P-Normalized Accepted Corrected Data



Fit parametric model to data using the PyPWA framework - (generalShell module) see <https://pypwa.jlab.org>

Performing an Unbinned Extended Likelihood fit :

$$-\ln \mathcal{L} = - \sum_{i=1}^N Q_i \ln [I(\vec{x}_i, \vec{a})] + \frac{1}{N_g} \sum_{i=1}^{N_a} I(\vec{x}_i, \vec{a})$$

Minimized the negative log-likelihood on the model parameters

In our case:

$$I(sD, tD, \theta_{Adair}, \phi_{Adair}, \text{parameters}) = \text{production} * \text{decay}$$

3 Fitted Function

$$I(sD, tD, uD, \theta, \phi, \mathbf{A1}, \mathbf{A2}, \mathbf{A3}, \mathbf{A4}, \mathbf{A5}) = \mathbf{A1} * W(\theta, \phi, \mathbf{A2}, \mathbf{A3}, \mathbf{A4}) * P(sD, tD, uD) * [F(sD, tD, uD, \mathbf{A5})]^2 \quad (15)$$

where W is the Schilling et al. spin density matrix (no-polarization):

$$W(\theta, \phi, \mathbf{A2}, \mathbf{A3}, \mathbf{A4}) = \quad (16)$$

$$\frac{3}{4\pi} [0.5 * (1 - \mathbf{A2}) + 0.5 * (3 * \mathbf{A2} - 1) \cos^2(\theta) - \sqrt{2} * \mathbf{A3} * \sin 2(\theta) \cos(\phi) - \mathbf{A4} * \sin^2(\theta) \cos 2(\phi)] \quad (17)$$

and θ, ϕ are Adair's angles. P is a kinematic factor given by:

$$P(sD, tD, uD) = sD * tD * uD - m_\pi^2 [M^2 - m_\pi^2]^2 \quad (18)$$

where sD, tD, uD are the Mandelstam variables of the decay such that:

$$sD = (p_X - p_{\pi^+})^2, \quad tD = (p_X - p_{\pi^-})^2 \quad \text{and} \quad uD = (p_X - p_{\pi^0})^2.$$

and $p_X = p_{\pi^+} + p_{\pi^-} + p_{\pi^0}$, M is the mass of the three pion system and m_π the mass of the pion (plus).

$F(sD, tD, uD, \mathbf{A5})$ is Igor Danilkin et al. amplitude given for a call to his fortran code.

MC was generated with a t -slope of 3 GeV^{-2} to match data low t distributions. - all t are included in current fits. (future analysis of t dependence is planned).

Preliminary Fits results

All	A2	A3	A4	A5
3.5-4.0	0.315	-0.016	-0.021	-12.53
4.0-4.5	0.315	-0.016	-0,021	-12.78
4.5-5.0	0.315	-0.016	-0.021	-12.82
5.0-5.5	0.191	0.018	-0.007	-12.8

nonF	A2	A3	A4
3.5-4.0	0.27	-0.018	-0.023
4.0-4.5	0.28	-0.026	-0,068
4.5-5.0	0.29	-0.022	-0.039
5.0-5.5	0.31	0.016	-0.021

F	A2	A3	A4	A5
3.5-4.0	0.194	-0.016	-0.007	0 (fixed)
3.5-4.0	0.299	-0.018	-0,08	1 (fixed)
3.5-4.0	0.191	-0.018	-0.007	-12.8 (fixed normalization)

Still studying stability of fits: **next steps**

(ii) Circular polarization of helicity $\lambda_\gamma = \pm 1$:

$$W^\pm(\cos\theta, \phi) = W^0(\cos\theta, \phi) \pm P_\gamma W^3(\cos\theta, \phi) .$$

$$W^0(\cos\theta, \phi) = \frac{3}{4\pi} \left(\frac{1}{2}(1 - \rho_{00}^0) + \frac{1}{2}(3\rho_{00}^0 - 1) \cos^2\theta - \sqrt{2} \operatorname{Re}\rho_{10}^0 \sin 2\theta \cos\phi - \rho_{1-1}^0 \sin^2\theta \cos 2\phi \right) ,$$

Schilling et al

$$W^3(\cos\theta, \phi) = \frac{3}{4\pi} \left(+\sqrt{2} \operatorname{Im}\rho_{10}^3 \sin 2\theta \sin\phi + \operatorname{Im}\rho_{1-1}^3 \sin^2\theta \sin 2\phi \right)$$

Parametrization of Dalitz intensity through:

$$x = \sqrt{z} \cos\vartheta, \quad y = \sqrt{z} \sin\vartheta, \quad (39)$$

and fit the following polynomial expansion

$$|F_{par}(z, \vartheta)|^2 = |N|^2 \left(1 + 2\alpha z + 2\beta z^{3/2} \sin(3\vartheta) + 2\gamma z^2 + 2\delta z^{5/2} \sin(3\vartheta) + \mathcal{O}(z^3) \right) \quad (40)$$

Kubis et al

Preliminary Results and Things still to be done

- $\omega \rightarrow \pi^+\pi^-(\pi^0)$ events for $3.6 < E_{\text{photon}} < 5.4$ GeV have been extracted given a mass for the ω of 783.5 MeV and width of 9.92 MeV (PDG: 782.65,8.49). Sample with very small background.
- Comparison with theory has started:
 - Data seems dominated by P-wave (as expected).
 - The extra-terms of the three-body decay models are important at the edges of the Dalitz plots where acceptance/statistics are very limited.

Next steps:

- Study Fit stability.
- Introduce other parametrization (and polarization,...)
- Study Energy and t dependancies.