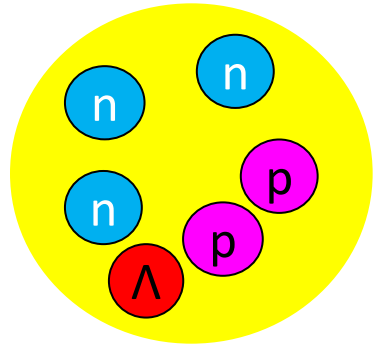


# P-shell to heavier $\Lambda$ hypernuclei

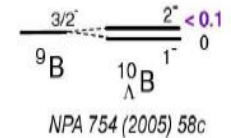
Emiko Hiyama (Tohoku/RIKEN)

Since 1998

# Hypernuclear $\gamma$ -ray data (2019)

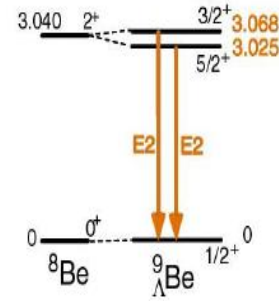


$^{10}\text{B} (K^-, \pi^+ \gamma)$  BNL E930('01)



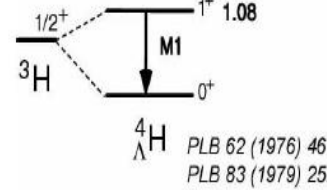
NPA 754 (2005) 58c

$^9\text{Be} (K^-, \pi^+ \gamma)$  BNL E930('98)



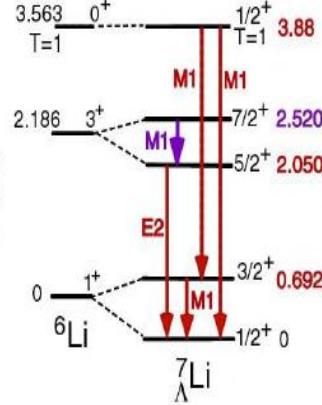
PRL 88 (2002) 082501  
NPA 754 (2005) 58c

$^7\text{Li}$  etc. ( $K^-_{\text{stop}}, \gamma \pi^-$ )



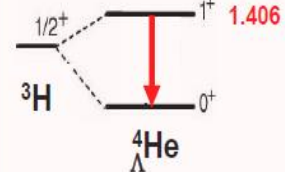
PLB 62 (1976) 46  
PLB 83 (1979) 25

$^7\text{Li} (\pi^+, K^+ \gamma)$  KEK E419



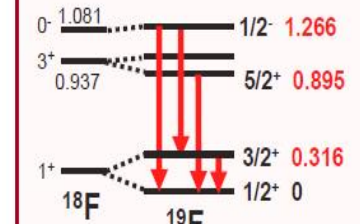
PRL 84 (2000) 5963  
PRL 86 (2001) 1982  
PLB 579 (2004) 258  
PRC 73 (2006) 012501

$^4\text{He} (K^-, \pi^+ \gamma)$  J-PARC E13



PRL 115 (2015) 222501

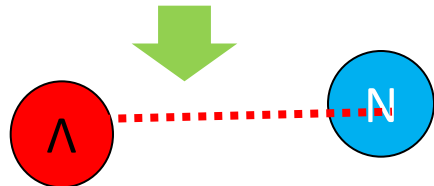
$^{19}\text{F} (K^-, \pi^+ \gamma)$  J-PARC E13



PRL 120 (2018) 132505

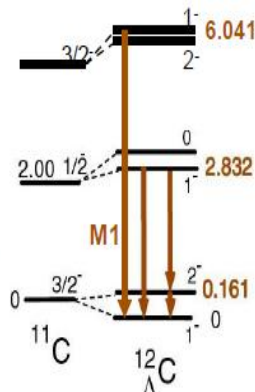
Few-body calculation  
Shell model calculation  
+

High-resolution experiments



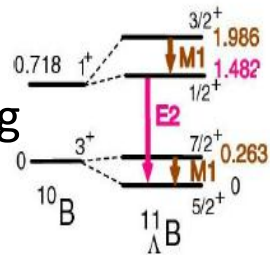
We have been obtaining information on  $\Lambda N$  two-body interaction.

$^{12}\text{C} (\pi^+, K^+ \gamma)$  KEK E566



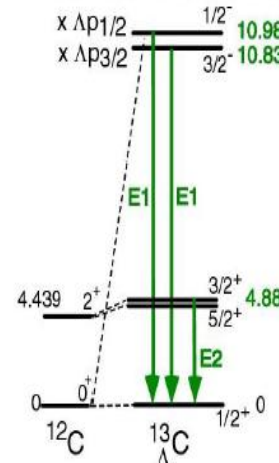
PTEP (2015) 081D01

$^{11}\text{B} (\pi^+, K^+ \gamma)$  KEK E518



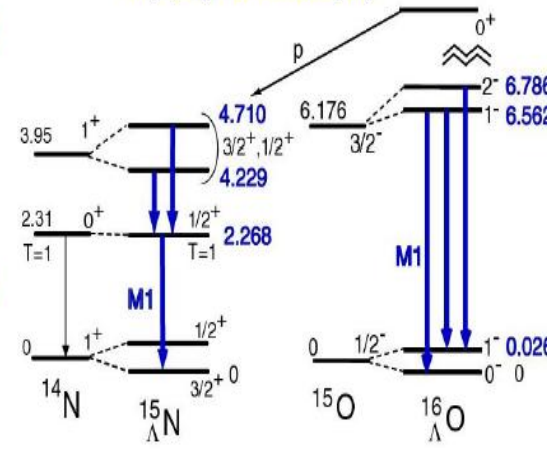
NPA835 (2010) 422

$^{13}\text{C} (K^-, \pi^+ \gamma)$  BNL E929 (Nal)



PRL 86 (2001) 4255  
PRC 65 (2002) 034607

$^{16}\text{O} (K^-, \pi^+ \gamma)$  BNL E930('01)



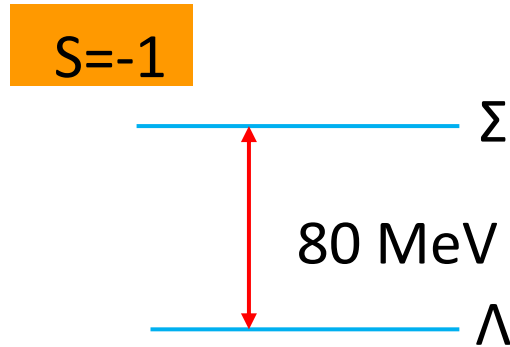
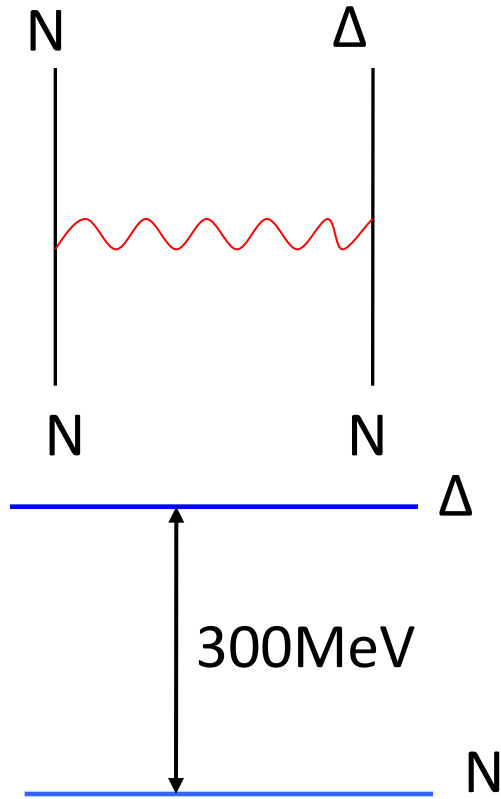
PRC 77 (2008) 054315

PRL 93 (2004) 232501  
EPJ A33 (2007) 247

$$V_{\Lambda N} = V_0 + \sigma_{\Lambda} \cdot \sigma_N V_{\sigma\sigma} + \mathbf{L} \cdot (\mathbf{s}_{\Lambda} + \mathbf{s}_N) V_{\text{SLS}} + \mathbf{L} \cdot (\mathbf{s}_{\Lambda} - \mathbf{s}_N) V_{\text{ALS}} + S_{12} V_{\text{tensor}} + \dots$$

# $\Lambda N - \Sigma N$ coupling

Non-strangeness sector



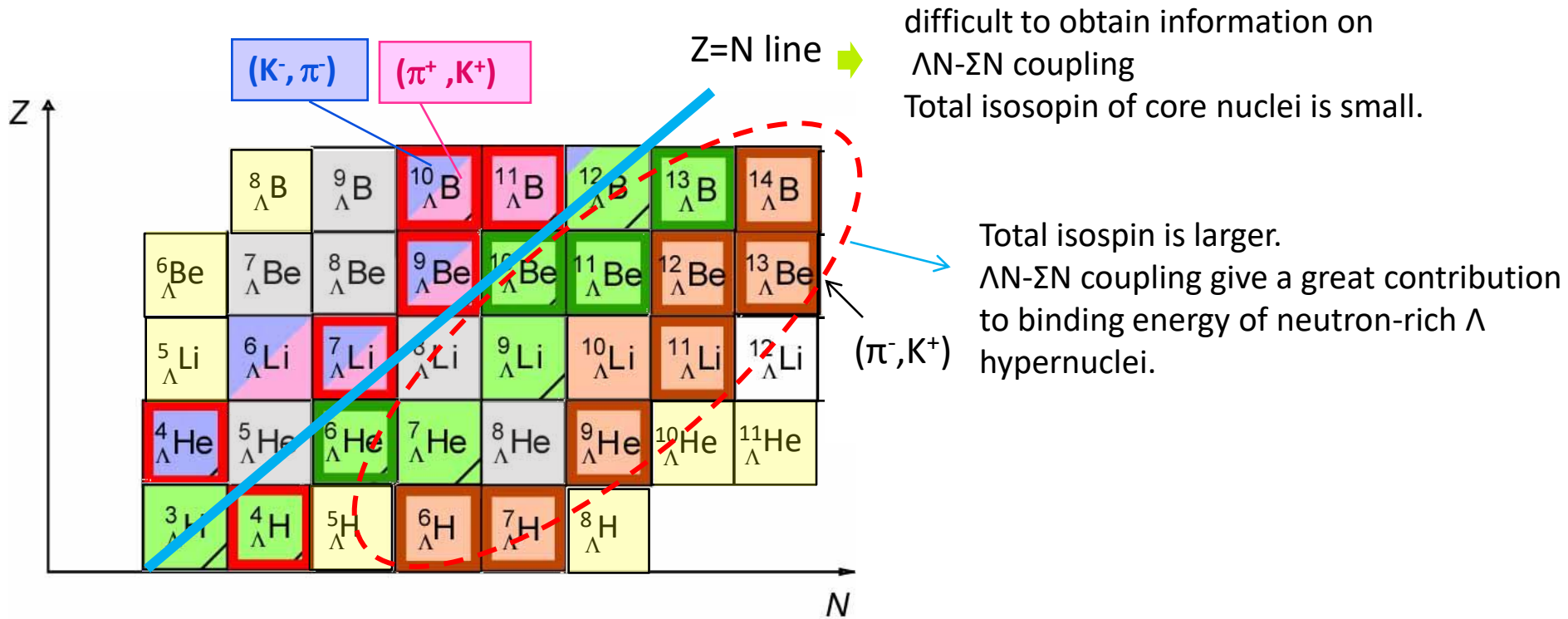
Mass is smaller.  
It is expected that  $\Lambda$ - $\Sigma$  conversion might affect in structure of  $\Lambda$  hypernuclei.

Probability of  $\Delta$  in nuclei is not large.

$\Lambda N$ - $\Sigma N$  coupling  $\Rightarrow$  three-body force and CSB interaction

# How do we obtain information on $\Lambda N$ - $\Sigma N$ coupling?

- (1)  $\Lambda N$  scattering experiment at J-PARC, Femtoscopic experiment
- (2) To study neutron-rich  $\Lambda$  hypernuclei at J-PARC



difficult to obtain information on  $\Lambda N$ - $\Sigma N$  coupling  
Total isospin of core nuclei is small.

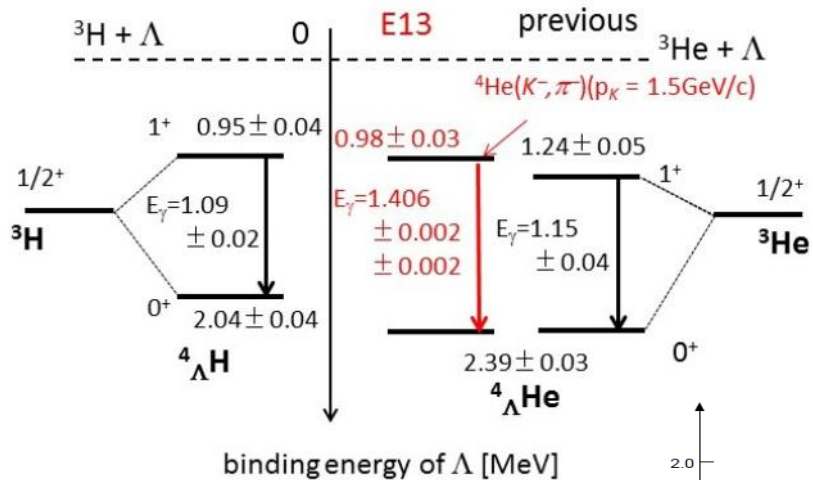
Total isospin is larger.  
 $\Lambda N$ - $\Sigma N$  coupling give a great contribution to binding energy of neutron-rich  $\Lambda$  hypernuclei.

These neutron-rich  $\Lambda$  hypernuclei are important.

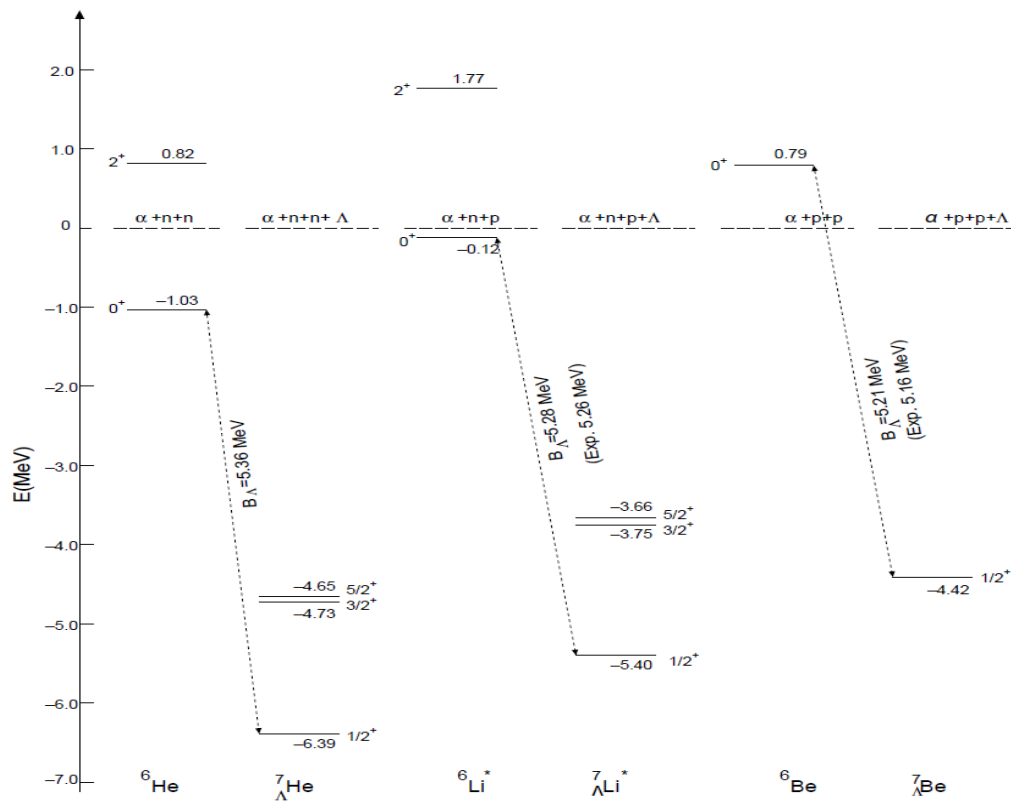
By neutron-rich  $\Lambda$  hypernuclei, we could obtain information on **long-range part** of  $\Lambda N$ - $\Sigma N$  coupling.  $\rightarrow$  **Long-range part** of  $\Lambda NN$  three-body force

Furthermore, we need short-range part of  $\Lambda NN$  three-body force.: important for **the study of neutron star**

# CSB interaction



How do we understand this difference?  
Odd state of CSB?



Now, it is interesting to see as follows:

(1) What is the level structure of  $A=7$  hypernuclei without CSB interaction?

(2) What is the level structure of  $A=7$  hypernuclei with CSB interaction?

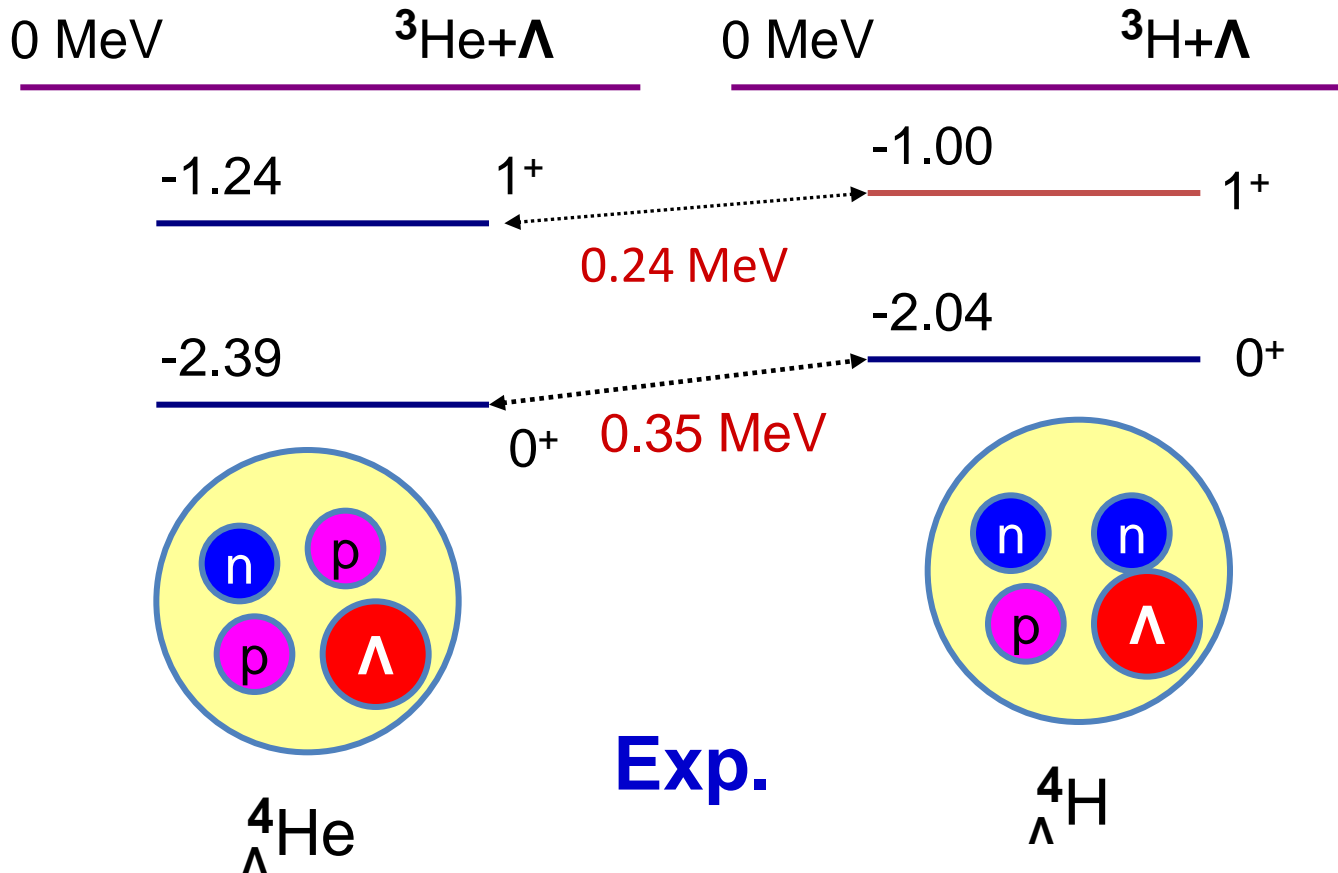
Next we introduce a phenomenological CSB potential with the central force component only.

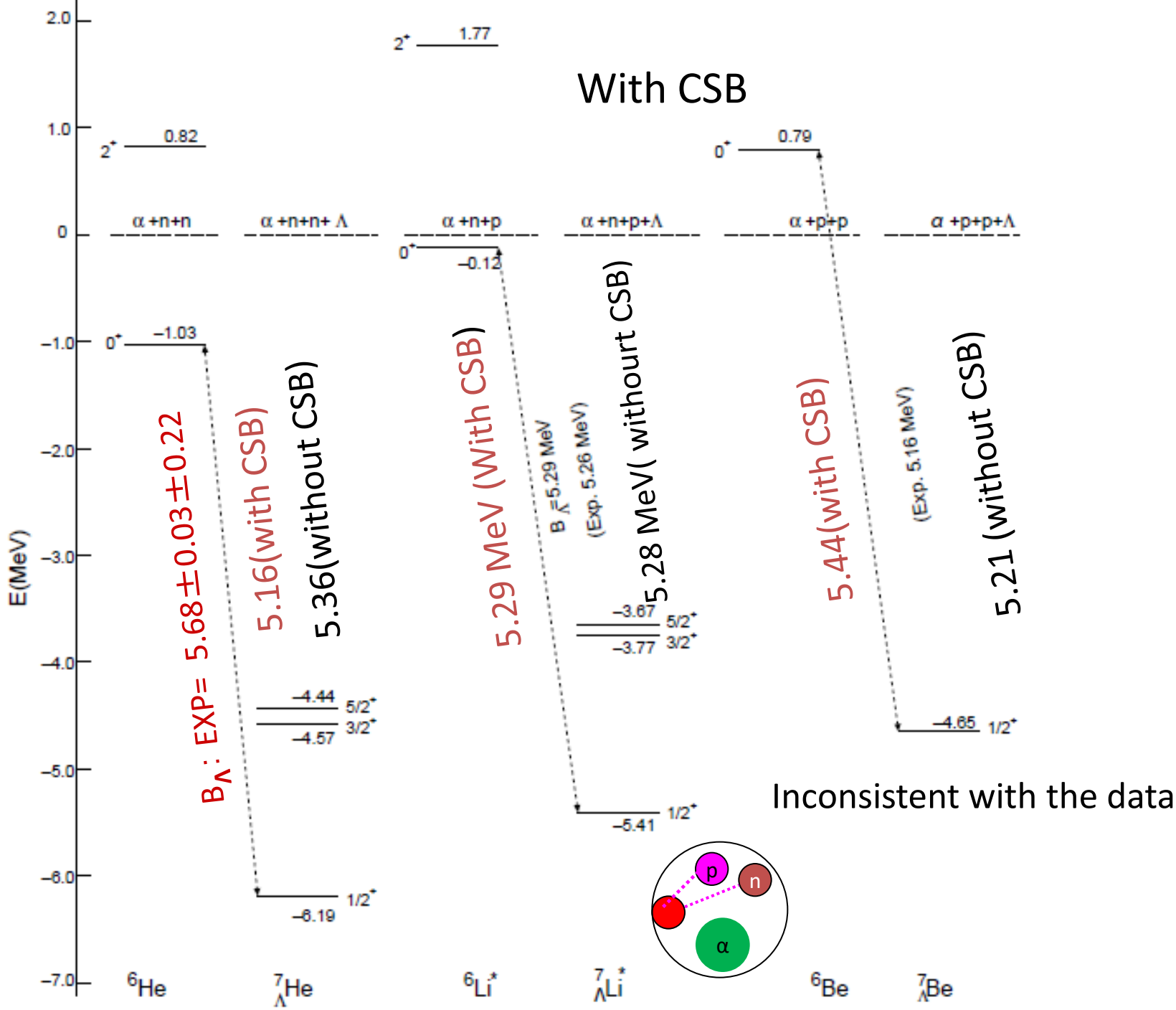
$$V_{\Lambda N}^{\text{CSB}}(r) = \quad (3.3)$$

$$-\frac{\tau_z}{2} \left[ \frac{1+P_r}{2} (v_0^{\text{even,CSB}} + \sigma_\Lambda \cdot \sigma_N v_{\sigma_\Lambda \cdot \sigma_N}^{\text{even,CSB}}) e^{-\beta_{\text{even}} r^2} \right.$$

$$\left. + \frac{1-P_r}{2} (v_0^{\text{odd,CSB}} + \sigma_\Lambda \cdot \sigma_N v_{\sigma_\Lambda \cdot \sigma_N}^{\text{odd,CSB}}) e^{-\beta_{\text{odd}} r^2} \right],$$

Strength, range are determined so as to reproduce the data.







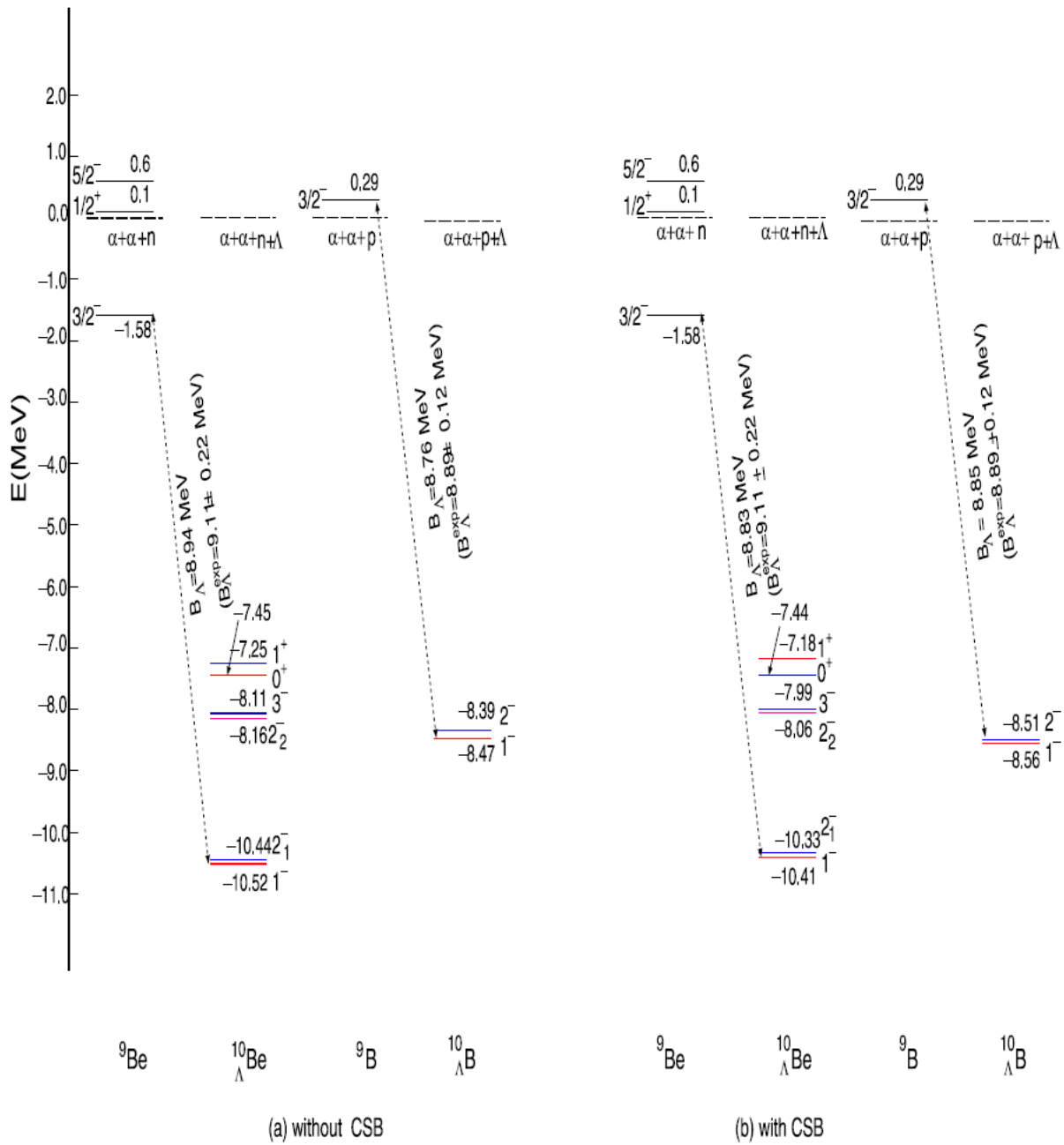
We include odd-state of CSB to reproduce the data of A=7  
 $\Lambda$  hypernuclei. =>apply to A=10  $\Lambda$  hypernuclei

Progress of Theoretical Physics, Vol. 128, No. 1, July 2012

## Structure of $^{10}_{\Lambda}\text{Be}$ and $^{10}_{\Lambda}\text{B}$ Hypernuclei Studied with the Four-Body Cluster Model

Emiko HIYAMA and Yasuo YAMAMOTO

$$V_{\Lambda N}^{\text{CSB}}(r) = -\frac{\tau_z}{2} \left[ \frac{1 + P_r}{2} (v_0^{\text{even,CSB}} + \boldsymbol{\sigma}_{\Lambda} \cdot \boldsymbol{\sigma}_N v_{\sigma_{\Lambda} \cdot \sigma_N}^{\text{even,CSB}}) e^{-\beta_{\text{even}} r^2} \right. \\ \left. + \frac{1 - P_r}{2} (v_0^{\text{odd,CSB}} + \boldsymbol{\sigma}_{\Lambda} \cdot \boldsymbol{\sigma}_N v_{\sigma_{\Lambda} \cdot \sigma_N}^{\text{odd,CSB}}) e^{-\beta_{\text{odd}} r^2} \right],$$



What about  
Experimental data?

Even CSB)

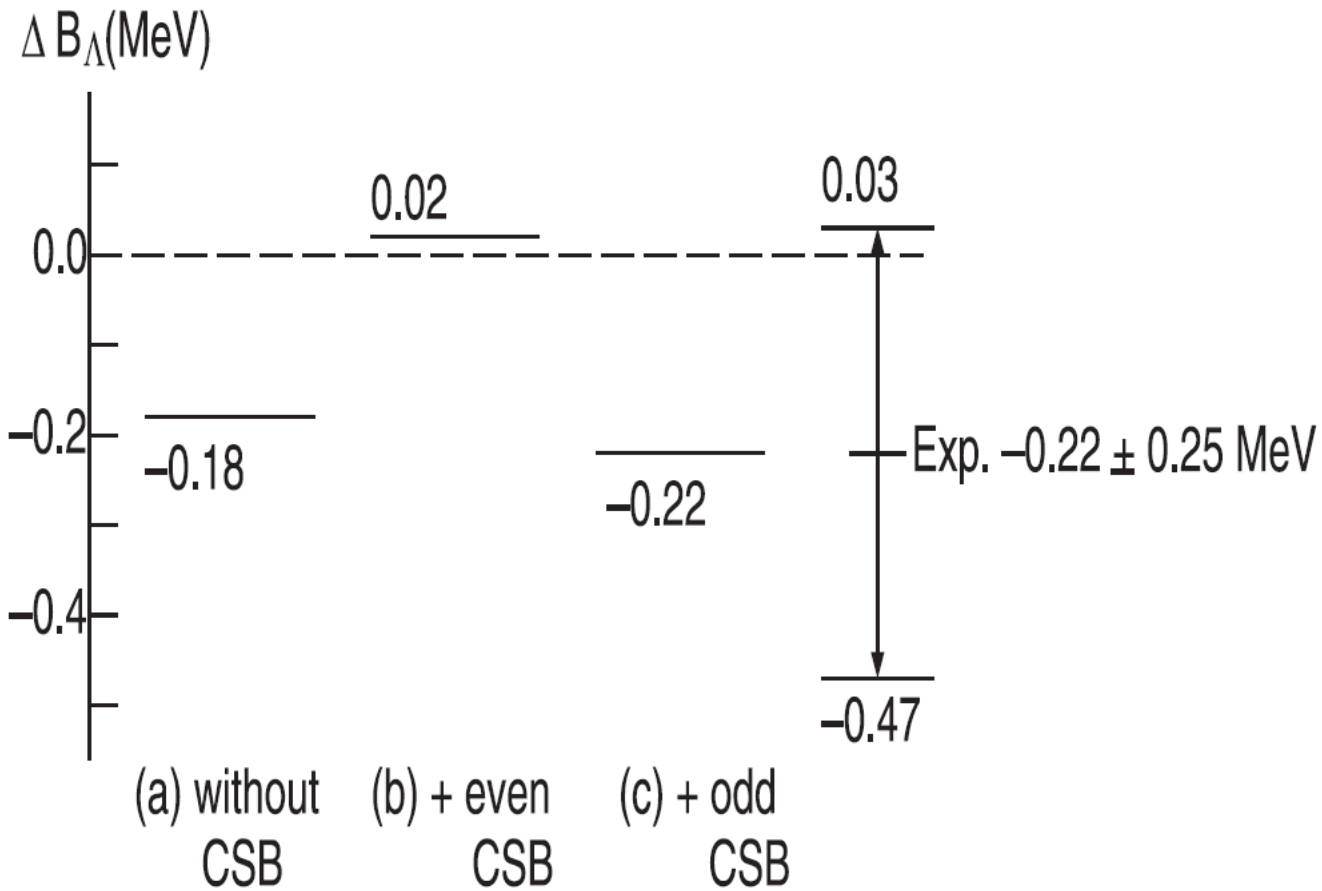


Fig. 7. Calculated energy difference of  ${}_{\Lambda}^{10}\text{Be}$  and  ${}_{\Lambda}^{10}\text{B}$ ,  $\Delta B_{\Lambda}(B_{\Lambda}({}_{\Lambda}^{10}\text{B}) - B_{\Lambda}({}_{\Lambda}^{10}\text{Be}))$ , (a) without CSB, (b) with even-state CSB, and (c) with both even- and odd-state CSB interactions.

What about data by Gogami san?

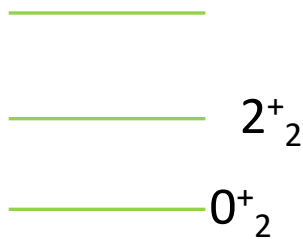
$^{10}_{\Lambda}\text{Be}$  is interesting to investigate structure modified by addition of  $\Lambda$  particle.

# Schematic illustration

shell-like states



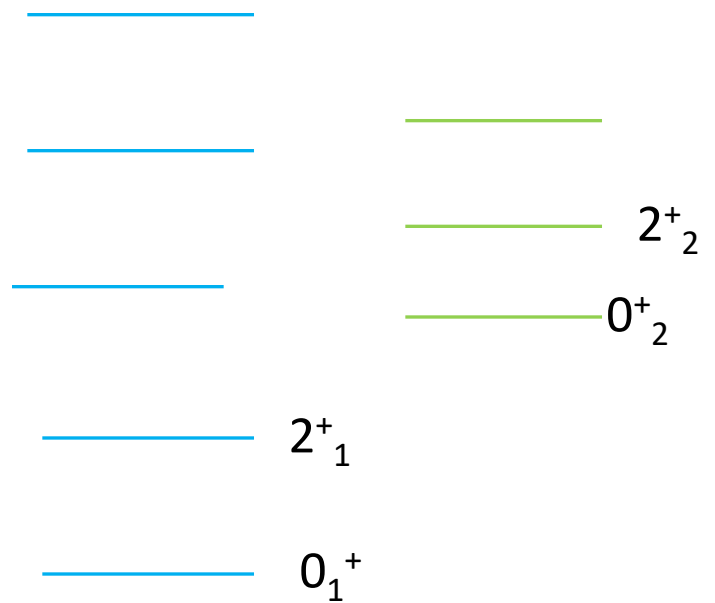
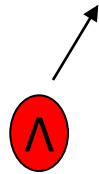
$\alpha$ -clustering states



**Question** : Does energy gain go in parallel way for all the states?

**No !**

$A \geq 9$  core nucleus



$A \geq 10$   $\Lambda$  hypernucleus

# Energy gain by $\Lambda$ -particle addition

$$\Delta E(\text{shell-like}) > \Delta E(\text{clustering})$$

shell-like  
state

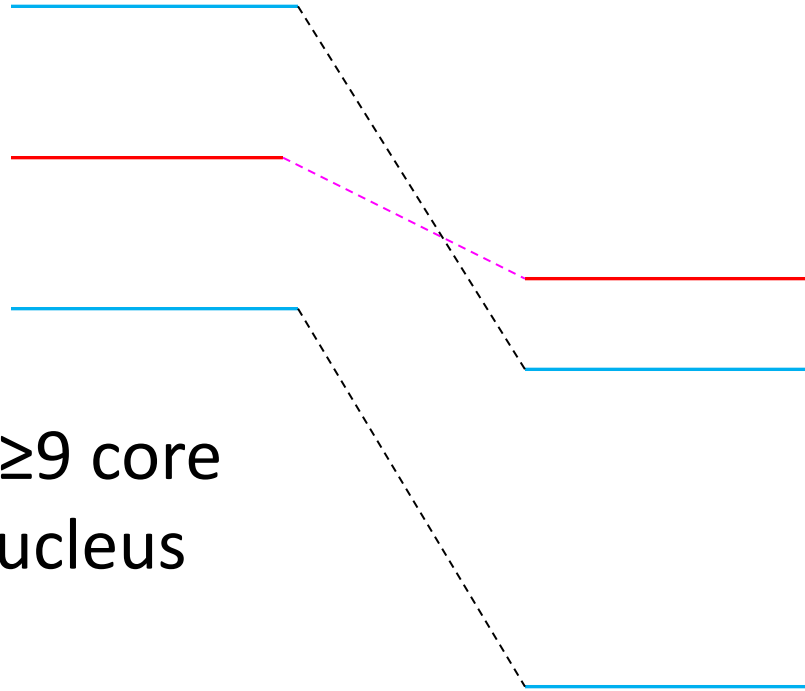
clustering  
state

shell-like  
state

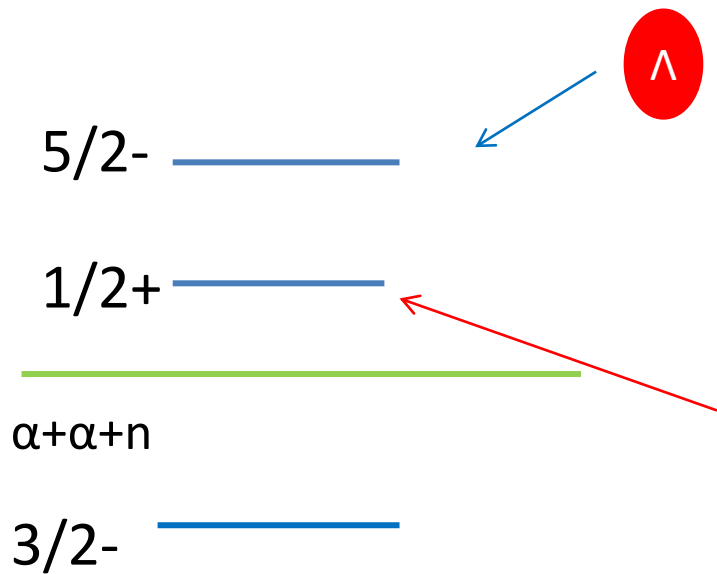
$A \geq 9$  core  
nucleus

$A \geq 10$   $\Lambda$  hypernucleus

} Level crossing

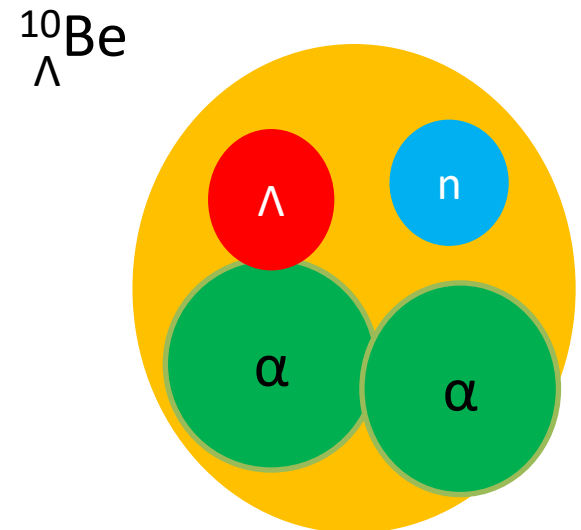
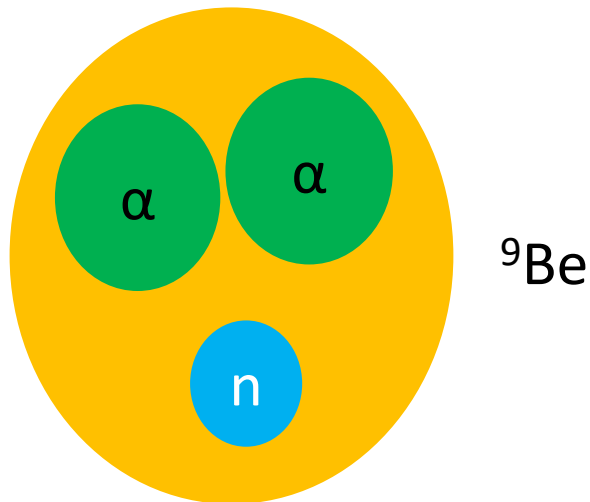


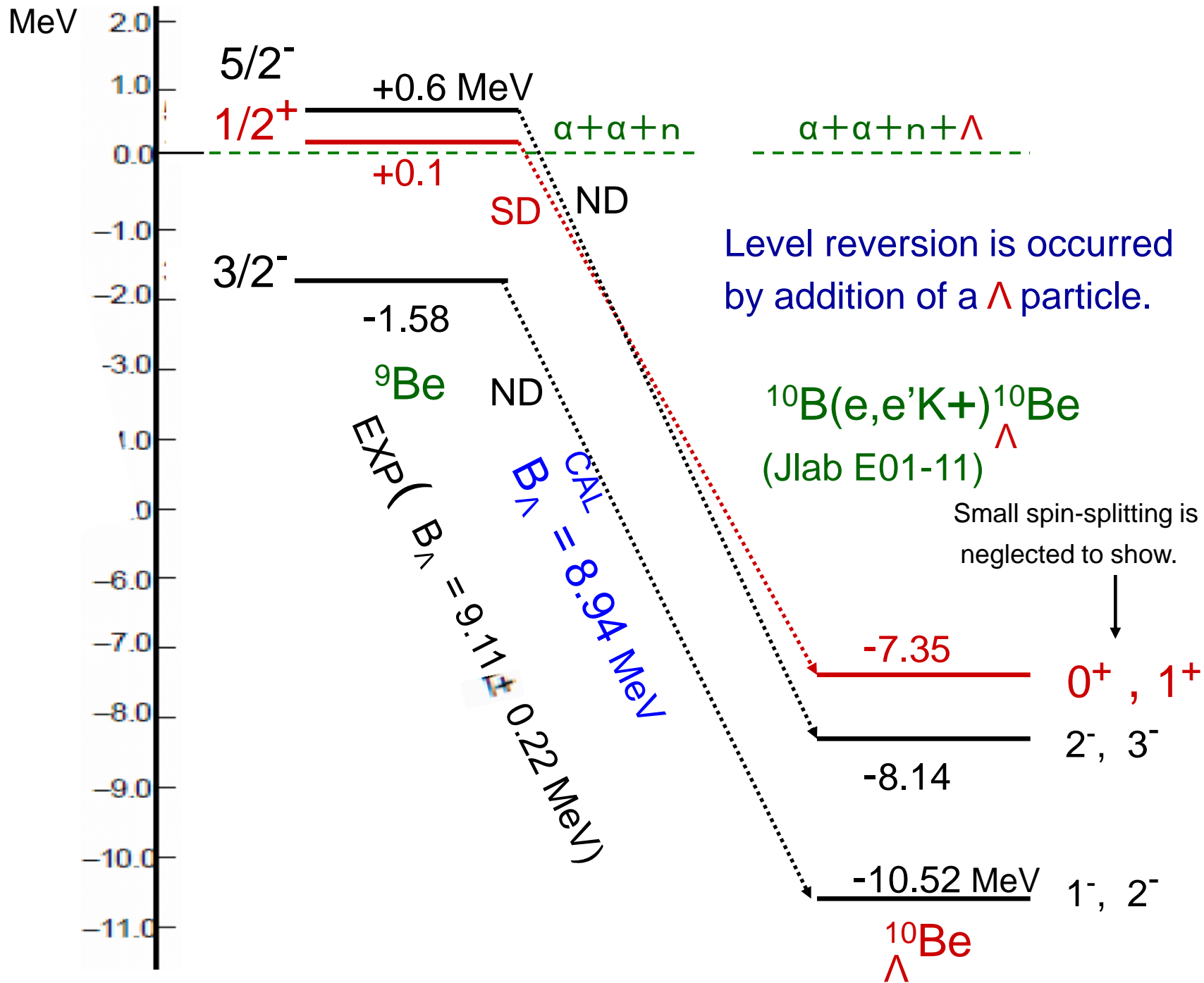
Interesting phenomena such as level crossing is seen in  ${}^9\text{Be}$  and  ${}^{10}\text{Be}$  combination.



When a  $\Lambda$  particle is injected into each state, is there change of energy spectra?

It is well known that a valence neutron is occupied in  $1S_{1/2}$  orbit in the simple shell model configuration.

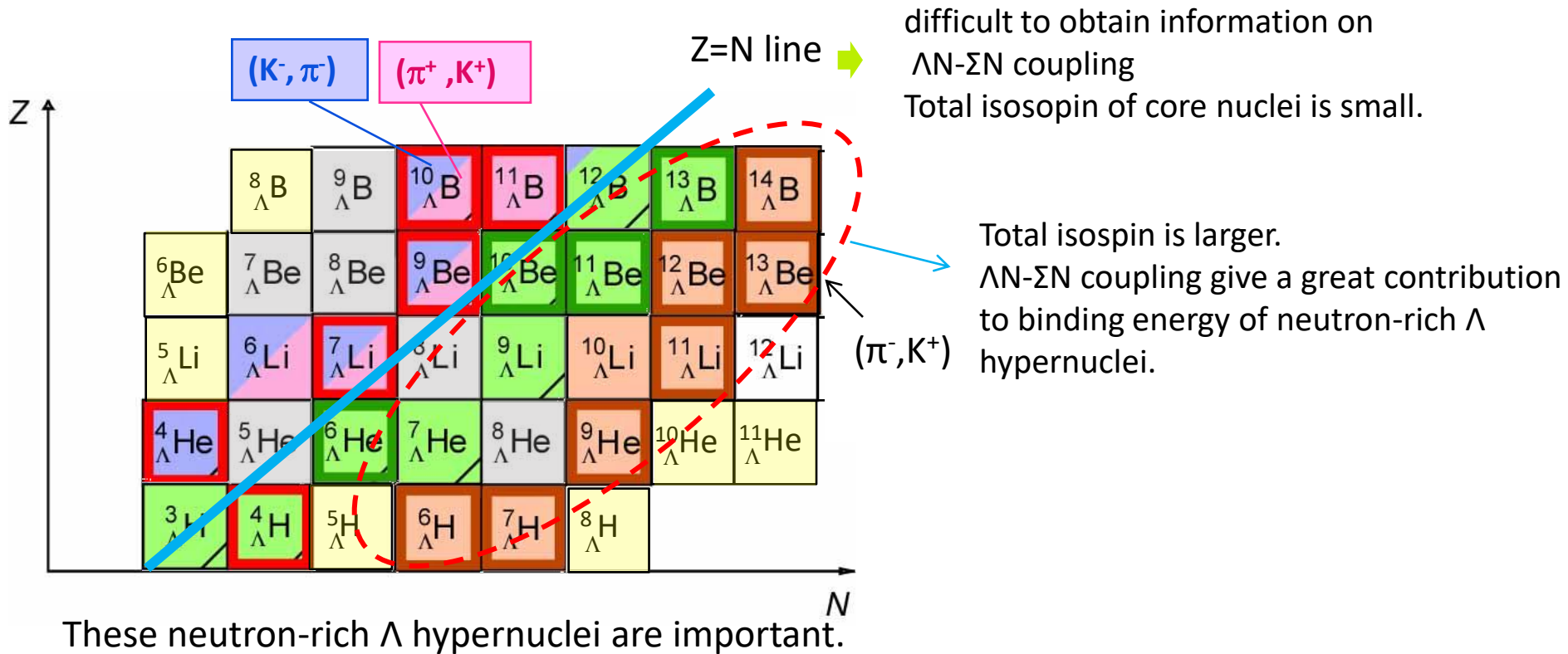






## How do we obtain information on $\Lambda N$ - $\Sigma N$ coupling?

- (1)  $\Lambda N$  scattering experiment at J-PARC, Femtoscopic experiment
- (2) To study neutron-rich  $\Lambda$  hypernuclei at J-PARC



By neutron-rich  $\Lambda$  hypernuclei, we could obtain information on

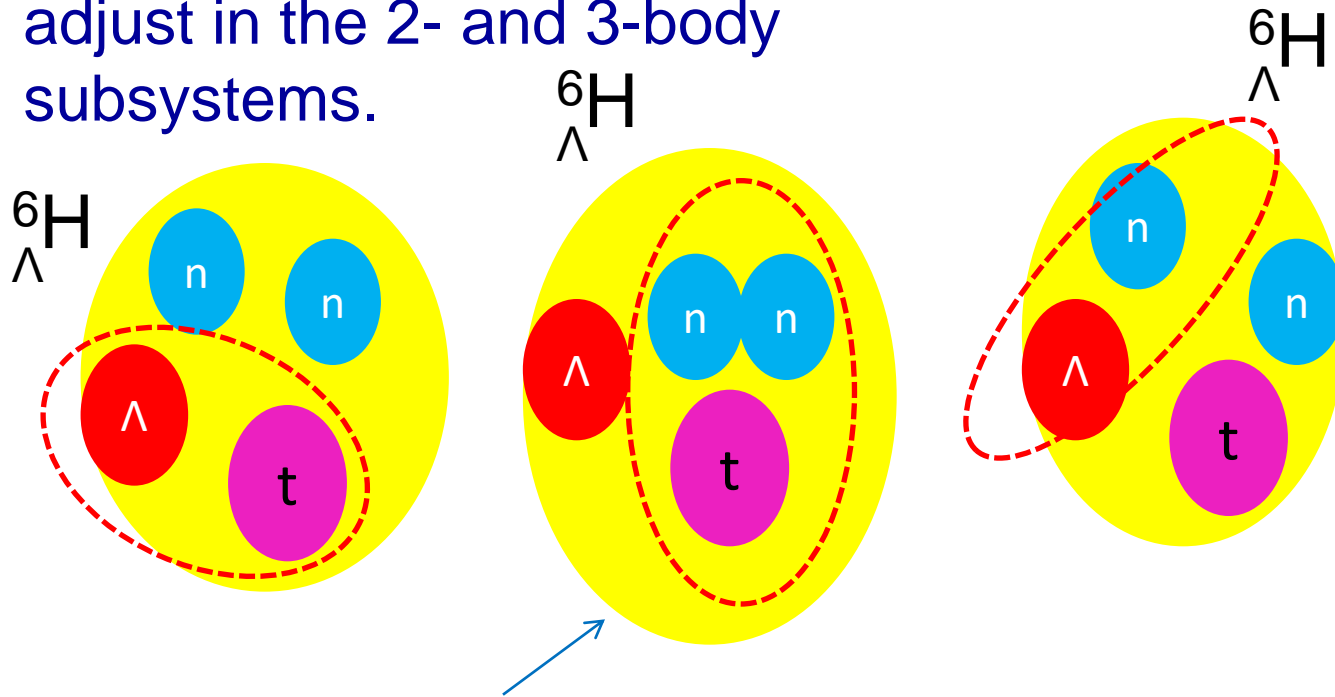
**long-range part** of  $\Lambda N$ - $\Sigma N$  coupling.  $\rightarrow$  **Long-range part** of  $\Lambda NN$  three-body force

Furthermore, we need short-range part of  $\Lambda NN$  three-body force.: important for **the study of neutron star**

So far, search  ${}^6_{\Lambda}\text{H}$  experiment has been done.

Conflict between FINUDA exp. and J-PARC exp.

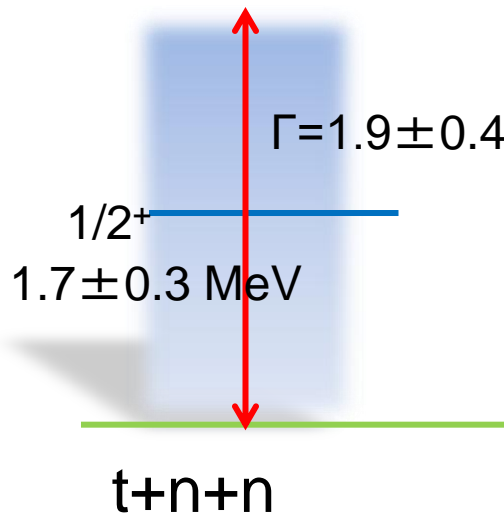
Before doing full 4-body calculation,  
 it is important and necessary to reproduce  
 the observed binding energies of all the sets  
 of subsystems in  ${}^6_{\Lambda}\text{H}$ .  
 Namely, All the potential parameters  
 are needed to  
 adjust in the 2- and 3-body  
 subsystems.



Among the subsystems, it is extremely important to  
 adjust the energy and decay width of  ${}^5\text{H}$  core nucleus.

## Framework:

To calculate the binding energy of  ${}_{\Lambda}^6\text{H}$ , it is very important to reproduce the binding energy of the core nucleus  ${}^5\text{H}$ .



transfer reaction  $p({}^6\text{He}, {}^2\text{He}){}^5\text{H}$

A. A. Korcheninnikov, et al. Phys. Rev. Lett. **87** (2001) 092501.

To reproduce the data, for example, [R. De Diego et al, Nucl. Phys. A786 \(2007\), 71.](#) calculated the energy and width of  ${}^5\text{H}$  with  $t+n+n$  three-body model using complex scaling method. The calculated binding energy for the ground state of  ${}^5\text{H}$  is 1.6 MeV with respect to  $t+n+n$  threshold and width has 1.5 MeV.

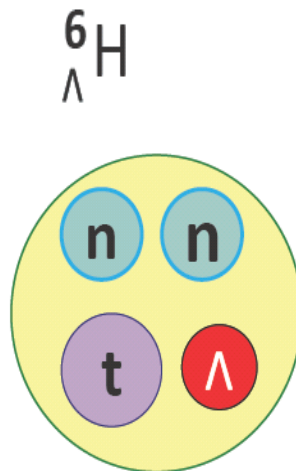
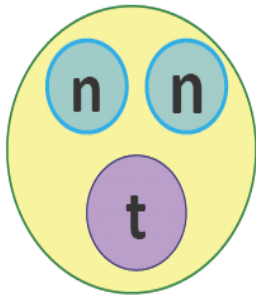
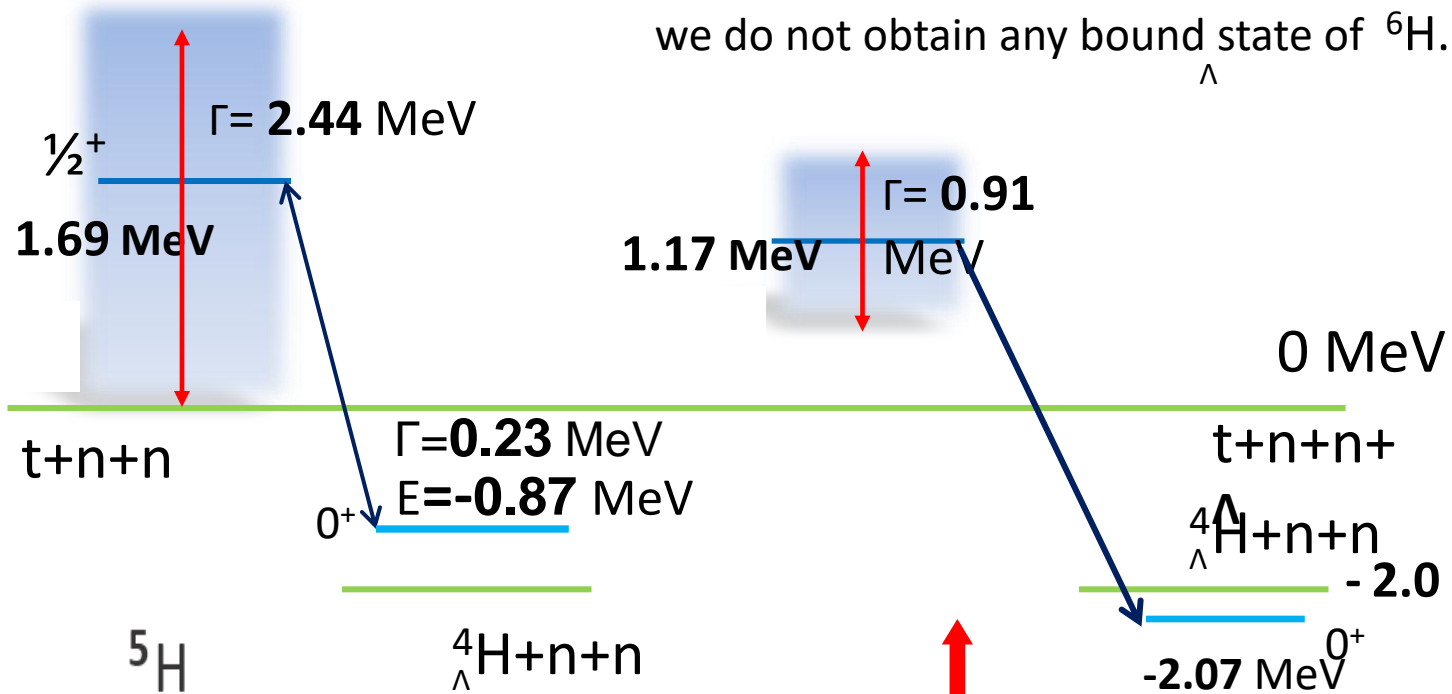
Exp:  $1.7 \pm 0.3$  MeV

$\Gamma = 1.9 \pm 0.4$  MeV



Even if the potential parameters were tuned so as to reproduce the lowest value of the Exp.,  $E = 1.4$  MeV,  $\Gamma = 1.5$  MeV,

we do not obtain any bound state of  ${}^6_{\Lambda}\text{H}$ .



On the contrary, if we tune the potentials to have a bound state in  ${}^6_{\Lambda}\text{H}$ , then what is the energy and width of  ${}^5\text{H}$ ?

$(E_R, \Gamma_R)$ (MeV)	
$J^\pi$	$1/2^+$
${}^5\text{H}$ (full)	(1.57, 1.53)
${}^5\text{H}$ ( $d = 0$ )	(1.55, 1.35)
Theor. [16]	(2.26, 2.93)
Theor. [12]	(2.5–3.0, 3–4)
Theor. [13]	(3.0–3.2, 1–4)
Theor. [15]	(1.59, 2.48)
Exp. [3]	$(1.7 \pm 0.3, 1.9 \pm 0.4)$ ←
Exp. [8]	$(1.8 \pm 0.1, < 0.5)$
Exp. [4]	(1.8, 1.3)
Exp. [5]	(2, 2.5)
Exp. [6]	(3, 6)
Exp. [9]	$(5.5 \pm 0.2, 5.4 \pm 0.6)$

We cited this experiment. However, you have many different decay widths. width is strongly related to the size of wavefunction. Then, I hope that The decay width will be determined at RCNP this year.

[3] A.A. Koroshennikov et al., PRL87 (2001) 092501

[8] S.I. Sidorchuk et al., NPA719 (2003) 13

[4] M.S. Golovkov et al. PRC 72 (2005) 064612

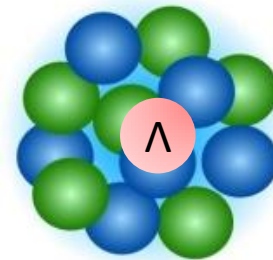
[5] G. M. Ter-Akopian et al., Eur. Phys. J A25 (2005) 315.

For the study of neutron-rich  $\Lambda$  hypernuclei,  
We should exclude ambiguity of the energy and decay width of  
core nuclei.

For this purpose,  ${}^8\text{He}$  and  ${}^9_{\Lambda}\text{He}$  is the best candidate to  
perform search experiment.

# heavier $\Lambda$ hypernuclei

For example: Pb, Sn, Zr, La, Y etc. Istopes



Density of heavier nuclei is high and then,  $\Lambda$  particle is acting in such high dense matter. => We could obtain information on the short-range part of  $\Lambda$ NN interaction.

In heavier nuclei, density becomes high.

$^{208}_{\Lambda}\text{Pb}$ ,  $^{139}_{\Lambda}\text{La}$ ,  $^{89}_{\Lambda}\text{Y}$ : plan in the project at HIHR

I request the experimentalists to produce  $^{112-124}_{\Lambda}\text{Sn}$ ,  $^{90-96}_{\Lambda}\text{Zr}$   $\Lambda$  hypernuclei.

Heavy  $\Lambda$  hypernuclei exp. + theoretical cal.

Isotope dependence of  $\Lambda$ NN three-body force

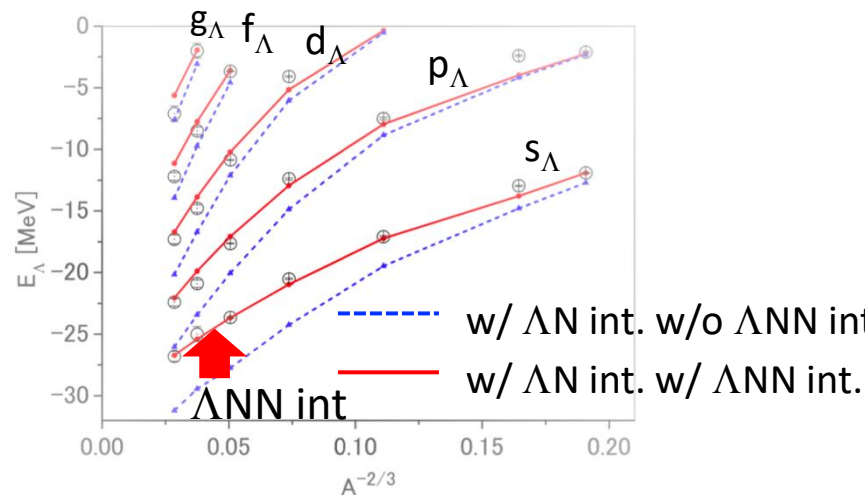


Determine  $\Lambda$ NN interaction



Reliable EOS

Calc. by Nijmegen ESC16 model





# Effects of Many-body Interactions in the Hyperon with Skyrme and KIDS Functionals for Finite Nuclei and Nuclear Matter

Soonchul Choi,<sup>1,2</sup> Emiko Hiyama,<sup>3</sup> Chang Ho Hyun,<sup>4,\*</sup> and Myung-Ki Cheoun<sup>1</sup>

<sup>1</sup>*Department of Physics, Soongsil University, Seoul 06978, Korea*

<sup>2</sup>*Center for Exotic Nuclear Studies,*

*Institute for Basic Science, Daejeon 34126, Korea*

<sup>3</sup>*Department of Physics, Tohoku University, Sendai 980-8578, Japan*

<sup>4</sup>*Department of Physics Education, Daegu University, Gyeongsan 38453, Korea*

Submitted to EPJA

$$H_{\Lambda\rho} = \frac{3}{8}\rho_{\Lambda} \sum_{i=1}^{N_f} u_{3i} \left(1 + \frac{1}{2}y_{3i}\right) \rho_N^{1+i/3}.$$

This new type of functional is added to the two-body interaction terms given as

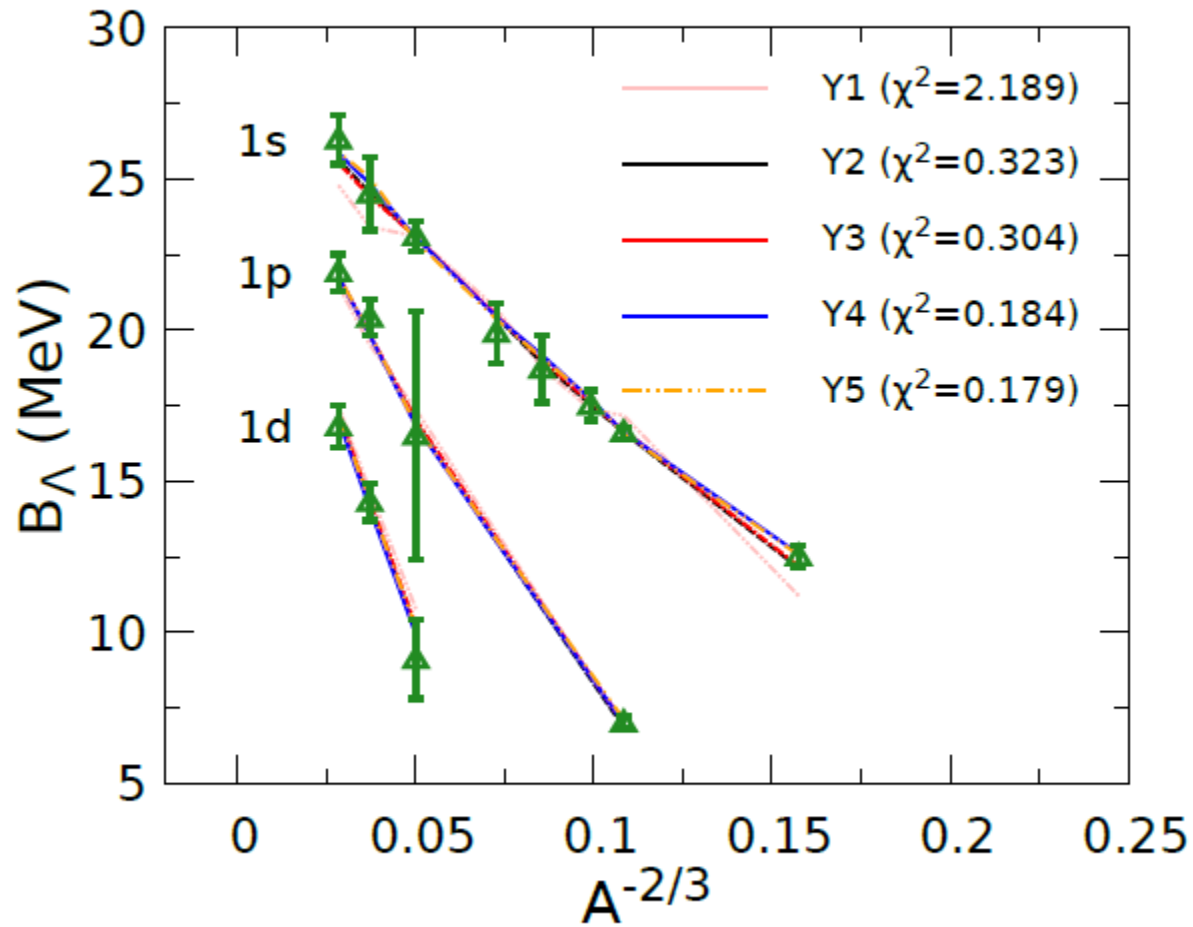
$$\begin{aligned} H_{\Lambda N} = & \frac{\hbar^2}{2m_{\Lambda}}\tau_{\Lambda} + u_0 \left(1 + \frac{1}{2}y_0\right) \rho_N \rho_{\Lambda} \\ & + \frac{1}{4}(u_1 + u_2)(\tau_{\Lambda}\rho_N + \tau_N\rho_{\Lambda}) + \frac{1}{8}(3u_1 - u_2)(\nabla\rho_N \cdot \nabla\rho_{\Lambda}) \\ & + \frac{1}{2}W_{\Lambda}(\nabla\rho_N \cdot \vec{J}_{\Lambda} + \nabla\rho_N \cdot \vec{J}_{\Lambda}). \end{aligned}$$

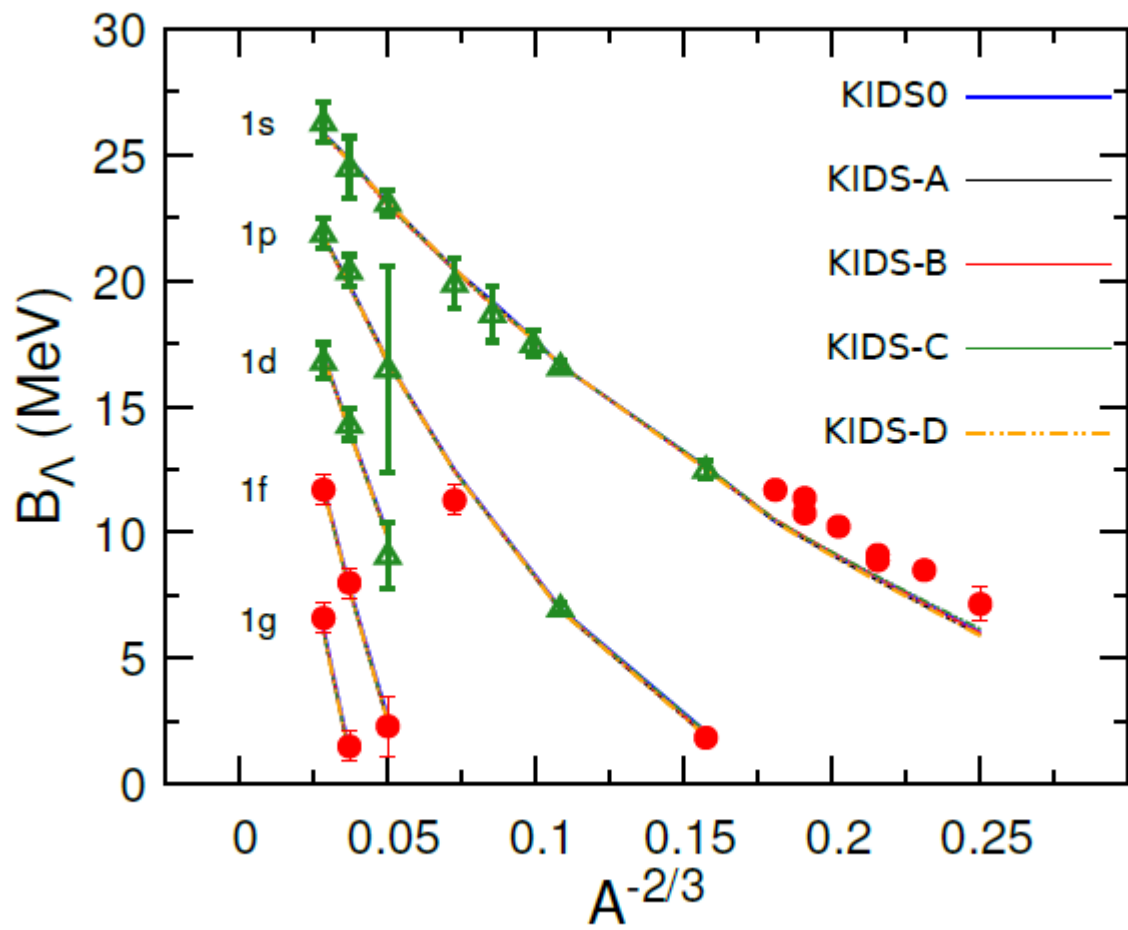
Nuclear density functional theory: two-body interaction + many body force

Model parameters  $u_0$ ,  $y_0$ ,  $u_1$ ,  $u_2$ ,  $W_{\Lambda}$ ,  $u_{3i}$  and  $y_{3i}$  :fitted so as to reproduce the data

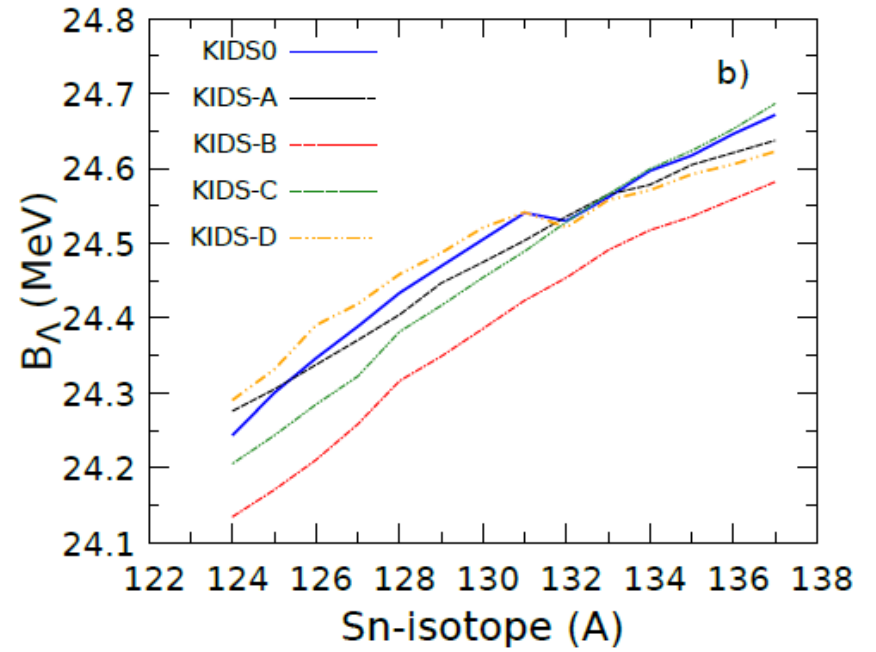
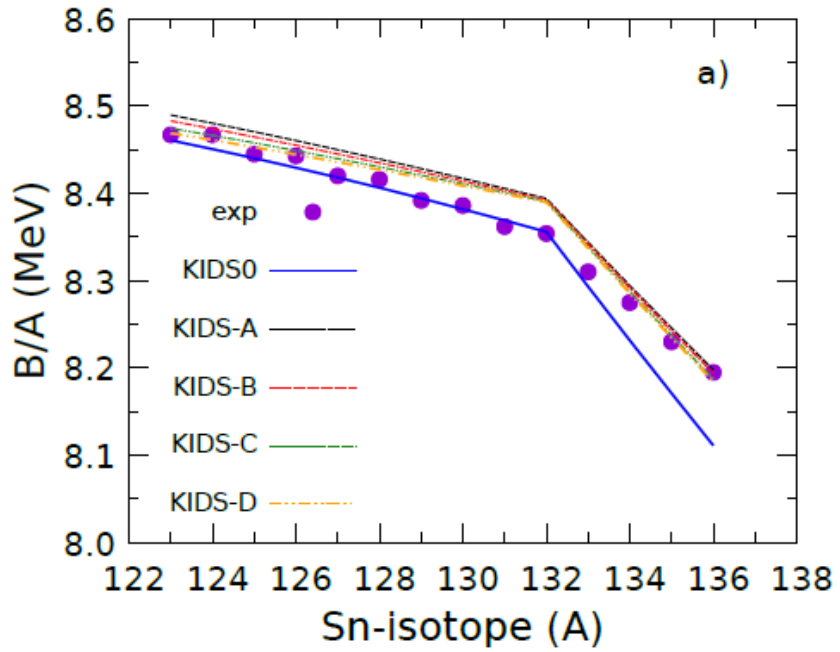
Nuclei	$1s$ (MeV)	$1p$ (MeV)	$1d$ (MeV)
${}_{\Lambda}^{16}\text{O}$ [13, 14]	$12.50 \pm 0.35$		
${}_{\Lambda}^{28}\text{Si}$ [2, 15]	$16.60 \pm 0.20$	$7.0 \pm 0.2$	
${}_{\Lambda}^{32}\text{S}$ [16]	$17.50 \pm 0.50$		
${}_{\Lambda}^{40}\text{Ca}$ [17, 18]	$18.70 \pm 1.1$		
${}_{\Lambda}^{51}\text{V}$ [19, 20]	$19.9 \pm 1.0$		
${}_{\Lambda}^{89}\text{Y}$ [2, 20]	$23.10 \pm 0.50$	$16.50 \pm 4.1$	$9.1 \pm 1.3$
${}_{\Lambda}^{139}\text{La}$ [2]	$24.50 \pm 1.20$	$20.40 \pm 0.6$	$14.3 \pm 0.6$
${}_{\Lambda}^{208}\text{Pb}$ [2]	$26.30 \pm 0.80$	$21.90 \pm 0.6$	$16.8 \pm 0.7$

We construct 5 types interaction.

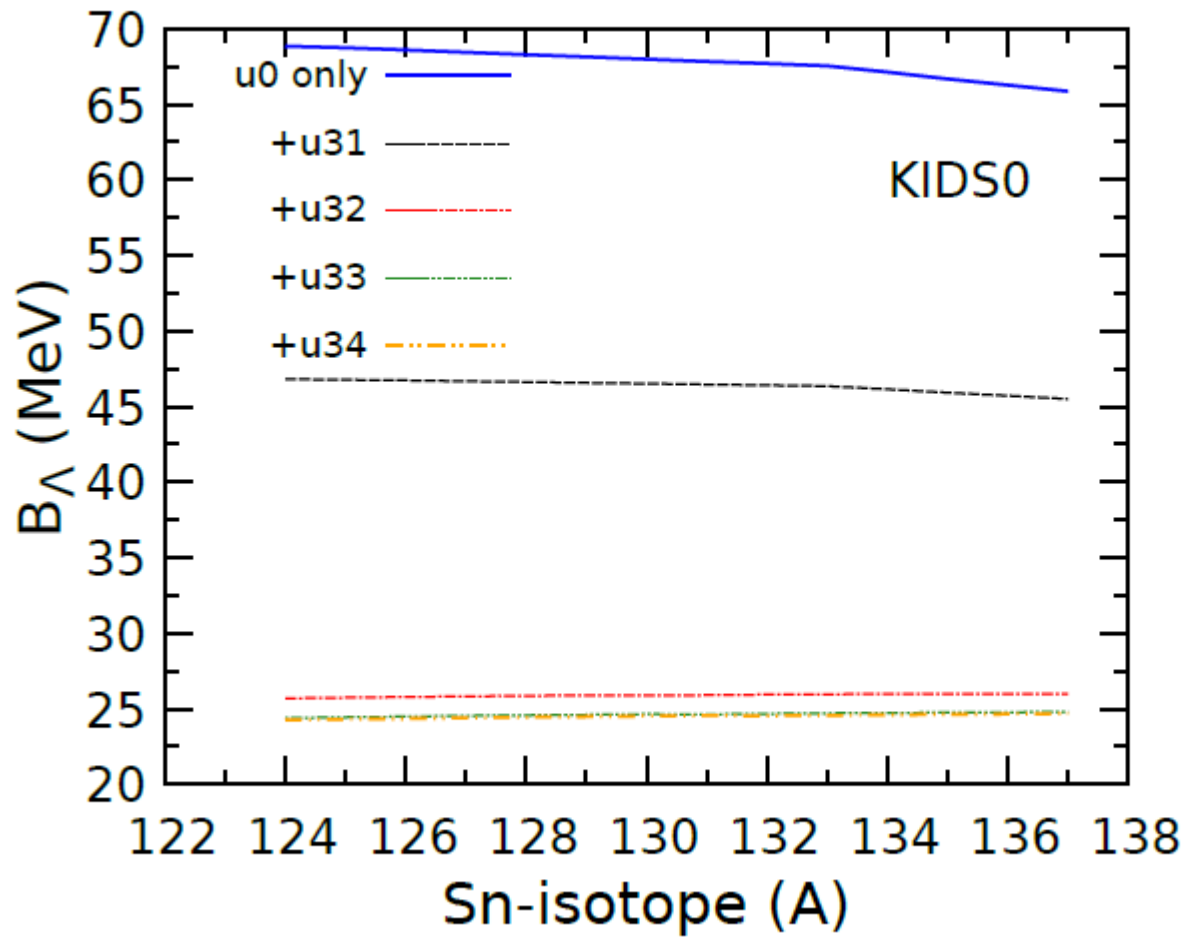




Apply to Sn



The energies of core nuclei, Sn isotope are reproduced.



Three-body  
Effect is repulsive.

# Summary

$\Lambda$ N- $\Sigma$ N coupling: neutron-rich  $\Lambda$  hypernuclei

It would be better to have deeper binding.

Heavier neutron-rich  $\Lambda$  hypernuclei

Sn Isotope?