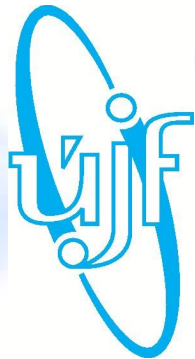


Equation of Motion Phonon Method in Hypernuclei - progress report



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**Collaboration Japan - Czech Republic,
HIEI2022, March 2022**

Motivation

Work on development of **many-body model(s)** suitable for description of **nuclear and hypernuclear structure** which aim to describe wide range of hypernuclei including medium-size & heavy systems.

This model can be used in calculations of **hypernuclear production** – especially the hypernuclei whose production is planned by experimentalists in close future (${}^{40}_{\Lambda}\text{K}$, ${}^{48}_{\Lambda}\text{K}$, ${}^{208}_{\Lambda}\text{Tl}$).

Progress report:

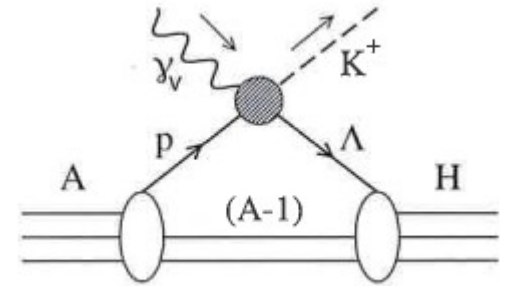
- **Electroproduction** of Hypernuclei
- Applications of **EMPM** for the Calculations of **Hypernuclei**
- **Results** – convergence test of EMPM & preliminary cross sections of the production of ${}^{12}_{\Lambda}\text{B}$
- **Summary** and Future Plans

Electroproduction of Hypernuclei

Hypernuclear production:

- Elementary process of electroproduction: $\mathbf{p (e, e' K^+) \Lambda}$
- **Kinematics** of the reaction
- Information about (many-body) **nuclear & hypernuclear structure**
- Information about the **$\Lambda\mathbf{N(N)}$ interactions**

All “ingredients“ important for description of hypernuclear production.



$$\frac{d^3\sigma}{dE_e d\Omega_e d\Omega_K} = \Gamma \left[\frac{d\sigma_U}{d\Omega_K} + \varepsilon_L \frac{d\sigma_L}{d\Omega_K} + \varepsilon \frac{d\sigma_P}{d\Omega_K} \cos 2\Phi_K + \sqrt{2\varepsilon_L(1+\varepsilon)} \frac{d\sigma_I}{d\Omega_K} \cos \Phi_K \right]$$

$$\frac{d\sigma_U}{d\Omega_K} = \frac{\beta}{2(2J_A+1)} \sum_{jm} \frac{1}{2j+1} (|A_{jm}^{+1}|^2 + |A_{jm}^{-1}|^2),$$

$$\frac{d\sigma_P}{d\Omega_K} = -\frac{\beta}{2J_A+1} \sum_{jm} \frac{1}{2j+1} \text{Re}\{A_{jm}^{+1} A_{jm}^{-1*}\},$$

$$\frac{d\sigma_L}{d\Omega_K} = \frac{\beta}{2J_A+1} \sum_{jm} \frac{1}{2j+1} |A_{jm}^0|^2,$$

$$\frac{d\sigma_I}{d\Omega_K} = \frac{\beta}{2J_A+1} \sum_{jm} \frac{1}{2j+1} \text{Re}\{A_{jm}^{0*} [A_{jm}^{+1} - A_{jm}^{-1}]\}$$

Transition amplitude

$$T_\lambda^{(1)} = \frac{Z}{[J_H]} \sum_{S\eta} \mathcal{F}_{\lambda\eta}^S \sum_{LM} \sum_{J_m} C_{LMS\eta}^{J_m} C_{J_A M_A J_m}^{J_H M_H} (J_H || F_{LM} [Y_L \otimes \sigma^S]^J || J_A)$$

Operator can be expressed in the second quantized form.

We evaluate

$$(\Phi_H || [b_\alpha^+ \otimes a_\alpha]^J || \Phi_A)$$

hypernucleus

creation Λ

annihilation p

nucleus

Electroproduction of Hypernuclei

Hypernuclear production:

$$\frac{d^3\sigma}{dE'_e d\Omega'_e d\Omega_K} = \Gamma \left[\frac{d\sigma_U}{d\Omega_K} + \epsilon_L \frac{d\sigma_L}{d\Omega_K} + \epsilon \frac{d\sigma_P}{d\Omega_K} \cos 2\Phi_K + \sqrt{2\epsilon_L(1+\epsilon)} \frac{d\sigma_I}{d\Omega_K} \cos \Phi_K \right]$$

$$M_\mu = \langle \Psi_H | \langle \chi_K | \sum_{j=1}^Z \hat{j}_\mu(j) | \chi_\gamma \rangle | \Psi_A \rangle \quad (\Phi_H || [b_\alpha^+ \otimes a_\alpha]^J || \Phi_A)$$

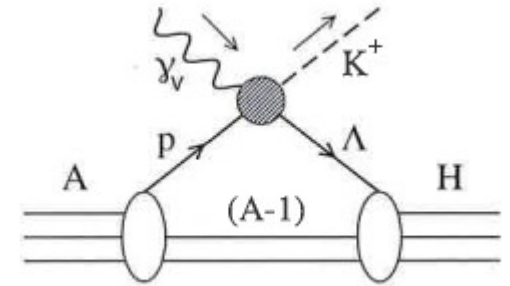
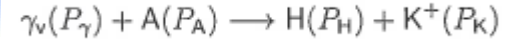


Fig. from M. Sotona, S. Frullani, *Prog. Theor. Phys. Suppl.* **117**, 151 (1994)

Study of the **effects** of the $\Lambda N(N)$ interactions & the used many-body model...

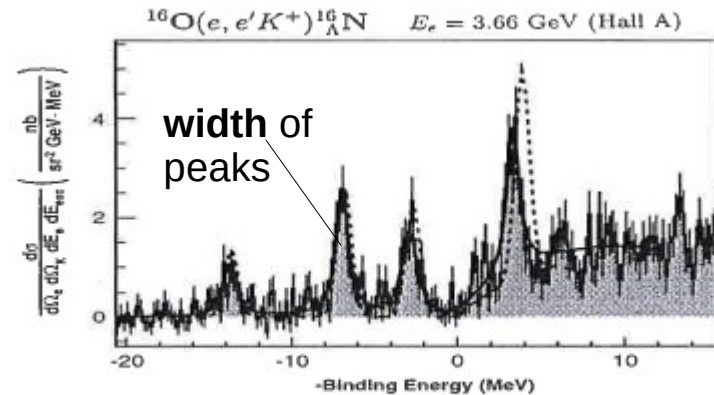
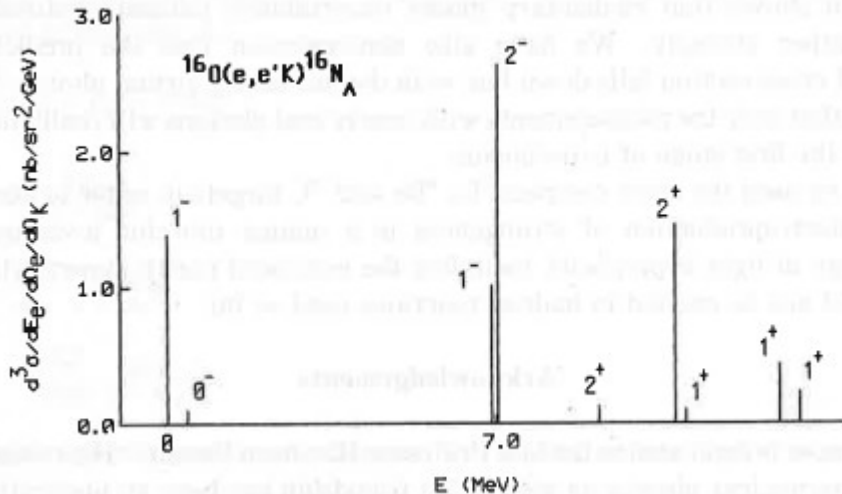


Fig. 17. Experimental spectrum of the $^{16}\text{O}(e, e'K^+)_{\Lambda}^{16}\text{N}$ reaction obtained at JLab Hall A. Taken from Ref. 7). 7) F. Cusanno et al., *Phys. Rev. Lett.* **103** (2009), 202501.

EMPM for Hypernuclei

EMPM extended on single- Λ hypernuclei

hypernuclei with Λ in even-odd nuclear cores

$$\mathcal{H} = \mathcal{H}_0 \oplus \mathcal{H}_1 \oplus \mathcal{H}_2 \oplus \dots \oplus \mathcal{H}_n$$

We couple Λ N TDA states with $|\text{HF}\rangle$ and (multi)phonon excitations of $|\text{HF}\rangle$

$$\begin{aligned}\mathcal{H}_0 &= \{R_\nu^\dagger |\text{HF}\rangle\} \\ \mathcal{H}_1 &= \{R_\nu^\dagger O_{\mu_1}^\dagger |\text{HF}\rangle\} \\ \mathcal{H}_2 &= \{R_\nu^\dagger O_{\mu_1}^\dagger O_{\nu_1}^\dagger |\text{HF}\rangle\}\end{aligned}$$

We can apply this method to calculate structure of ${}^{12}_\Lambda\text{B}$, ${}^{12}_\Lambda\text{C}$, ${}^{16}_\Lambda\text{N}$, ${}^{16}_\Lambda\text{O}$, ${}^{40}_\Lambda\text{K}$, ${}^{40}_\Lambda\text{Ca}$, ${}^{48}_\Lambda\text{K}$, ${}^{48}_\Lambda\text{Ca}$ etc.

Nuclei:	HF	→	TDA	→	EMPM
Hypernuclei:	p-n-Λ HF	→	NΛ TDA	→	ext. EMPM

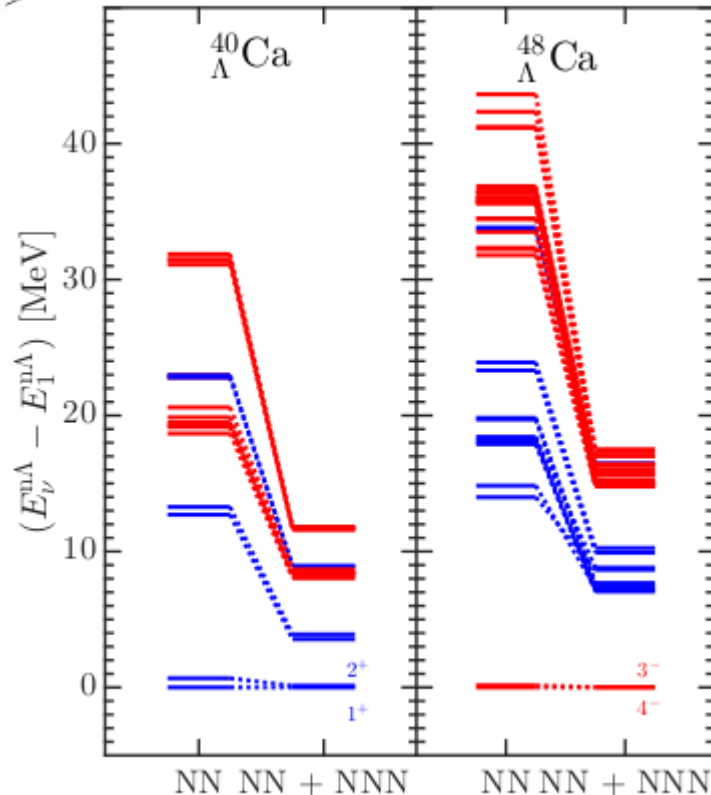
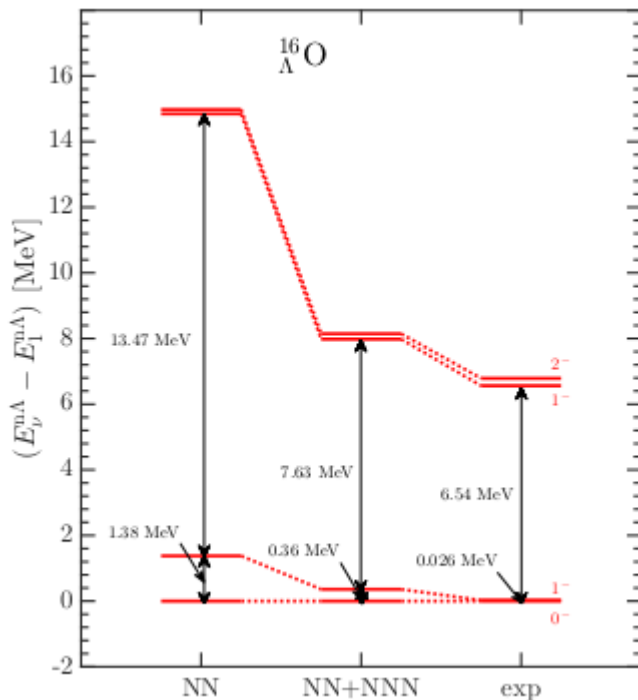
NΛ Tamm-Dancoff

NΛ TDA = Nucleon-Λ Tamm-Dancoff Approximation

- diploma thesis of **J. Pokorný** “*Three-body Interactions in Mean-Field Model of Nuclei and Hypernuclei*”, Czech Technical University, (2018)
- **Phys. Scr. 94, 014006, (2019)**; **Acta Phys. Pol. B Proc. Supl. 12, 657, (2019)**

Suitable for hypernuclei with Λ in even-odd nuclear cores

NΛ TDA Phonon $R_{\nu}^{\dagger} | \text{HF} \rangle = \sum_{\text{ph}} r_{\text{ph}}^{\nu} c_{\text{p}}^{\dagger} a_{\text{h}} | \text{HF} \rangle$



realistic
chiral
NN+NNN
potential
NNLO_{sat}

realistic chiral LO
YN potential
($\Lambda\text{N}-\Lambda\text{N}$ channel)
with different
regulator cut-off λ

EMPM for Hypernuclei

EMPM extended on single- Λ hypernuclei

II) hypernuclei with Λ in even-odd nuclear cores

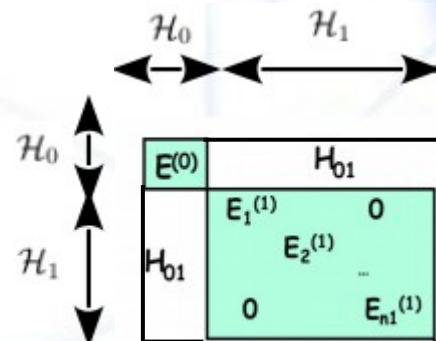
$$\hat{H} = \hat{T}^N + \hat{T}^\Lambda + \hat{V}^{NN} + \hat{V}^{NNN} + \hat{V}^{\Lambda N} + \hat{V}^{\Lambda NN} - \hat{T}_{CM}$$

It is more important to study such hypernuclei from the point of view of experiment (production of hypernuclei $^{12}_\Lambda\text{B}$, $^{16}_\Lambda\text{O}$, $^{16}_\Lambda\text{N}$, $^{40}_\Lambda\text{K}$, $^{48}_\Lambda\text{K}$,...)

Our theoretical formalism:

$$|\nu\rangle = R_\nu^\dagger |\text{HF}\rangle$$

$$|\beta\rangle = \sum_{\nu\mu} X_{\nu\mu}^\beta R_\nu^\dagger Q_\mu^\dagger |\text{HF}\rangle$$



We construct the **Hamiltonian** matrix:

- 1) The diagonal block $\mathcal{H}_0 \times \mathcal{H}_0$ = **N Λ TDA energies**
- 2) The diagonal block $\mathcal{H}_1 \times \mathcal{H}_1$ = **Equation of Motion**
- 3) The nondiagonal block $\mathcal{H}_0 \times \mathcal{H}_1$ => not difficult to calculate

Equation of Motion:

$$AX = EX$$

A-matrix

$$A = \langle \beta | [\hat{H}, R_\nu^\dagger] | \mu \rangle + E_\mu \langle \beta | R_\nu^\dagger | \mu \rangle$$

Eigen-value problem in an overcomplete non-orthogonal basis...

$$\overline{ADC} = E\overline{DC}$$

Eigen-value problem in the reduced space (linearly independent subset of states)

$$-B_\Lambda = E_i + \varepsilon_F^N$$

D-matrix = overlap matrix of the basis states
(**A.D**) – must be hermitian

$$\mathcal{H} = \mathcal{H}_0 \oplus \mathcal{H}_1 \oplus \mathcal{H}_2 \oplus \dots \oplus \mathcal{H}_n$$

$$\mathcal{H}_0 = \{R_\nu^\dagger |\text{HF}\rangle\}$$

$$\mathcal{H}_1 = \{R_\nu^\dagger O_{\mu_1}^\dagger |\text{HF}\rangle\}$$

$$\mathcal{H}_2 = \{R_\nu^\dagger O_{\mu_1}^\dagger O_{\nu_1}^\dagger |\text{HF}\rangle\}$$

Description of Hypernuclei

Hypernuclei with **single- Λ** particle:
Hamiltonian

$$\hat{H} = \hat{T}^N + \hat{T}^\Lambda + \hat{V}^{NN} + \hat{V}^{NNN} + \hat{V}^{\Lambda N} + \hat{V}^{\Lambda NN} - \hat{T}_{CM}$$

Example: NN(+NNN) potential

- phenomenological **NN** potential
Brink-Boeker
 Nucl. Phys. **A91**, 1, (1967)
- realistic chiral **NN** potential
NNLO_{opt}
 Phys. Rev. Lett. **110**, 192502, (2013)
- realistic chiral **NN+NNN** potential
NNLO_{sat}
 Phys. Rev. **C91**, 051301(R), (2015)

ΛN potential

- **G-matrix** effective ΛN potential derived from **Juelich & Nijmegen YN**
 Prog. Theor. Phys. Suppl. **117**, 361 (1994)

$$V_{\Lambda N}(r) = \sum_{i=1}^3 (a_i + b_i k_F + c_i k_F^2) \exp[-r^2/\beta_i^2]$$

Juelich-A YN:

	$\beta_i(\text{fm})$	1.25	0.70	0.45
1E	a	-25.82	-389.4	859.0
	b	-12.51	401.2	-303.2
	c	2.437	-136.0	188.8
3E	a	-45.01	-296.6	1094.
	b	4.620	218.3	-504.6
	c	.7500	-92.50	230.0
1O	a	-14.54	144.7	734.6
	b	3.615	27.50	76.37
	c	-.8750	-5.000	3.125
3O	a	-25.91	248.1	615.3
	b	5.410	210.9	-1260.
	c	.5000	-123.1	734.8

Description of Hypernuclei

- **G-matrix** effective ΛN potential derived from **Juelich & Nijmegen YN**
Prog. Theor. Phys. Suppl. **117**, 361 (1994)
Gaussian-like form – easy to implement, interaction is effective (we can take just ΛN - ΛN part)
but dependent on a parameter k_F

$$V_{\Lambda N}(r) = \sum_{i=1}^3 (a_i + b_i k_F + c_i k_F^2) \exp[-r^2/\beta_i^2]$$

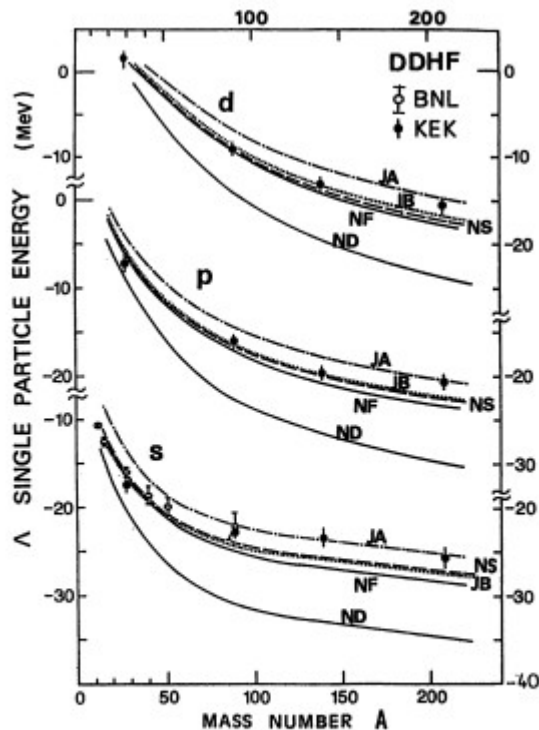


Fig. 2. The Λ single-particle energies $\epsilon_{\Lambda}(nl; A)$, $nl = 0s, 0p$ and $0d$, calculated in DDHF with five effective interactions. The experimental data

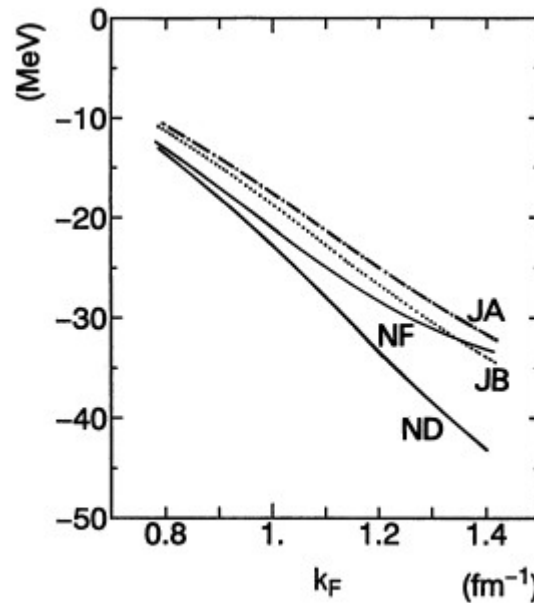


Fig. 1. The Λ single-particle potentials calculated as a function of k_F . The different interaction

Dependence of the Λ **single-particle** energies on k_F

k_F as a parameter to tune the proper effective ΛN interaction. But tuning should be done at the level of the beyond mean-field calculation.

Description of Hypernuclei

- **G-matrix effective ΛN potential derived from Juelich & Nijmegen YN**
Prog. Theor. Phys. Suppl. **117**, 361 (1994)

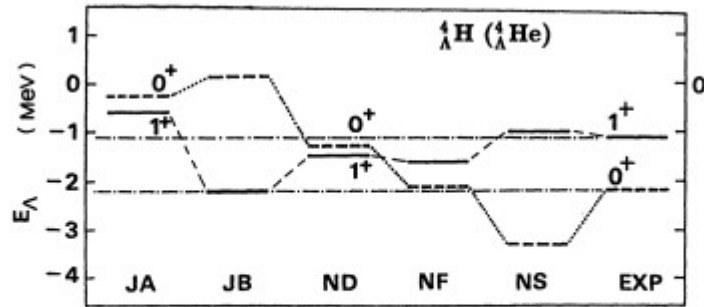
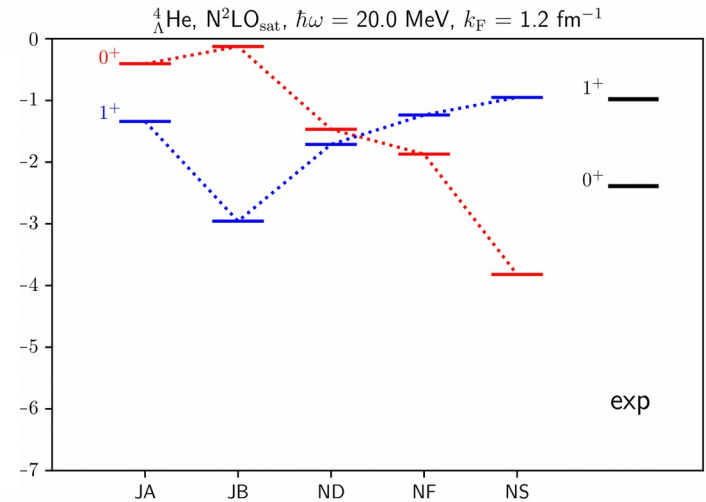


Fig. 3. Energy levels of ${}^4_\Lambda\text{H}$ (${}^4_\Lambda\text{He}$), 0^+ (dashed) and 1^+ (solid). The horizontal dash-dotted lines indicate the experimental positions. The GCM is applied by adopting the size parameter $b=0.5$ fm with five effective interactions $\text{YNG}(k_F=0.8 \text{ fm}^{-1})$.



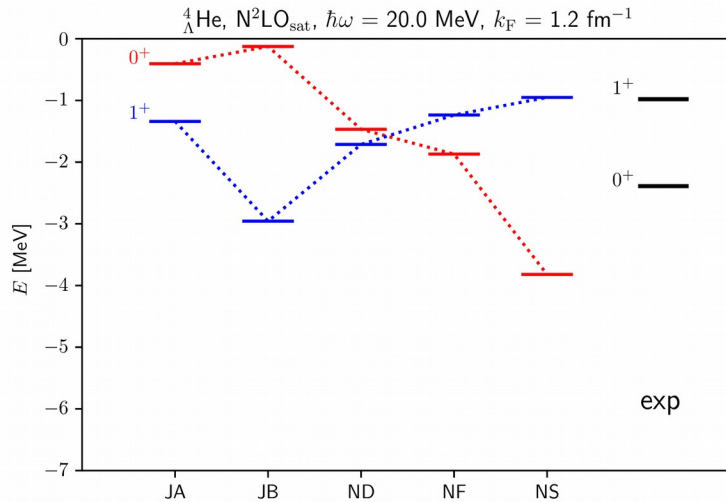
Generator Coordinate Method (GCM)
Prog. Theor. Phys. Suppl. **117**, 361 (1994)

$N\Lambda$ TDA
(Fig. by J. Pokorny)

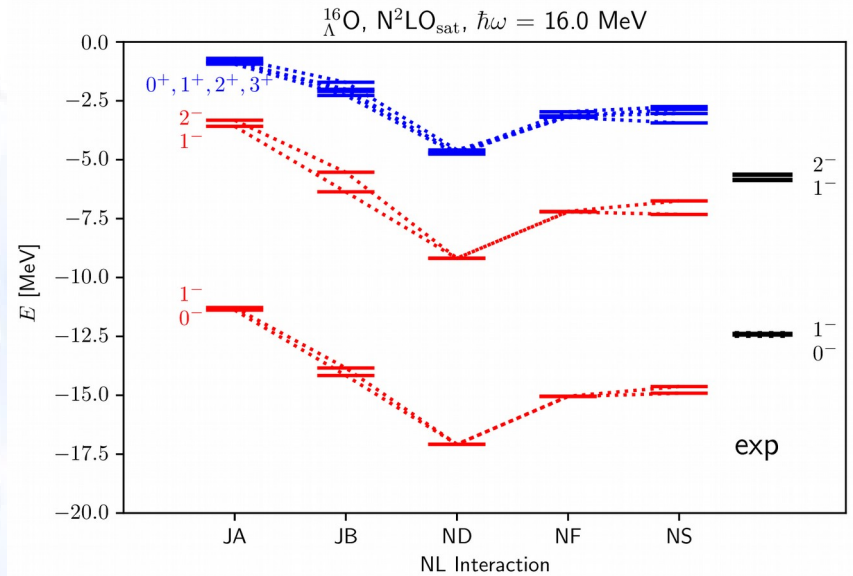
Big variability among different $N\Lambda$ potentials

Description of Hypernuclei

- **G-matrix** effective ΛN potential derived from **Juelich & Nijmegen YN**
Prog. Theor. Phys. Suppl. **117**, 361 (1994)



$N\Lambda$ TDA
(Fig. by J. Pokorny)



$N\Lambda$ TDA
(Fig. by J. Pokorny)

Big variability among different $N\Lambda$ potentials

NF – seems to have good behavior for **$N\Lambda$ TDA**
(if k_F properly tuned)

Convergence of EMPM

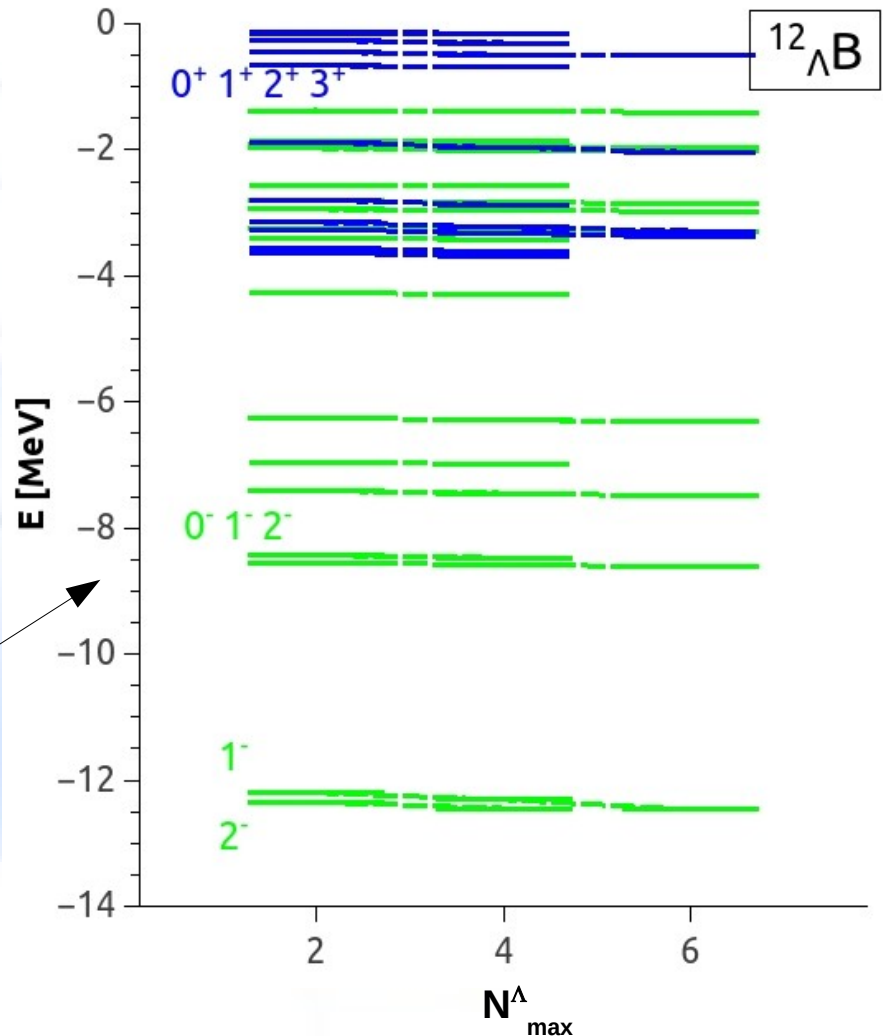
- **G-matrix** effective ΛN potentials “**JA**” and “**NF**” derived from **Juelich** and **Nijmegen YN** Prog. Theor. Phys. Suppl. **117**, 361 (1994)

S.p. basis defined by N_{\max} – number of major oscillator shells ($N = 0, 1, 2, \dots, N_{\max}$)

- For the **p-n- Λ Hartree-Fock** we took N_{\max}
- For the **TDA** we took N_{\max}
- For the **$N\Lambda$ TDA** we took N_{\max}^{Λ}
(in general $N_{\max}^{\Lambda} < N_{\max}$)

It is too computationally costly to apply both **TDA** and **$N\Lambda$ TDA** in the maximal space...

Calculation for $N_{\max} = 10$
Potential **JA** with $k_F = 1.2 \text{ fm}^{-1}$



Convergence of EMPM

- **G-matrix** effective ΔN potentials “**JA**” and “**NF**” derived from **Juelich** and **Nijmegen YN** Prog. Theor. Phys. Suppl. **117**, 361 (1994)

S.p. basis defined by N_{\max} – number of major oscillator shells ($N = 0, 1, 2, \dots, N_{\max}$)

- For the **p-n- Δ Hartree-Fock** we took N_{\max}

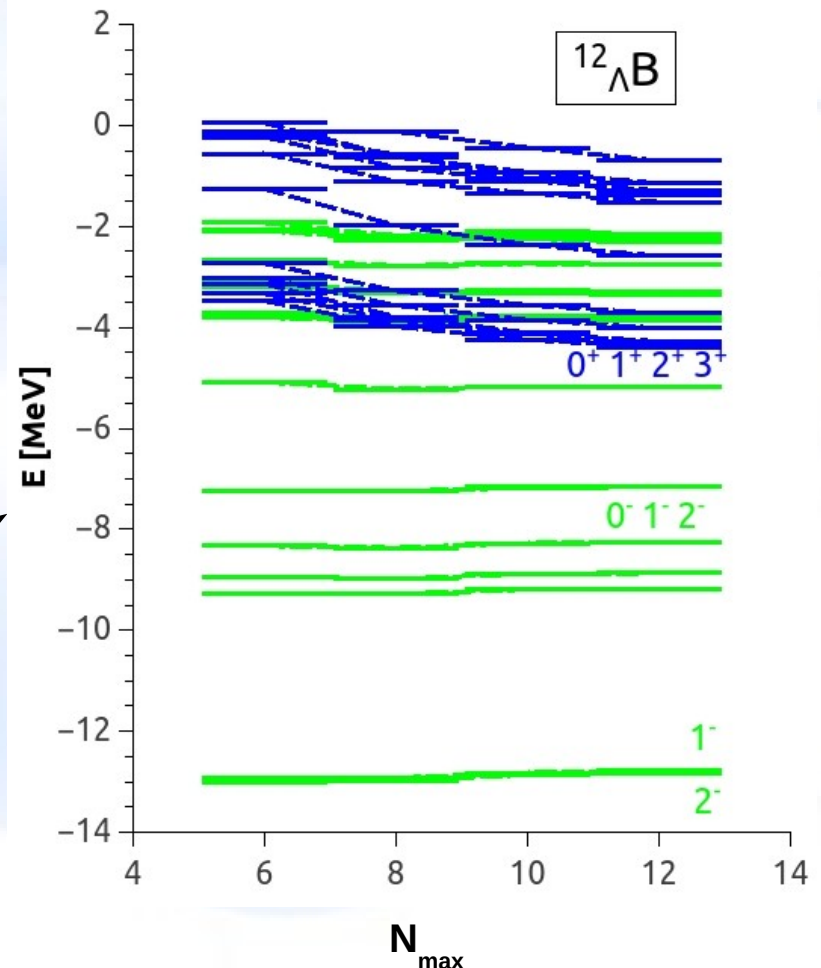
- For the **TDA** we took N_{\max}

- For the **$N\Delta$ TDA** we took N_{\max}^{Δ}
(in general $N_{\max}^{\Delta} < N_{\max}$)

It is too computationally costly to apply both **TDA** and **$N\Delta$ TDA** in the maximal space...

Calculation for $N_{\max}^{\Delta} = 4$

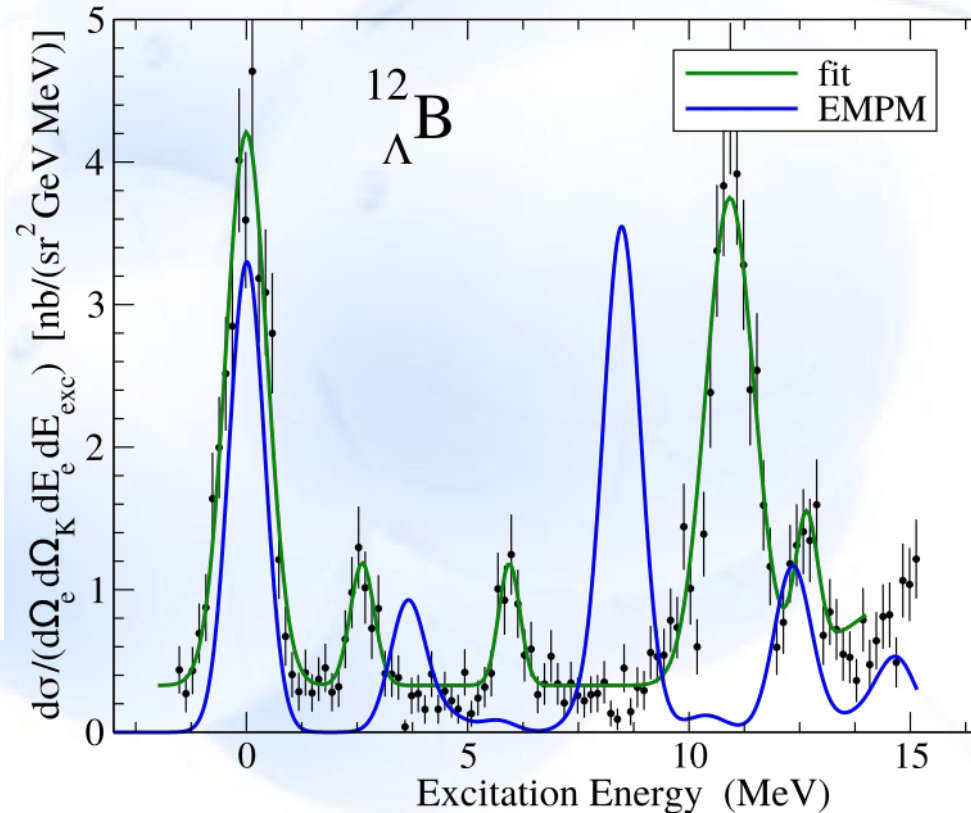
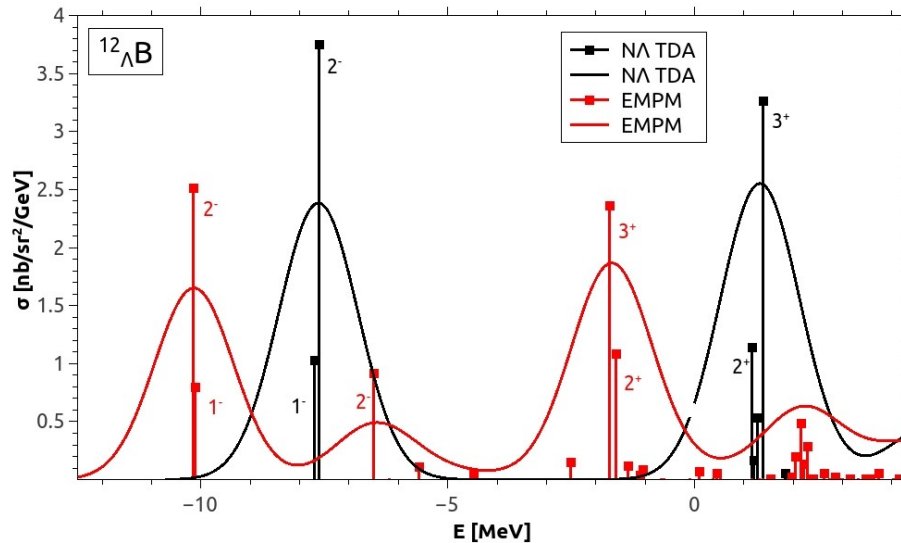
Potential **NF** with $k_F = 1.25 \text{ fm}^{-1}$



Cross section of electroproduction of ${}_{\Lambda}^{12}\text{B}$

- **G-matrix** effective ΛN potentials “**JA**” and “**NF**” derived from **Juelich** and **Nijmegen YN** Prog. Theor. Phys. Suppl. **117**, 361 (1994)

Calculation for $N_{\text{max}} = 12$, $N_{\Lambda}^{\text{max}} = 4$
 Potential **NF** with $k_F = 1.25 \text{ fm}^{-1}$



Summary

- **Extensions** of EMPM to calculate single- Λ **hypernuclei**
- Both nuclear & hypernuclear calcs. useful to study **production of hypernuclei**
- Convergence study of **EMPM (N Λ TDA + TDA)** for $^{12}_{\Lambda}\mathbf{B}$
- Preliminary results of the cross section of the electroproduction of $^{12}_{\Lambda}\mathbf{B}$
- **Tasks to be addressed:**
 - study of $^{40}_{\Lambda}\mathbf{K}$, $^{48}_{\Lambda}\mathbf{K}$, $^{208}_{\Lambda}\mathbf{Tl}$ (their exper. measurement is planned in close future)
 - formalism to study isospin dependence of $\Lambda\mathbf{NN}$ interaction ($^{40}_{\Lambda}\mathbf{K}$ & $^{48}_{\Lambda}\mathbf{K}$)
- More **long-term tasks:**
 - further **development of EMPM** itself (coupling to 2-phonon states)
 - formulation of whole method in **deformed HF** basis

Many thanks to all my collaborators!!

P. Bydžovský, G. De Gregorio, D. Denisova, F. Knapp, N. Lo Iudice, D. Petrellis, J. Pokorný, D. Skoupil