Investigation of Lambda Hypernuclei within Equation of Motion Phonon Method – progress report

Petr Veselý
Nuclear Physics Institute, Czech Academy of Sciences
gemma.ujf.cas.cz/~p.vesely/

Collaboration Japan – Czech Republic,
WEPH2022, October 2022

Motivation

Work on development of **many-body model**(s) suitable for description of **nuclear and hypernuclear structure** which aim to describe wide range of hypernuclei including medium-size & heavy systems.

This model can be used in calculations of **hypernuclear production** – especially the hypernuclei whose production is planned by experimentalists in close future (**$^{40}{}_{\Lambda}$K, $^{48}{}_{\Lambda}$K, $^{208}{}_{\Lambda}$Tl**).

Progress report:

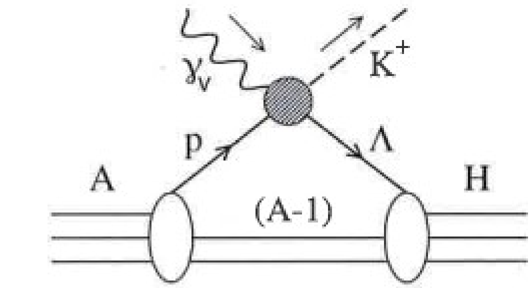
- **Electroproduction** of Hypernuclei
- Applications of **EMPM** for the Calculations of **Hypernuclei**
- **Results** – cross sections of the production of **$^{12}{}_{\Lambda}$B** and preliminary results for **$^{40,48}{}_{\Lambda}$K**
- **Summary** and Future Plans

Electroproduction of Hypernuclei

Hypernuclear production:

$$\gamma_v(P_\gamma) + \mathrm{A}(P_\mathrm{A}) \longrightarrow \mathrm{H}(P_\mathrm{H}) + \mathrm{K}^+(P_\mathrm{K})$$

- Elementary process of electroproduction: **p (e,e‘ K⁺) Λ**
- **Kinematics** of the reaction
- Information about (many-body) **nuclear & hypernuclear structure**
- Information about the **ΛN(N) interactions**



All “ingredients“ important for description of hypernuclear production.

$$\frac{d^3\sigma}{dE'_e d\Omega'_e d\Omega_K} = \Gamma\left[\frac{d\sigma_U}{d\Omega_K} + \varepsilon_L \frac{d\sigma_L}{d\Omega_K} + \varepsilon \frac{d\sigma_P}{d\Omega_K}\cos 2\Phi_K + \sqrt{2\varepsilon_L(1+\varepsilon)}\frac{d\sigma_I}{d\Omega_K}\cos\Phi_K\right]$$

$$\frac{d\sigma_U}{d\Omega_K} = \frac{\beta}{2(2J_A+1)}\sum_{jm}\frac{1}{2j+1}\{|A^{+1}_{jm}|^2 + |A^{-1}_{jm}|^2\},$$

$$\frac{d\sigma_P}{d\Omega_K} = -\frac{\beta}{2J_A+1}\sum_{jm}\frac{1}{2j+1}\mathrm{Re}\{A^{+1}_{jm}A^{-1*}_{jm}\},$$

$$\frac{d\sigma_L}{d\Omega_K} = \frac{\beta}{2J_A+1}\sum_{jm}\frac{1}{2j+1}|A^{0}_{jm}|^2,$$

$$\frac{d\sigma_I}{d\Omega_K} = \frac{\beta}{2J_A+1}\sum_{jm}\frac{1}{2j+1}\mathrm{Re}\{A^{0*}_{jm}[A^{+1}_{jm} - A^{-1}_{jm}]\}$$

Transition amplitude

$$T^{(1)}_\lambda = \frac{Z}{[J_H]}\sum_{S\eta}\mathcal{F}^S_{\lambda\eta}\sum_{LM}\sum_{Jm} C^{Jm}_{LMS\eta}\, C^{J_H M_H}_{J_A M_A Jm}\,(\, J_H \|\, F_{LM}[Y_L \otimes \sigma^S]^J \,\|\, J_A \,)$$

Operator can be expressed in the second quantized form.

We evaluate $(\,\Phi_H \,\|\, [b^+_{\alpha'} \otimes a_\alpha]^J \,\|\, \Phi_A\,)$

nucleus

hypernucleus

annihilation **p**

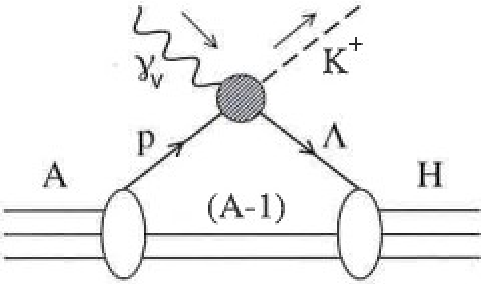
creation Λ

P. Bydžovský et al., **arXiv:2209.00881**

Electroproduction of Hypernuclei

Hypernuclear production:

$$\gamma_v(P_\gamma) + A(P_A) \longrightarrow H(P_H) + K^+(P_K)$$

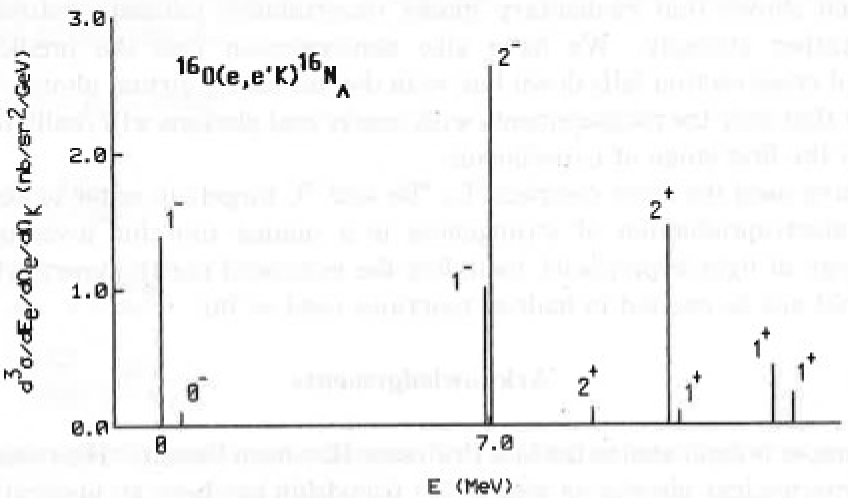


$$\frac{d^3\sigma}{dE'_e d\Omega'_e d\Omega_K} = \Gamma\left[\frac{d\sigma_U}{d\Omega_K} + \varepsilon_L \frac{d\sigma_L}{d\Omega_K} + \varepsilon \frac{d\sigma_P}{d\Omega_K}\cos 2\Phi_K + \sqrt{2\varepsilon_L(1+\varepsilon)}\frac{d\sigma_I}{d\Omega_K}\cos\Phi_K\right]$$

$$M_\mu = \langle \Psi_H | \langle \chi_K | \sum_{j=1}^{Z} \hat{J}_\mu(j) \, | \chi_\gamma \rangle | \Psi_A \rangle \qquad (\, \Phi_H \, || \, [b^+_{\alpha'} \otimes a_\alpha]^J \, || \, \Phi_A \,)$$

Fig. from M. **Sotona**, S. **Frullani**,
Prog. Theor. Phys. Suppl. 117, 151 (1994)

Study of the **effects** of the **ΛN(N) interactions** & the used **many-body model**...



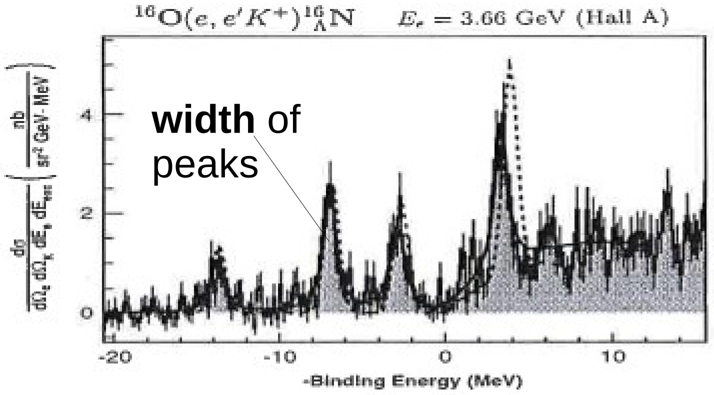


Fig. 17. Experimental spectrum of the $^{16}O(e,e'K^+)^{16}_{\Lambda}N$ reaction obtained at JLab Hall A. Taken from Ref. 7).

7) F. Cusanno et al., Phys. Rev. Lett. **103** (2009), 202501.

EMPM for Hypernuclei

EMPM extended on single-Λ hypernuclei

hypernuclei with Λ in **even-odd nuclear cores**

$$\mathcal{H} = \mathcal{H}_0 \oplus \mathcal{H}_1 \oplus \mathcal{H}_2 \oplus \ldots \oplus \mathcal{H}_n$$

We **couple ΛN TDA** states with |HF> and **(multi)phonon** excitations of |HF>

$$\begin{aligned}
\mathcal{H}_0 &= \{R^\dagger_\nu |\mathrm{HF}>\} \\
\mathcal{H}_1 &= \{R^\dagger_\nu O^\dagger_{\mu_1} |\mathrm{HF}>\} \\
\mathcal{H}_2 &= \{R^\dagger_\nu O^\dagger_{\mu_1} O^\dagger_{\nu_1} |\mathrm{HF}>\}
\end{aligned}$$

We can apply this method to calculate structure of $^{12}_{\Lambda}$**B**, $^{12}_{\Lambda}$**C**, $^{16}_{\Lambda}$**N**, $^{16}_{\Lambda}$**O**, $^{40}_{\Lambda}$**K**, $^{40}_{\Lambda}$**Ca**, $^{48}_{\Lambda}$**K**, $^{48}_{\Lambda}$**Ca** etc.

Nuclei:	**HF**	→	**TDA**	→	**EMPM**
Hypernuclei:	**p-n-Λ HF**	→	**NΛ TDA**	→	**ext. EMPM**

EMPM for Hypernuclei

EMPM extended on single-Λ hypernuclei

II) hypernuclei with Λ in **even-odd nuclear cores**

$$\hat{H} = \hat{T}^N + \hat{T}^\Lambda + \hat{V}^{NN} + \hat{V}^{NNN} + \hat{V}^{\Lambda N} + \hat{V}^{\Lambda NN} - \hat{T}_{CM}$$

$$\mathcal{H} = \boxed{\mathcal{H}_0 \oplus \mathcal{H}_1} \oplus \mathcal{H}_2 \oplus ... \oplus \mathcal{H}_n$$

It is more important to study such hypernuclei from the point of view of experiment (production of hypernuclei **$^{12}_{\Lambda}$B, $^{16}_{\Lambda}$O, $^{16}_{\Lambda}$N, $^{40}_{\Lambda}$K, $^{48}_{\Lambda}$K,...**)

$$\boxed{\begin{aligned} \mathcal{H}_0 &= \{R^\dagger_\nu|\mathrm{HF}>\} \\ \mathcal{H}_1 &= \{R^\dagger_\nu O^\dagger_{\mu_1}|\mathrm{HF}>\} \end{aligned}}$$
$$\mathcal{H}_2 = \{R^\dagger_\nu O^\dagger_{\mu_1} O^\dagger_{\nu_1}|\mathrm{HF}>\}$$

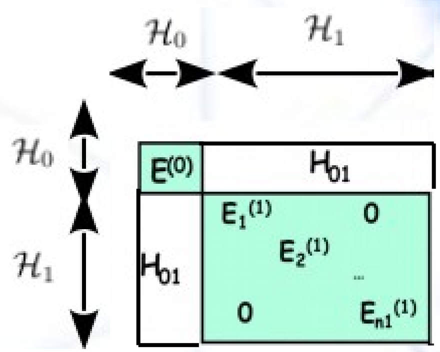
Our theoretical **formalism**:

$$|\nu>= R^\dagger_\nu|\mathrm{HF}>$$

$$|\beta>= \sum_{\nu\mu} X^\beta_{\nu\mu} R^\dagger_\nu Q^\dagger_\mu |\mathrm{HF}>$$

We construct the **Hamiltonian** matrix:

1) The diagonal block $\mathcal{H}_0 \times \mathcal{H}_0$ = **NΛ TDA** energies
2) The diagonal block $\mathcal{H}_1 \times \mathcal{H}_1$ = **Equation of Motion**
3) The nondiagonal block $\mathcal{H}_0 \times \mathcal{H}_1$ => not difficult to calculate



Equation of Motion:

$$\mathcal{A}X = EX$$

A-matrix $\mathcal{A} =< \beta|[\hat{H}, R^\dagger_\nu]|\mu> + E_\mu < \beta|R^\dagger_\nu|\mu>$

Eigen-value problem in an **overcomplete non-orthogonal basis**…

$$\overline{\mathcal{A}\mathcal{D}}C = E\overline{\mathcal{D}}C$$

Eigen-value problem in the **reduced space** (linearly independent subset of states)

$-B_\Lambda = E_i + \varepsilon_F^N$

D-matrix = overlap matrix of the basis states
(A.D) – must be hermitian

Description of Hypernuclei

Hypernuclei with **single-Λ** particle:
Hamiltonian

$$\widehat{H} = \widehat{T}^{N} + \widehat{T}^{\Lambda} + \widehat{V}^{NN} + \widehat{V}^{NNN} + \widehat{V}^{\Lambda N} + \widehat{V}^{\Lambda NN} - \widehat{T}_{CM}.$$

Example:

NN(+NNN) potential

- phenomenological **NN** potential
  Brink-Boeker
  Nucl. Phys. **A91**, 1, (1967)

- realistic chiral **NN** potential
  NNLO$_{\text{opt}}$
  Phys. Rev. Lett. **110**, 192502, (2013)

- realistic chiral **NN+NNN** potential
  NNLO$_{\text{sat}}$
  Phys. Rev. **C91**, 051301(R), (2015)

ΛN potential

- **G-matrix** effective **ΛN** potential derived from **Juelich & Nijmegen YN**
  Prog. Theor. Phys. Suppl. **117**, 361 (1994)

$$V_{\Lambda N}(r) = \sum_{i=1}^{3}(a_i + b_i k_F + c_i k_F^2)\exp[-r^2/\beta_i^2]$$

Juelich-A YN:

	β_i(fm)	1.25	0.70	0.45
	a	−25.82	−389.4	859.0
^1E	b	−12.51	401.2	−303.2
	c	2.437	−136.0	188.8
	a	−45.01	−296.6	1094.
^3E	b	4.620	218.3	−504.6
	c	.7500	−92.50	230.0
	a	−14.54	144.7	734.6
^1O	b	3.615	27.50	76.37
	c	−.8750	−5.000	3.125
	a	−25.91	248.1	615.3
^3O	b	5.410	210.9	−1260.
	c	.5000	−123.1	734.8

Description of Hypernuclei

- **G-matrix** effective **ΛN** potential derived from **Juelich & Nijmegen YN**
  Prog. Theor. Phys. Suppl. **117**, 361 (1994)
  Gaussian-like form – easy to implement, interaction is effective (we can take just ΛN-ΛN part) but dependent on a parameter $\mathbf{k_F}$

$$V_{\Lambda N}(r)=\sum_{i=1}^{3}(a_i+b_ik_F+c_ik_F^2)\exp[-r^2/\beta_i^2]$$

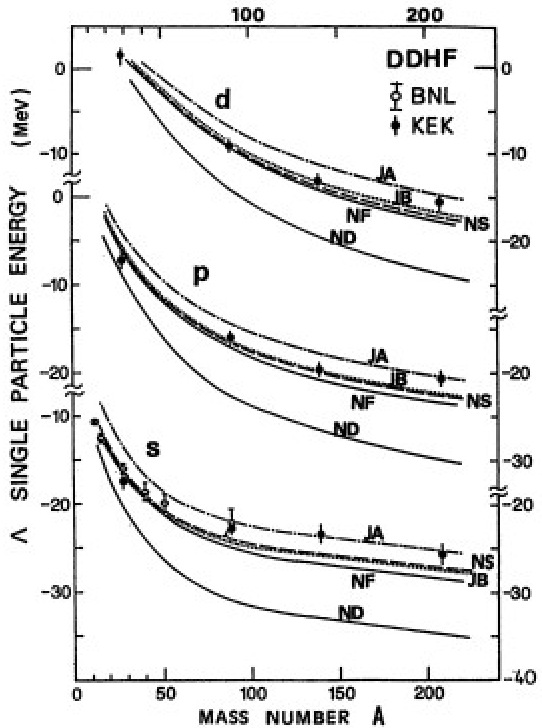


Fig. 2. The Λ single-particle energies $\epsilon_\Lambda(nl; A)$, nl $=0s$, $0p$ and $0d$, calculated in DDHF with five effective interactions. The experimental data

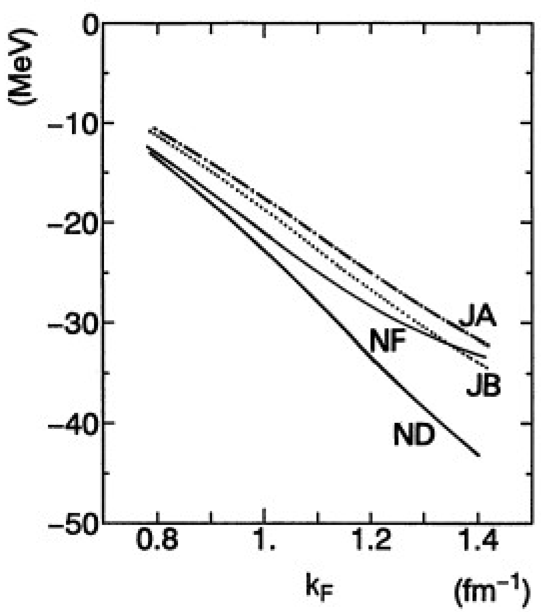


Fig. 1. The Λ single-particle potentials calculated as a function of k_F. The different interaction

Dependence of the **Λ single-particle** energies on $\mathbf{k_F}$

$\mathbf{k_F}$ as a parameter to tune the proper effective ΛN interaction. But tuning should be done at the level of the beyond mean-field calculation.

Prog. Theor. Phys. Suppl. **117**, 361 (1994)

Description of Hypernuclei

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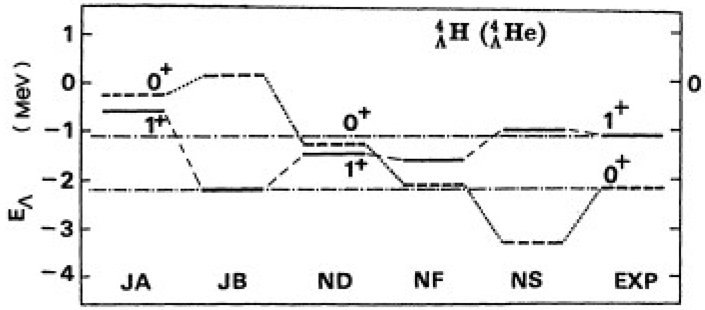
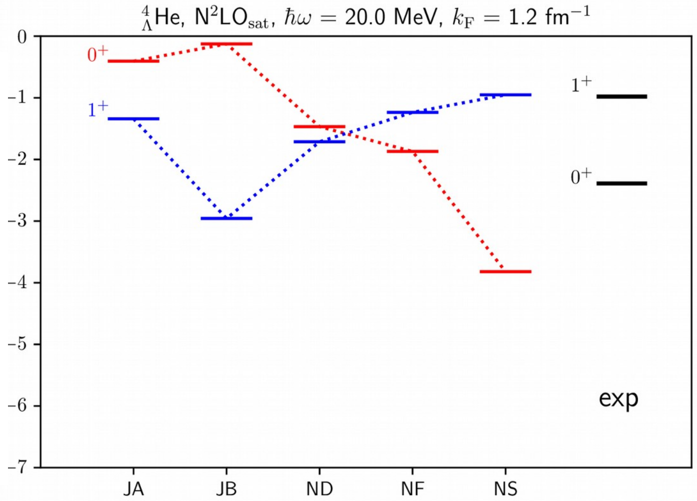


Fig. 3. Energy levels of $^4_\Lambda$H($^4_\Lambda$He), 0^+(dashed) and 1^+(solid). The holizontal dash-dotted lines indicate the experimental positions. The GCM is applied by adopting the size parameter b=0.5 fm with five effective interactions YNG(k_F=0.8 fm^{-1}).



Generator Coordinate Method (GCM)
Prog. Theor. Phys. Suppl. **117**, 361 (1994)

NΛ TDA
(Fig. by J. Pokorny)

Big variability among different NΛ potentials

Description of Hypernuclei

- **G-matrix** effective **ΛN** potential derived from **Juelich & Nijmegen YN**
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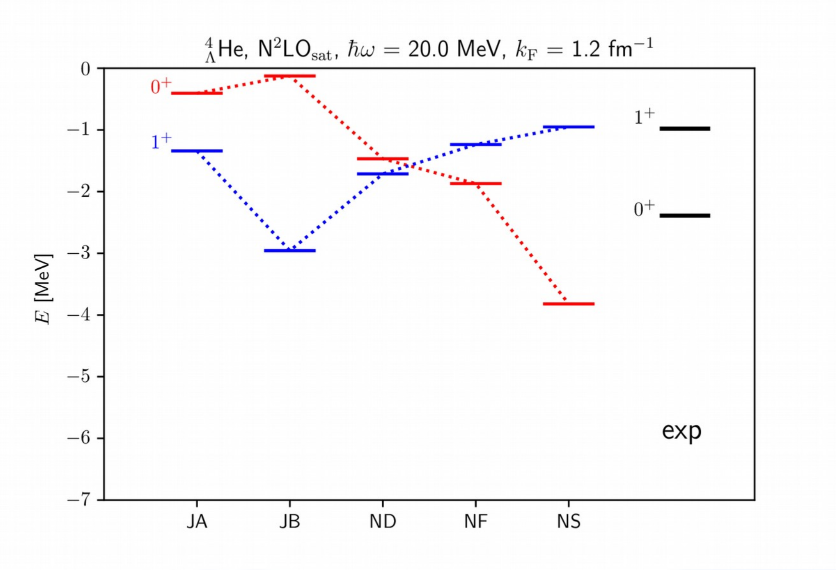
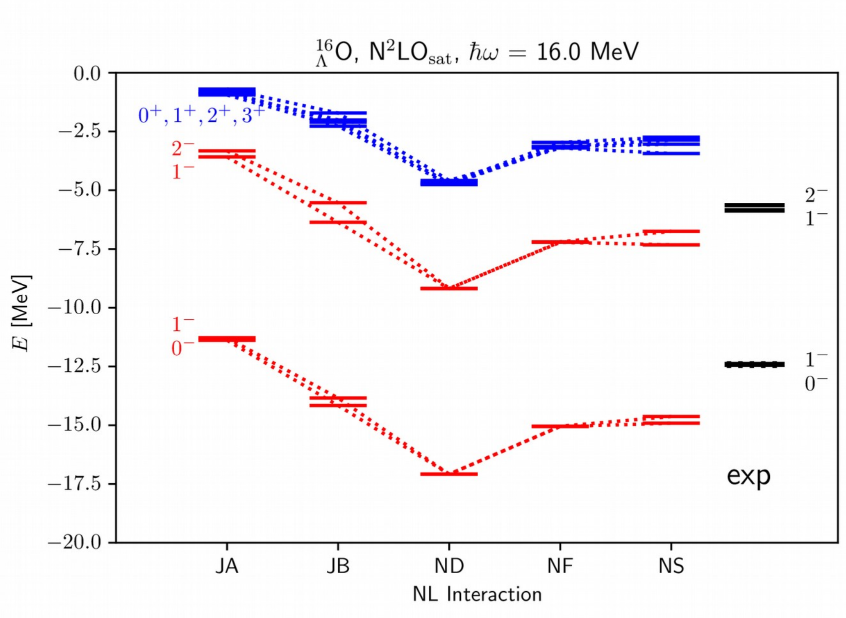


NΛ TDA
(Fig. by J. Pokorny)



NΛ TDA
(Fig. by J. Pokorny)

Big variability among different NΛ potentials

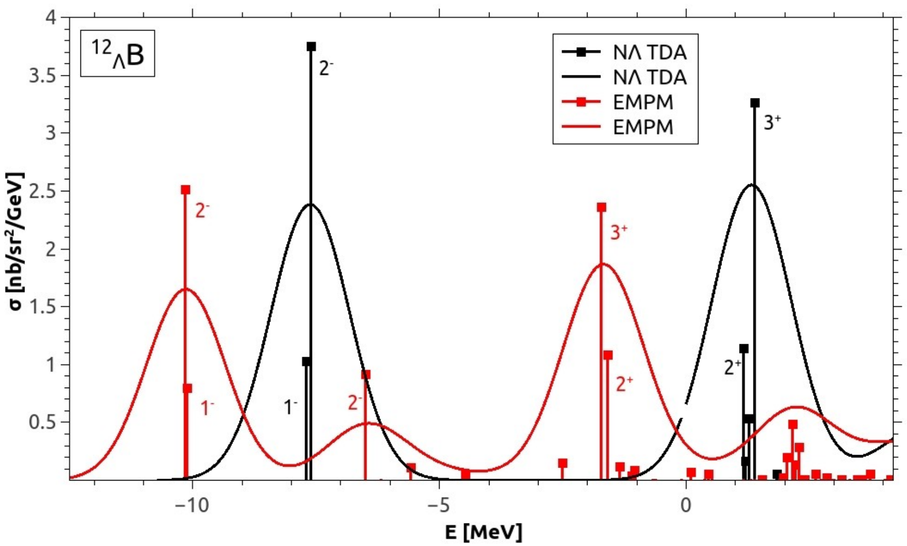
NF – seems to have good behavior for **NΛ TDA**
(if $\mathbf{k_F}$ properly tuned)

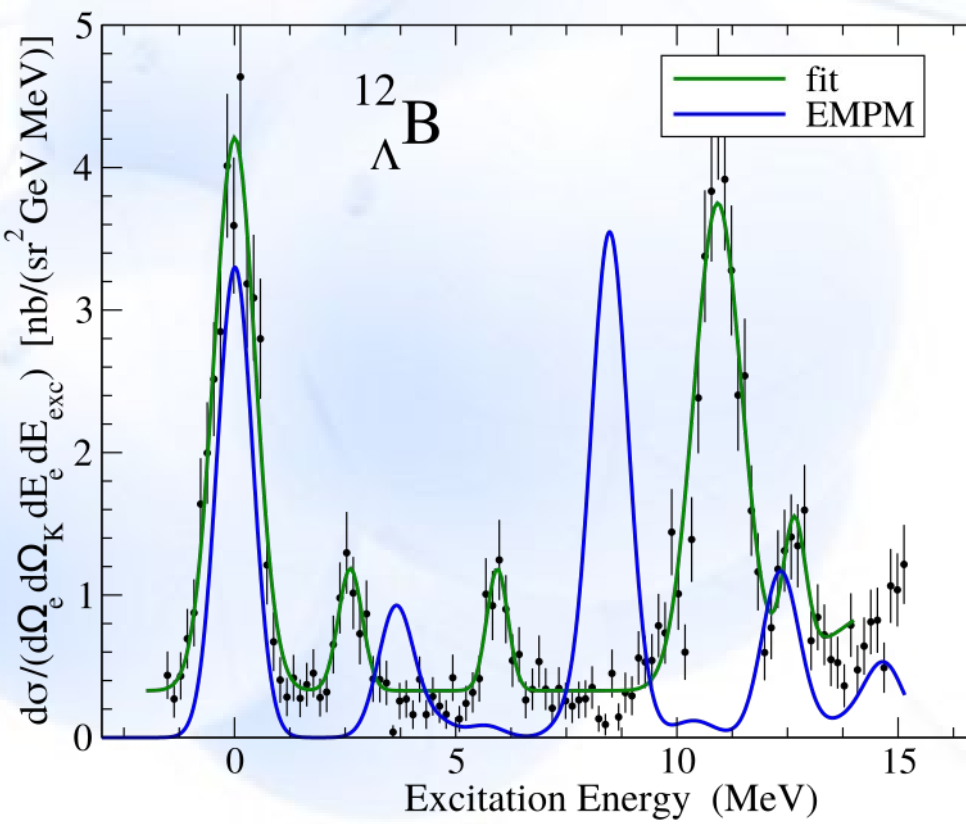
Cross section of electroproduction of $^{12}_{\Lambda}\mathrm{B}$

- **G-matrix** effective **ΛN** potentials “**NF**“ derived from **Juelich** and **Nijmegen YN**
  Prog. Theor. Phys. Suppl. **117**, 361 (1994)

Calculation for $\mathbf{N}_{\mathbf{max}}$ =12 , $\mathbf{N}^{\Lambda}_{\mathbf{max}}$ = 4

Potential **NF** with k_F = 1.25 fm^{-1}





- **EMPM** with respect to **NΛ TDA** has **richer spectrum** (at least two additional peaks), crucial for description of experimental results
- However, the detailed cross section should be improved. We can **tune parameters** of interaction or **add a phenom. term**

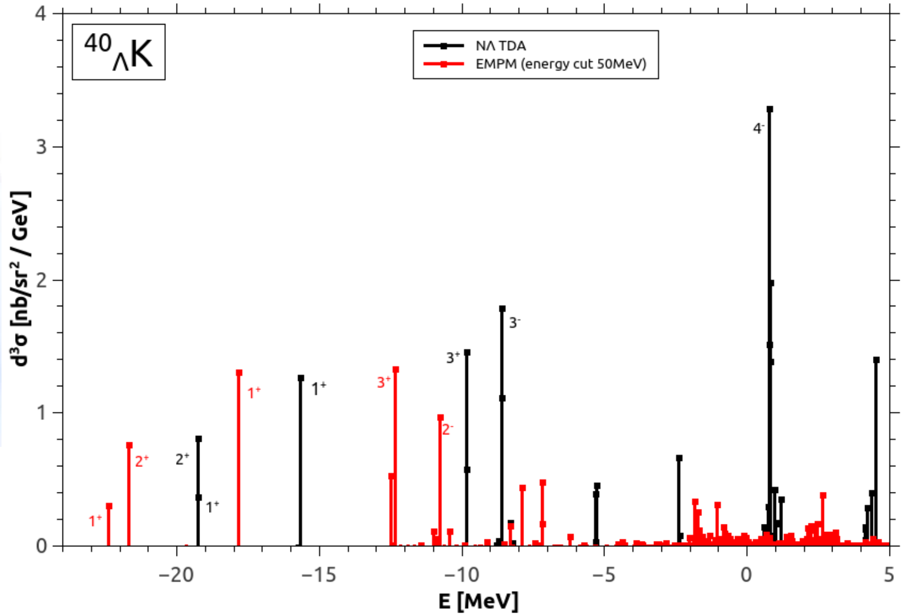
Cross section of electroproduction of ${}^{40,48}_{\Lambda}K$

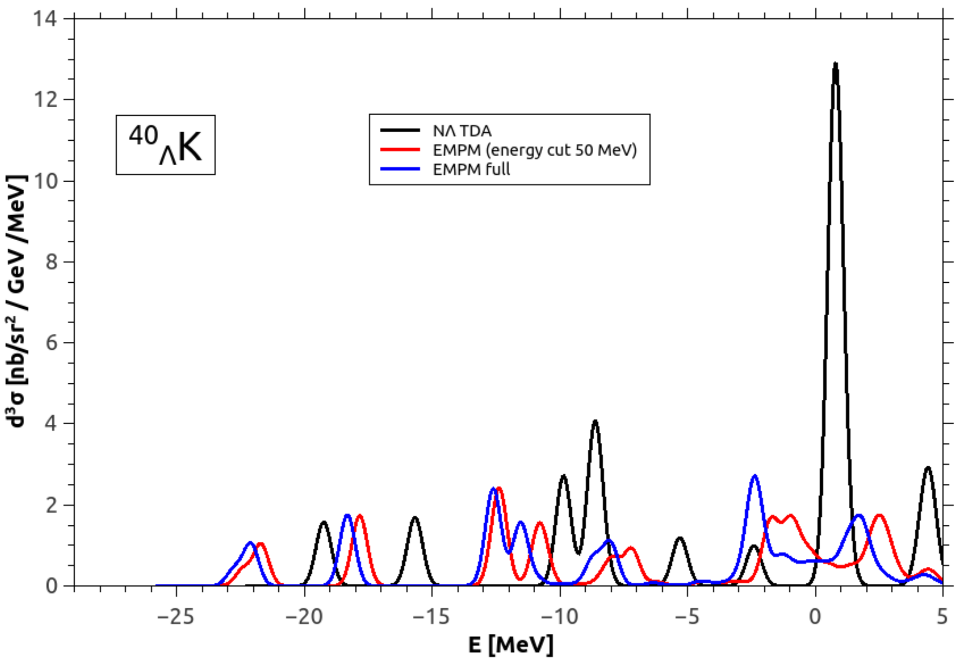
- **G-matrix** effective **ΛN** potentials "**NF**" derived from **Juelich** and **Nijmegen YN**
  Prog. Theor. Phys. Suppl. **117**, 361 (1994)

Calculation for $\mathbf{N_{max}}$ =12 , $\mathbf{N^{\Lambda}_{max}}$ = 4

Potential **NF** with k_F = 1.25 fm^{-1}

Preliminary results!
EMPM frgments the strength to small peaks.
multifragmentation





We test reasonable approximations to make calculations more feasible.

$\mathbf{N^{\Lambda}_{max}} < \mathbf{N_{max}}$

Energy cut-off < 50 MeV.

Summary

- **Extensions** of EMPM to calculate single-Λ **hypernuclei**

- Both nuclear & hypernuclear calcs. useful to study **production of hypernuclei**

- Preliminary results of the cross section of the electroproduction of $^{12}_{\Lambda}\mathbf{B}$, $^{40,48}_{\Lambda}\mathbf{K}$

- **Tasks to be addressed**:

  - study of $^{40}_{\Lambda}\mathbf{K}$, $^{48}_{\Lambda}\mathbf{K}$, $^{208}_{\Lambda}\mathbf{Tl}$ (their exper. measurement is planned in close future)

  - formalism to study isospin dependence of **ΛNN** interaction ($^{40}_{\Lambda}\mathbf{K}$ & $^{48}_{\Lambda}\mathbf{K}$)

- More **long-term tasks**:

  - further **development of EMPM** itself (coupling to 2-phonon states)

  - formulation of whole method in **deformed HF** basis

  - study of **structure Ξ hypernuclei**

Many thanks to all my collaborators!!
P. Bydžovský, G. De Gregorio, D. Denisova, F. Knapp, N. Lo Iudice, D. Petrellis, J. Pokorný, D. Skoupil