

JLab meeting (Pb target)

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Nov 15, 2023



GRADUATE
SCHOOL OF
FACULTY OF **SCIENCE**

KYOTO UNIVERSITY

Required beam time simulation (LOI12-23-013)

https://researchmap.jp/gogami/published_papers/42361620/attachment_file.pdf

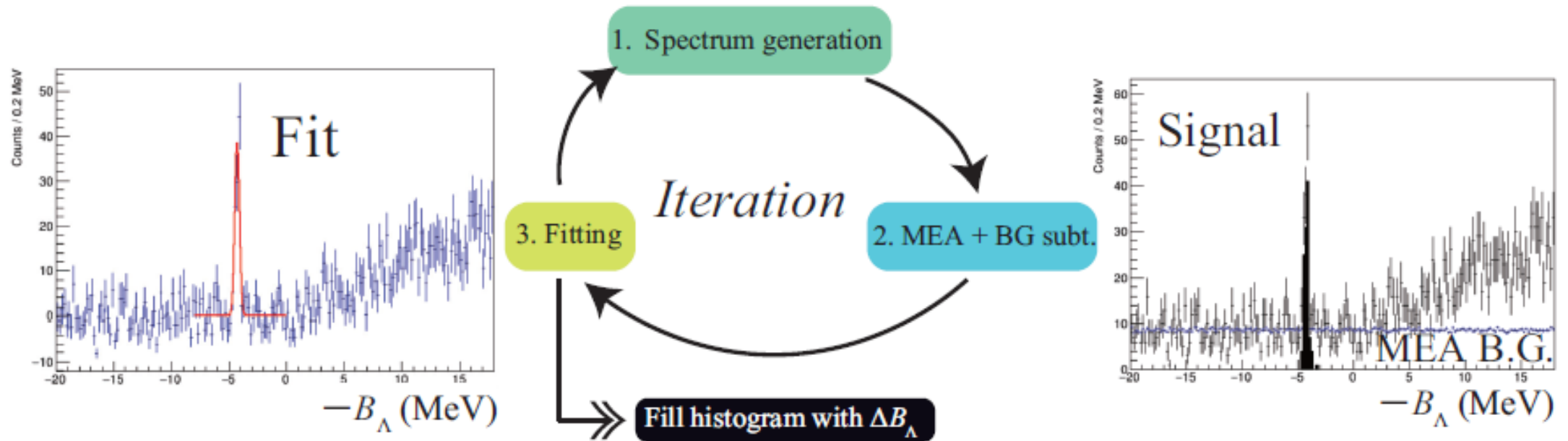


FIG. 8. Flow chart of the simulation to evaluate the required beam times.

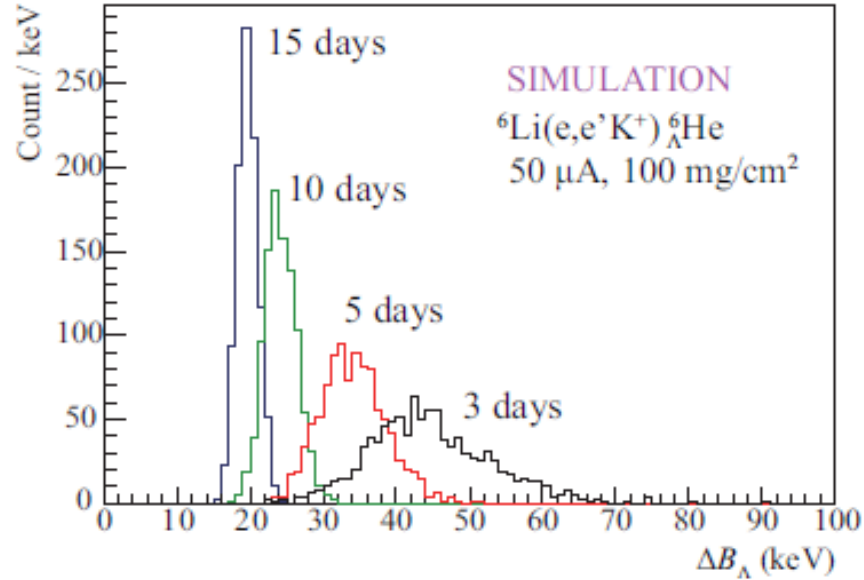


FIG. 9. Statistical uncertainties ($\Delta B_{\Lambda}^{\text{stat.}}$) obtained in the simulation for the ${}^6\text{Li}(e, e'K^+){}^6_{\Lambda}\text{He}$ reaction. The assumed beam current and target thickness are $50 \mu\text{A}$ and 100 mg/cm^2 , respectively. Peak fits to simulated spectra were performed 1000 times for each beam-time condition (3, 5, 10, and 15 days), and the obtained statistical uncertainties are filled in the histogram.

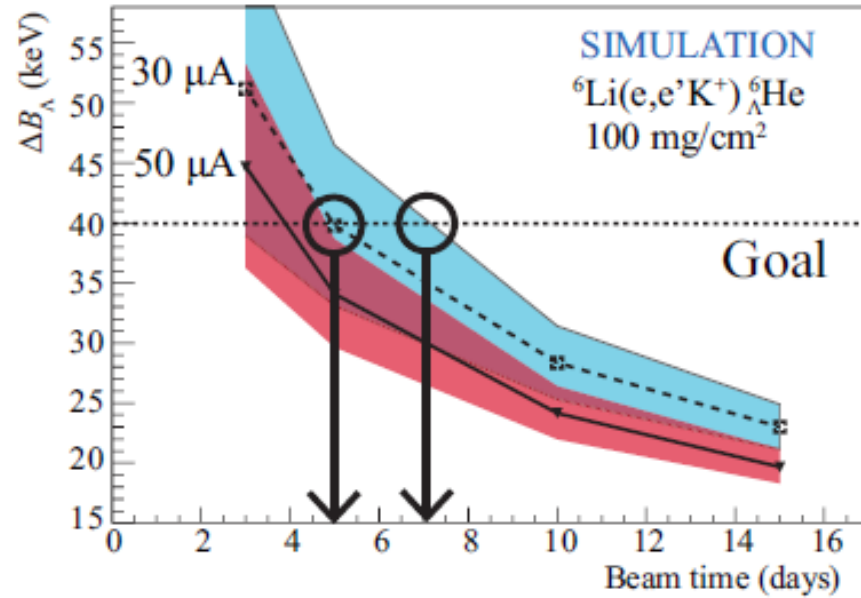
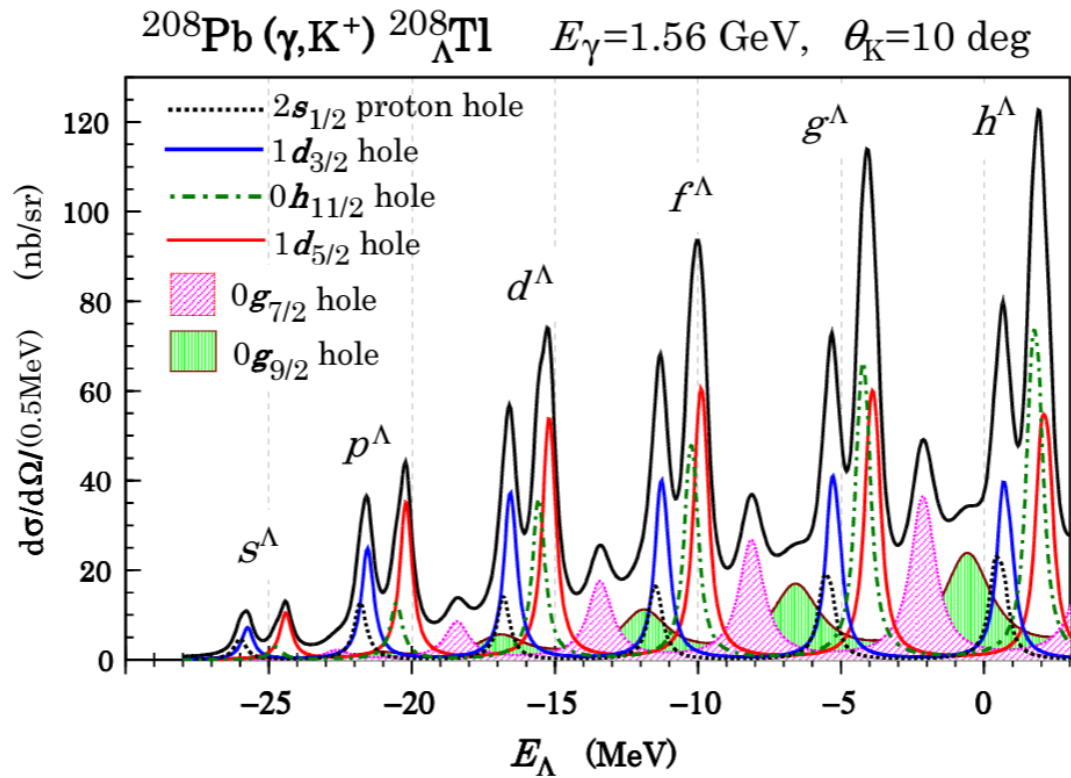


FIG. 10. Mean value of the statistical uncertainty on B_{Λ} (Fig. 9) as a function of a beam time for the ${}^6\text{Li}(e, e'K^+){}^6_{\Lambda}\text{He}$ reaction. Bands colored by blue and red correspond to standard deviations for the cases of 30- and 50- μA beam currents, respectively. Necessary beam times, which meet the goal uncertainty $|\Delta B_{\Lambda}^{\text{stat.}}| = 40 \text{ keV}$, are 5 and 7 days for the beam currents of 30 and 50 μA , respectively.

The same simulation
was applied to the Pb
target

Assumption for the simulation ($^{208}\text{Pb} \rightarrow ^{208}_{\Lambda}\text{Tl}$)



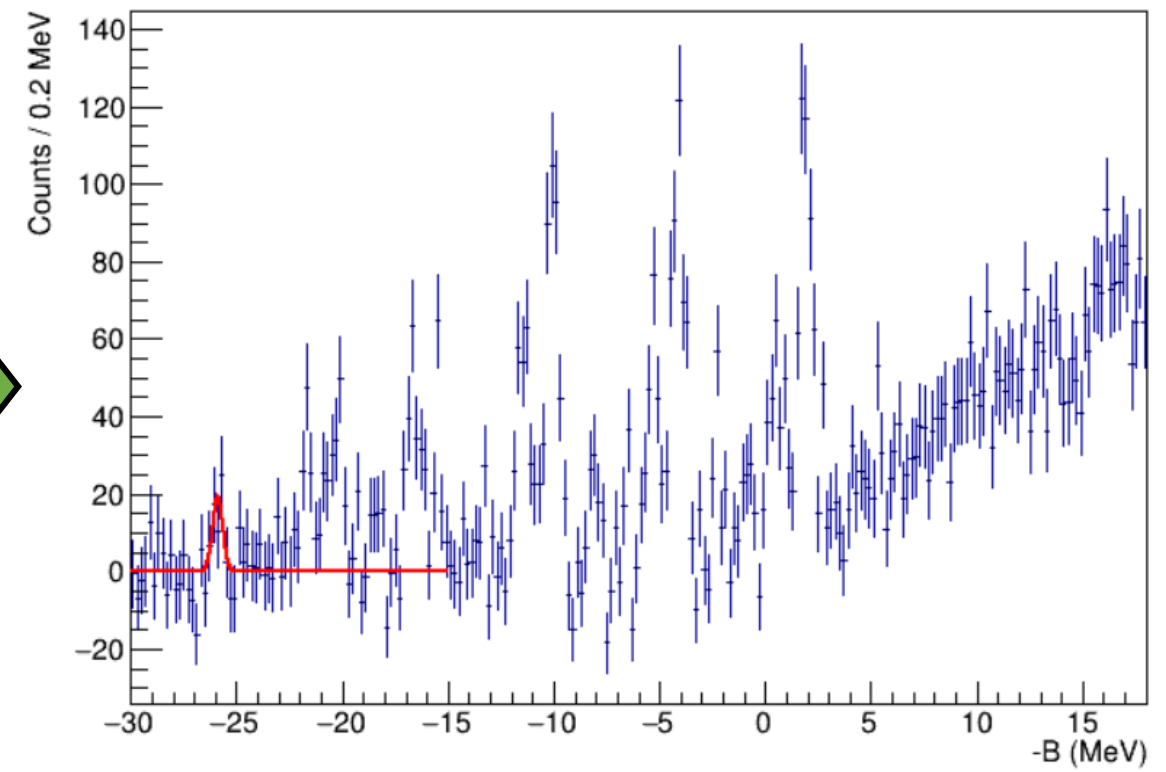
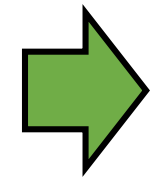
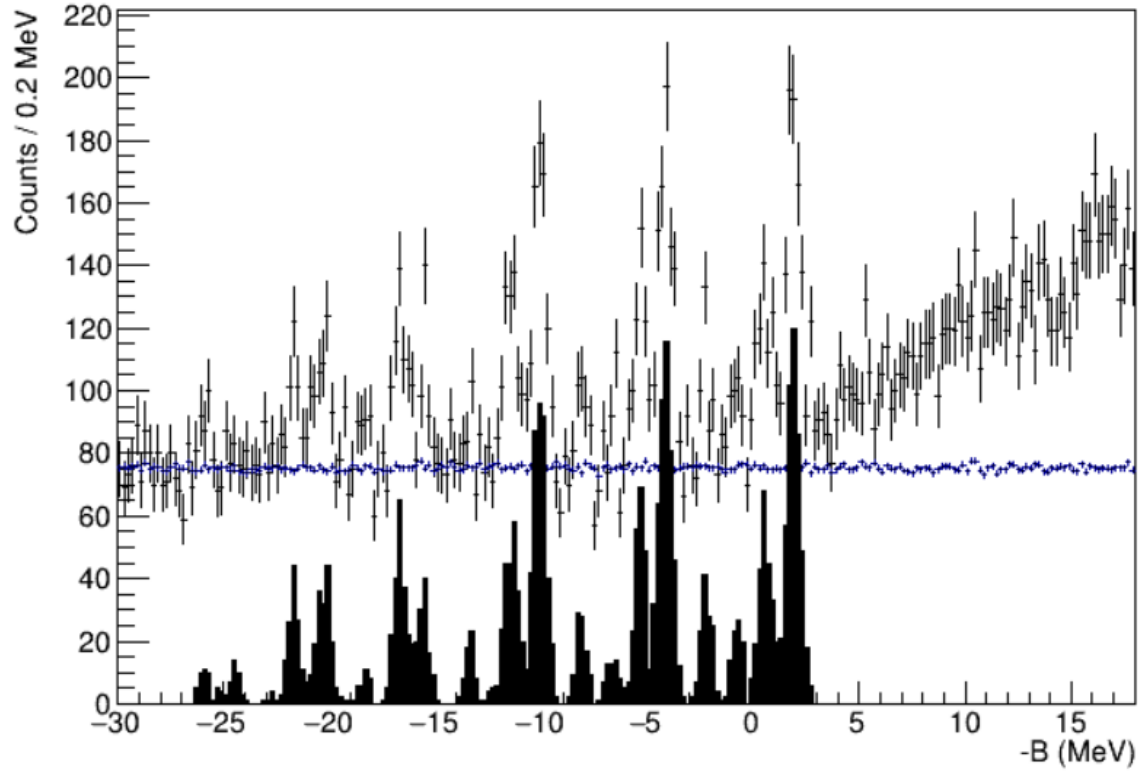
T. Motoba, JPS Conf. Proc. 17, 011003 (2017);
<https://doi.org/10.7566/JPSCP.17.011003>

Table III. Calculated cross sections for the $^{208}\text{Pb}(\gamma, K^+) ^{208}_{\Lambda}\text{Tl}$ reaction at $E_{\gamma}=1.56$ GeV and $\theta_K^{\text{lab}}=5$ and 10 deg. Each entry (nb/sr) is for the 1p-1h multiplet $[(nlj)_p^{-1}(nlj)_{\Lambda}]_J$ summed up over $j^{\Lambda} = \ell \pm 1/2$ and J .

	$\theta_K = 5^{\circ}$							$\theta_K = 10^{\circ}$				
	$0s^{\Lambda}$	$0p^{\Lambda}$	$0d^{\Lambda}$	$1s^{\Lambda}$	$0f^{\Lambda}$	$1p^{\Lambda}$	$(0g^{\Lambda})$	$0s^{\Lambda}$	$0p^{\Lambda}$	$0d1s^{\Lambda}$	$0f1p^{\Lambda}$	$0g1d2s^{\Lambda}$
$(2s_{1/2})_p^{-1}$	15.2	18.2	21.6	2.1	38.8	1.09	31.2	3.6	10.2	11.4	14.5	20.0
$(1d_{3/2})_p^{-1}$	19.5	43.9	54.0	4.7	46.8	16.7	37.6	5.4	19.6	30.9	34.8	37.9
$(0h_{11/2})_p^{-1}$	2.2	14.8	28.9	8.8	44.2	20.7	49.9	2.3	10.1	29.3	41.7	60.5
$(1d_{5/2})_p^{-1}$	29.0	64.0	80.7	6.9	69.0	24.9	46.5	7.9	28.0	45.2	53.8	55.9
$(0g_{7/2})_p^{-1}$	5.2	22.2	36.2	9.6	46.2	27.1	49.0	3.1	12.9	27.2	41.9	58.2
$(0g_{9/2})_p^{-1}$	6.5	26.6	44.9	11.9	55.0	32.8	60.9	3.8	15.1	31.5	49.1	70.6

Sample figures ($20 \mu A$, 150 mg/cm^2 , 20 days)

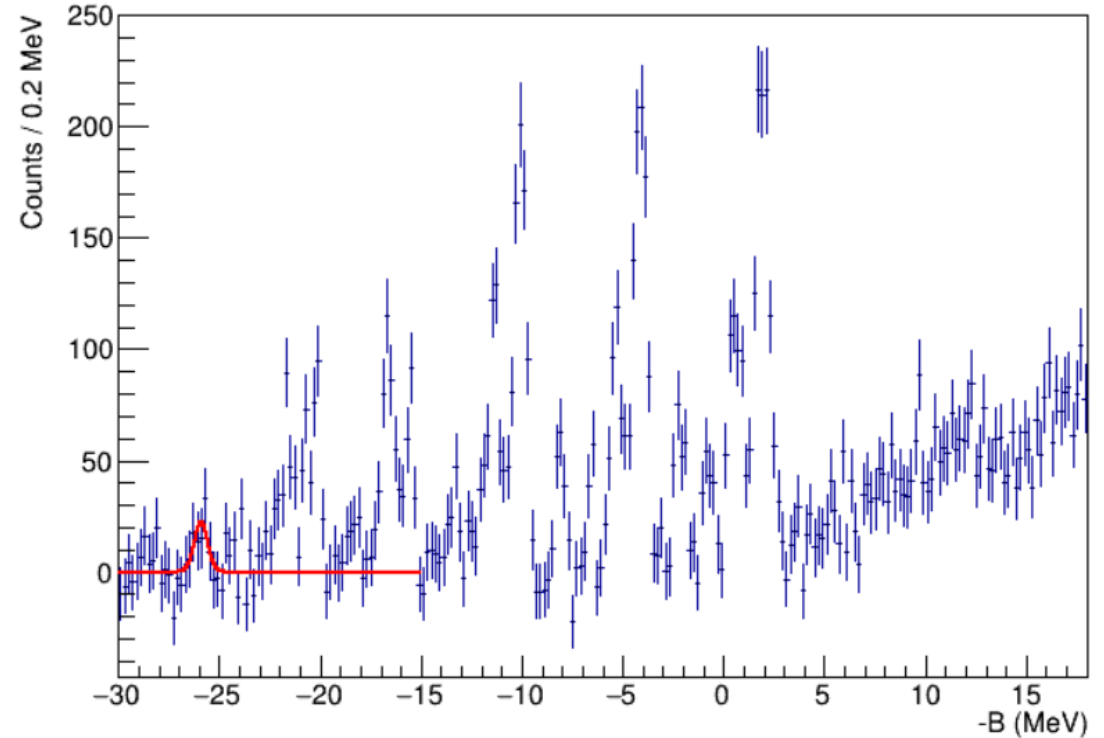
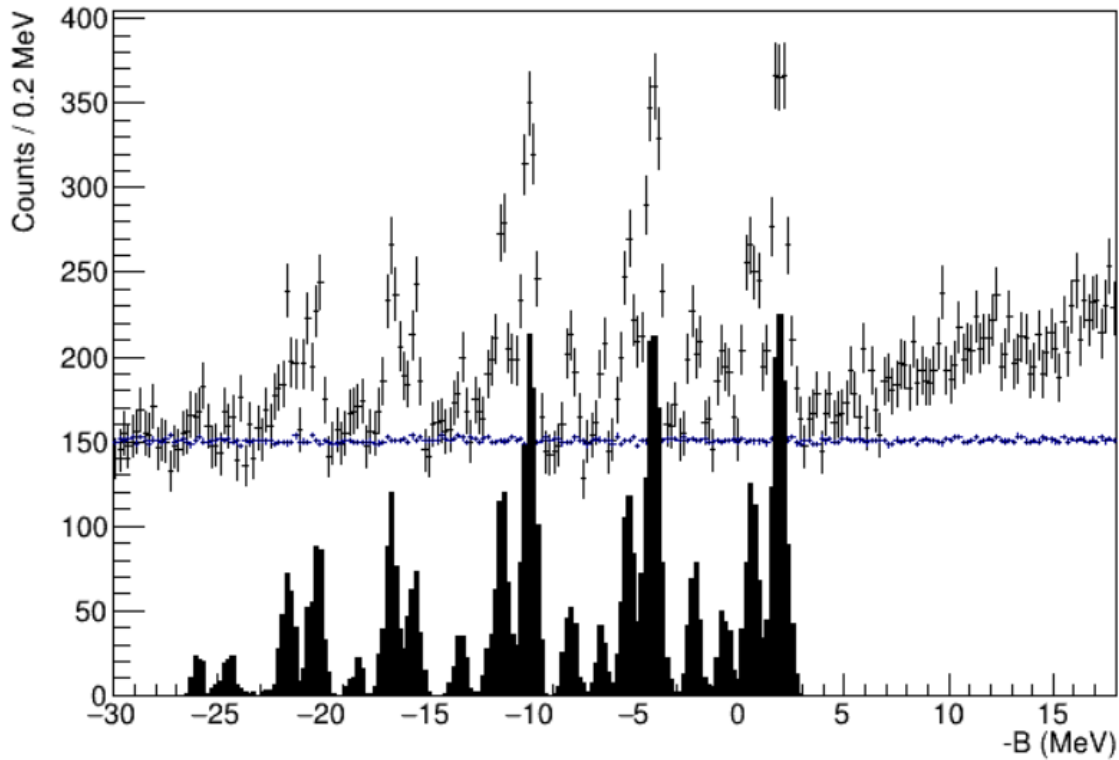
0.6 MeV (FWHM)



Accidental background subtraction
+ fitting

Sample figures ($20 \mu A$, 150 mg/cm^2 , 40 days)

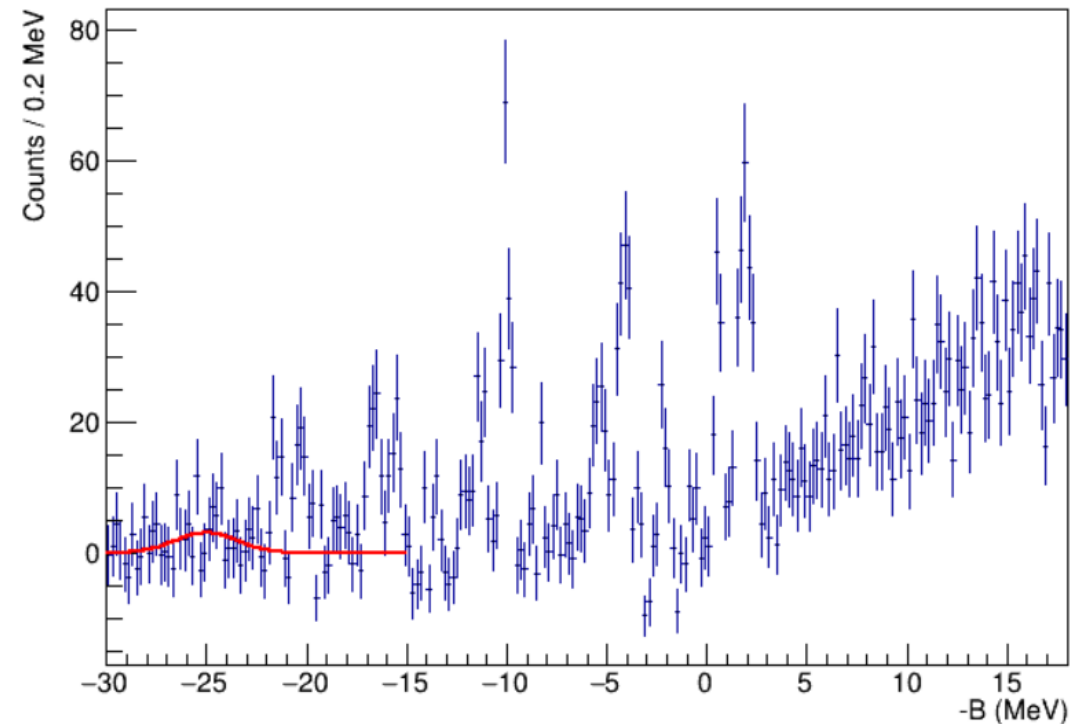
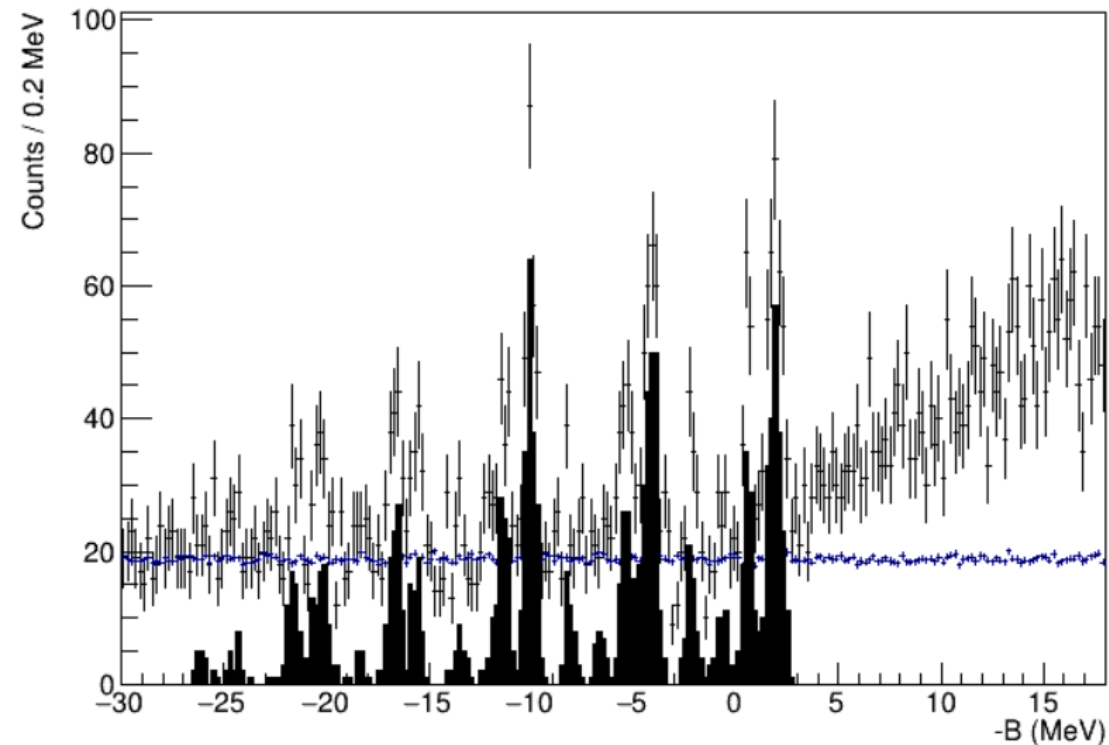
0.6 MeV (FWHM)



Accidental background subtraction
+ fitting

Sample figures ($10 \mu A$, 150 mg/cm^2 , 20 days)

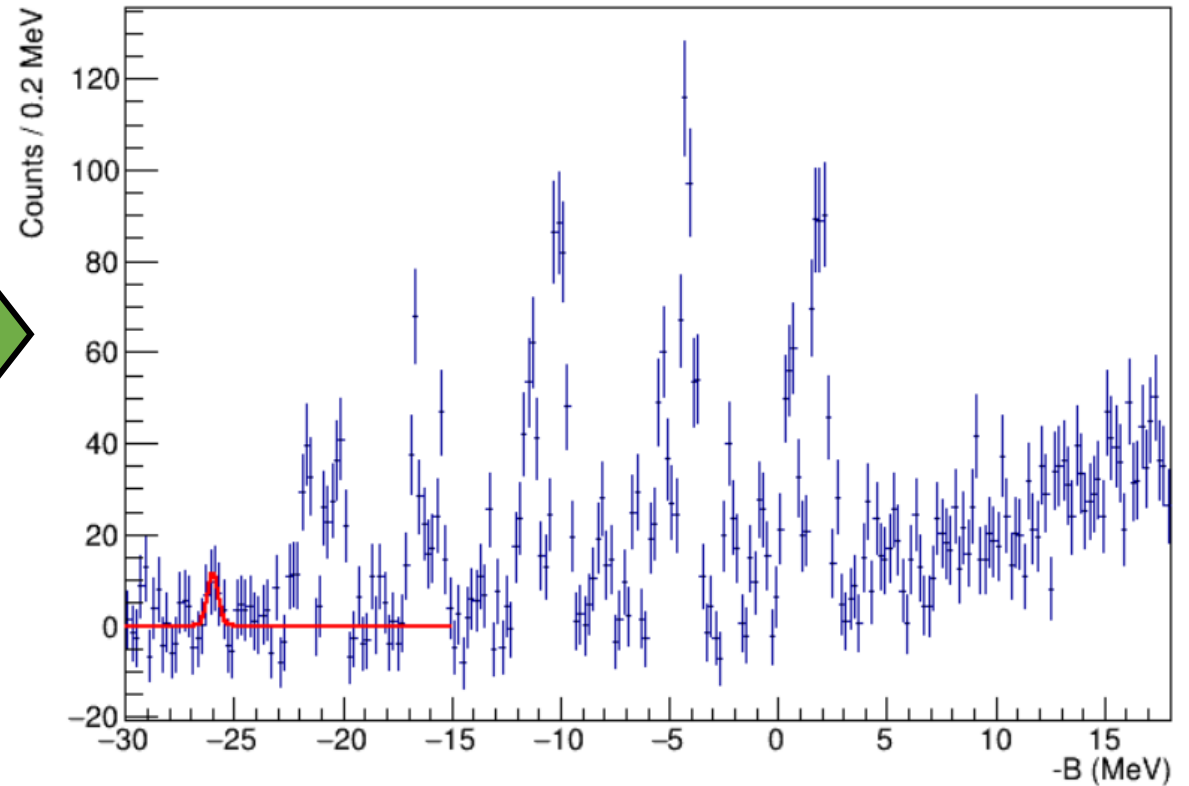
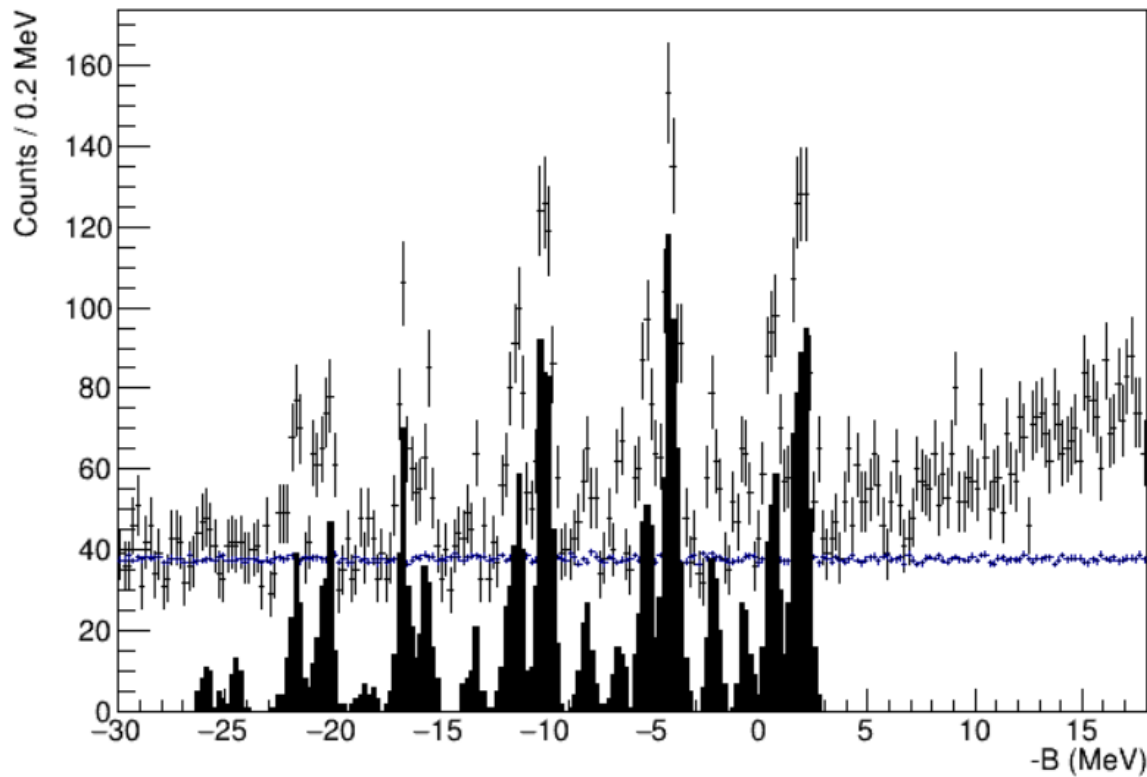
0.6 MeV (FWHM)



Accidental background subtraction
+ fitting

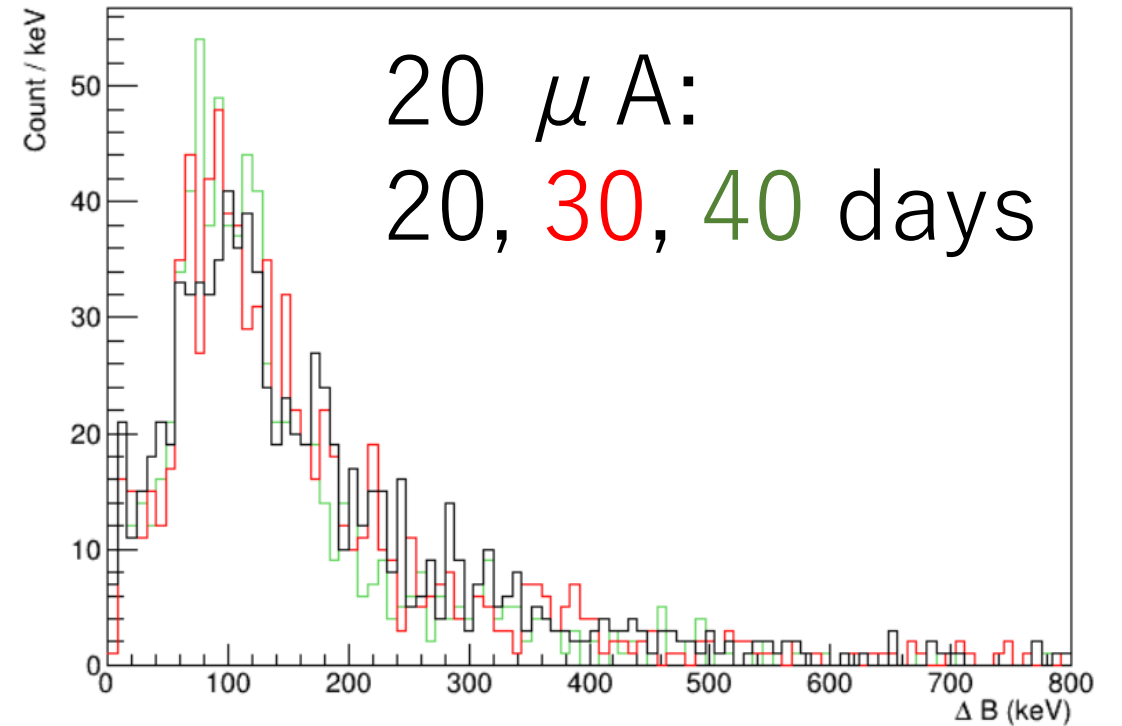
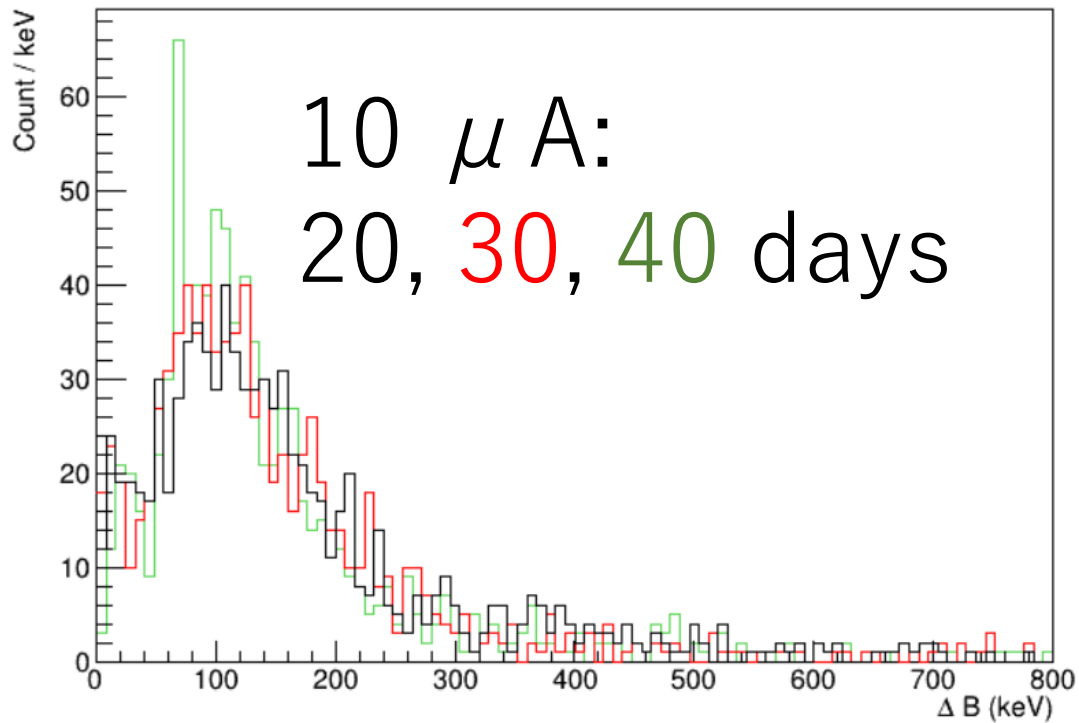
Sample figures ($10 \mu A$, 150 mg/cm^2 , 40 days)

0.6 MeV (FWHM)

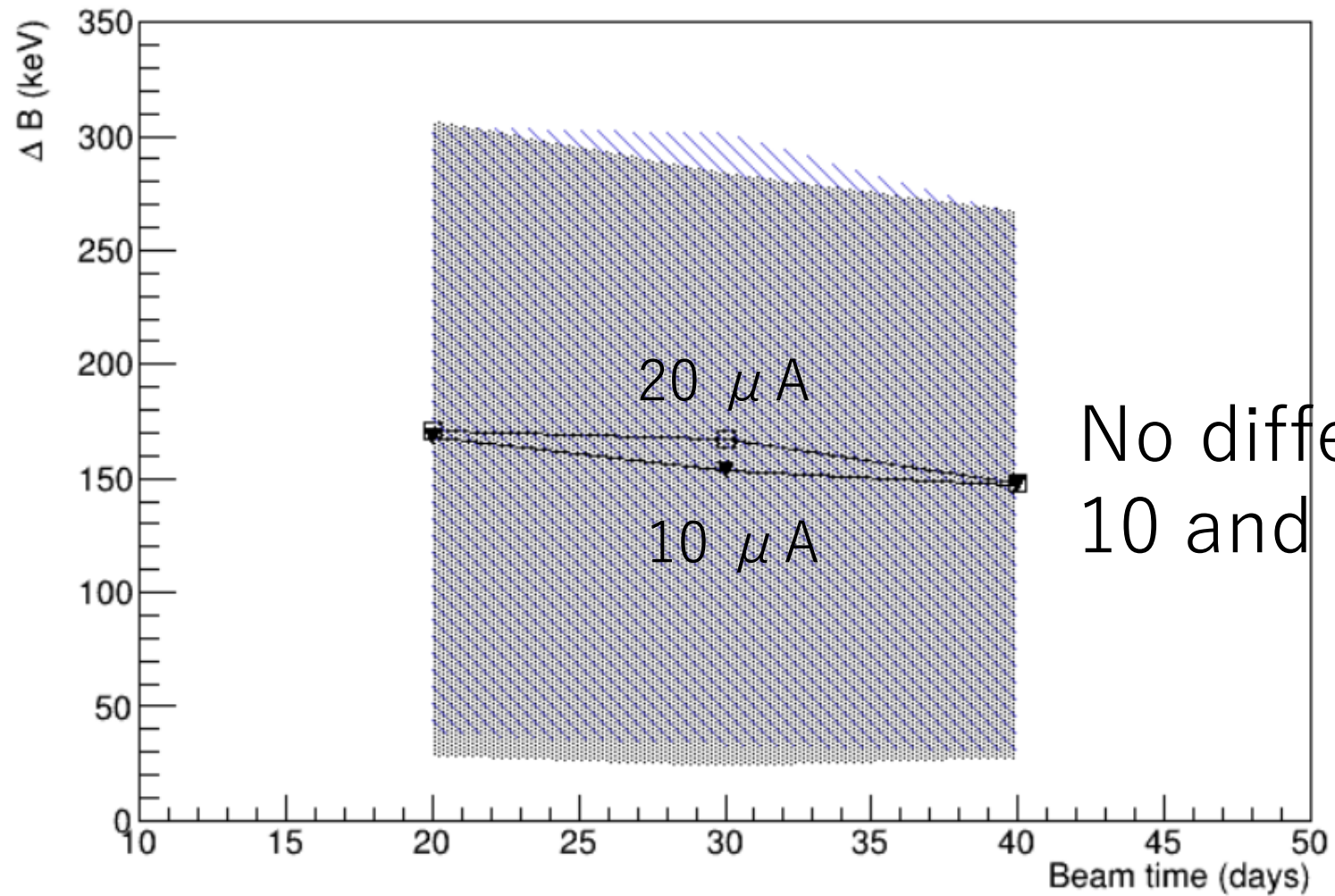


Accidental background subtraction
+ fitting

Statistical errors on B_{\wedge} obtained by fittings for 1000 simulated spectra

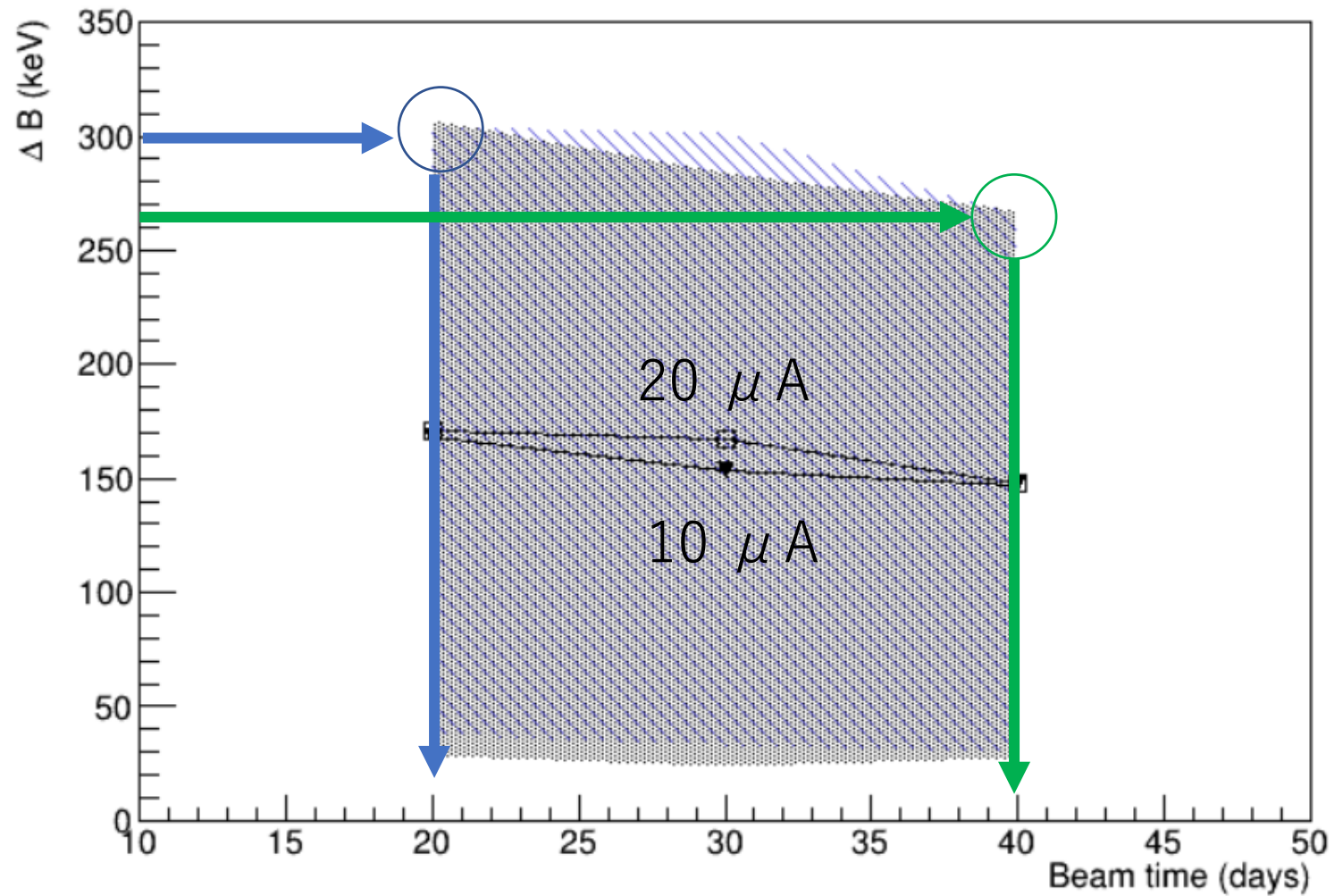


Result



No difference between
10 and $20 \mu A$

Result



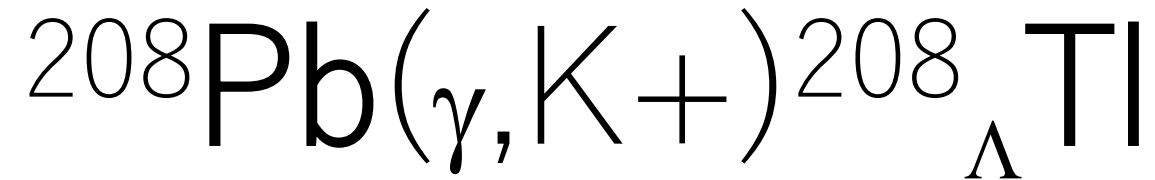
$\Delta B = 300 \text{ keV} \rightarrow 20 \text{ days}$

$\Delta B = 270 \text{ keV} \rightarrow 40 \text{ days}$

Schedule

- Jan: Consensus meeting
- April: 70% Draft
- May: Proposal submission

Backup



Note by T. Motoba (figures updated June 12)

The following three figures summarize the theoretical estimate of the DWIA cross sections calculated at $p(\gamma) = 1.5 \text{ GeV}/c$ and $\theta(K) = 0.5 \text{ deg}$, in which the Saclay-Lyon A amplitudes and the nuclear HO wave functions are employed. Three cases of resolutions (FWHM=0.6, 0.8, 1.0 MeV) are assumed for PR12-18-004. JPG figures are inserted in this .ppt slides, so that one may copy them easily.

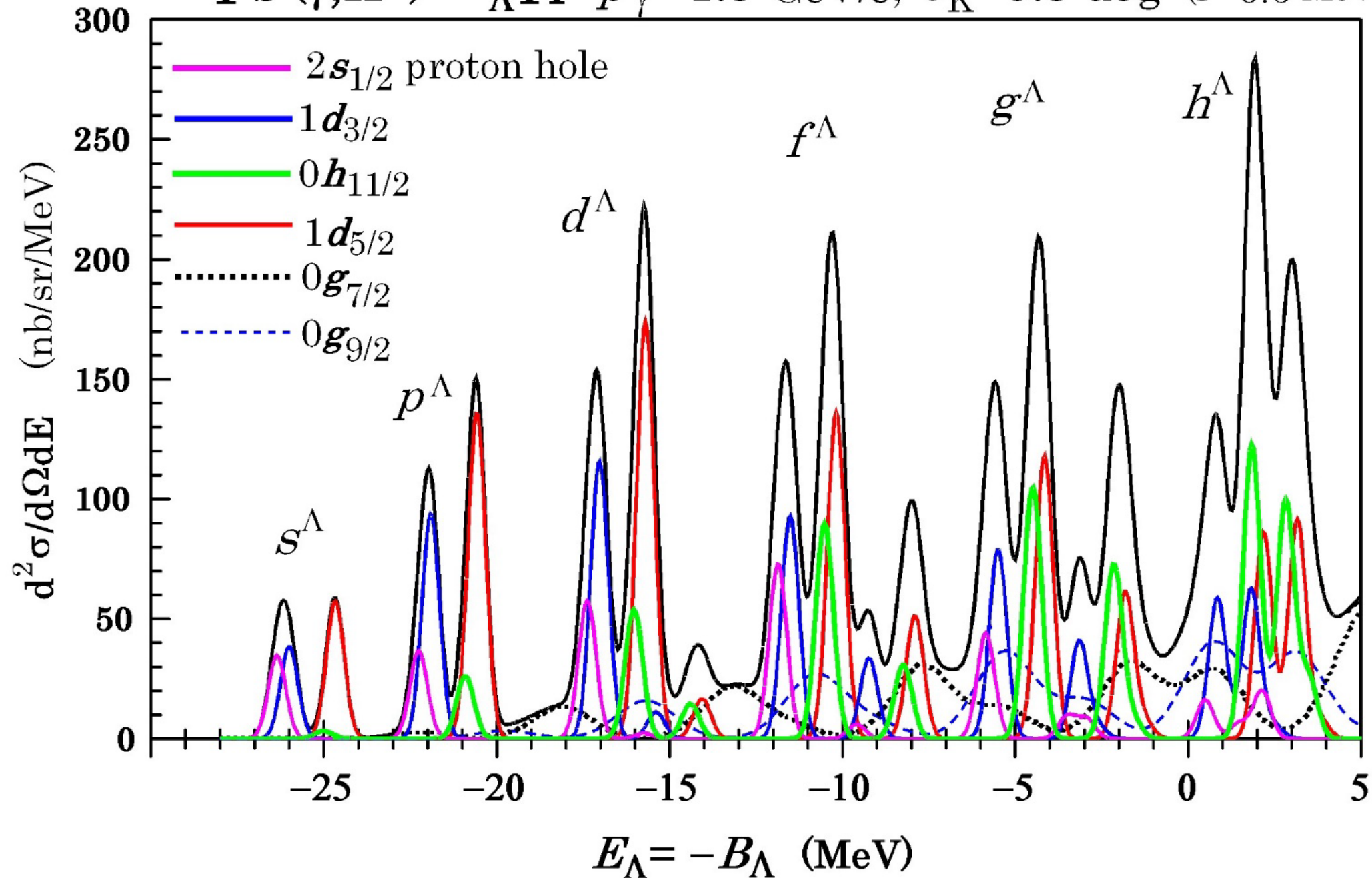
In drawing the spectra, however, the Λ single-particle energies from the Woods-Saxon potential (Millener et al., PRC38) are used instead of HO ones so as to be more realistic: $E(L) = -25.99 \text{ MeV}(0s), -21.90(0p), -17.02(0d), -15.38(1s), -11.50(0f), -9.22(1p), -5.48(0g), -3.14(1d), -2.58(2s), +0.86(0h), +1.84(1f), +2.50(2p)$. These are slightly shifted to have $E(L) = -26.35 \text{ MeV}(0s)$.

On the other hand, the proton single-hole energies are taken from the observed level energies of ^{207}Tl : $E_x = 0.0 \text{ MeV}(2s_{1/2}^{\text{hole}}), 0.351 \text{ MeV}(1d_{3/2}^{\wedge}), 1.348 \text{ MeV}(h_{11/2}^{\wedge}), 1.682 \text{ MeV}(1d_{5/2}^{\wedge}), 4.18 \text{ MeV}(\text{approx. centroid of } 0g_{7/2}^{\wedge}), \text{ and } 6.57 \text{ MeV}(\text{no observed value, but centroid assumed for } 0g_{9/2})$. Note that the constant spreading widths of 2 MeV are assumed to take account of the fragmented proton $0g_{7/2}$ and $0g_{9/2}$ orbits.

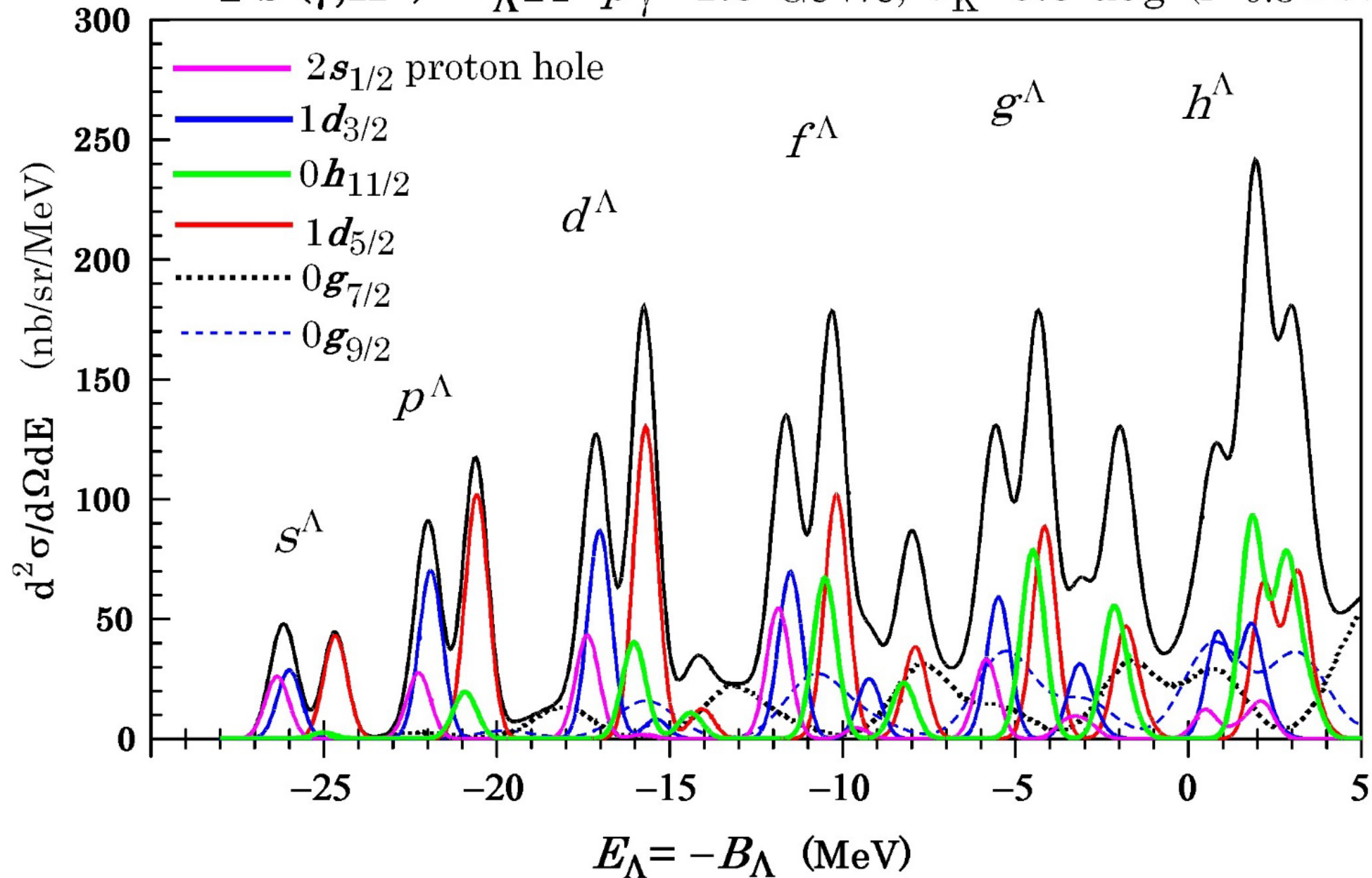
The cross section estimates should be further improved if we use the followings:

- 1) the Woods-Saxon wave functions for both p and Λ in X-S estimates, and
- 2) the better elementary amplitudes which explains forward scatterings.

$^{208}\text{Pb} (\gamma, \text{K}^+) ^{208}_{\Lambda}\text{Tl} \quad p_{\gamma}=1.5 \text{ GeV}/c, \theta_{\text{K}}=0.5 \text{ deg} \quad (\Gamma=0.6 \text{ MeV})$



$^{208}\text{Pb} (\gamma, \text{K}^+) ^{208}_{\Lambda}\text{Tl}$ $p_{\gamma}=1.5 \text{ GeV}/c$, $\theta_{\text{K}}=0.5 \text{ deg}$ ($\Gamma=0.8 \text{ MeV}$)



$^{208}\text{Pb} (\gamma, \text{K}^+) ^{208}_{\Lambda}\text{Tl} \quad p_{\gamma}=1.5 \text{ GeV}/c, \theta_{\text{K}}=0.5 \text{ deg} \quad (\Gamma=1.0 \text{ MeV})$

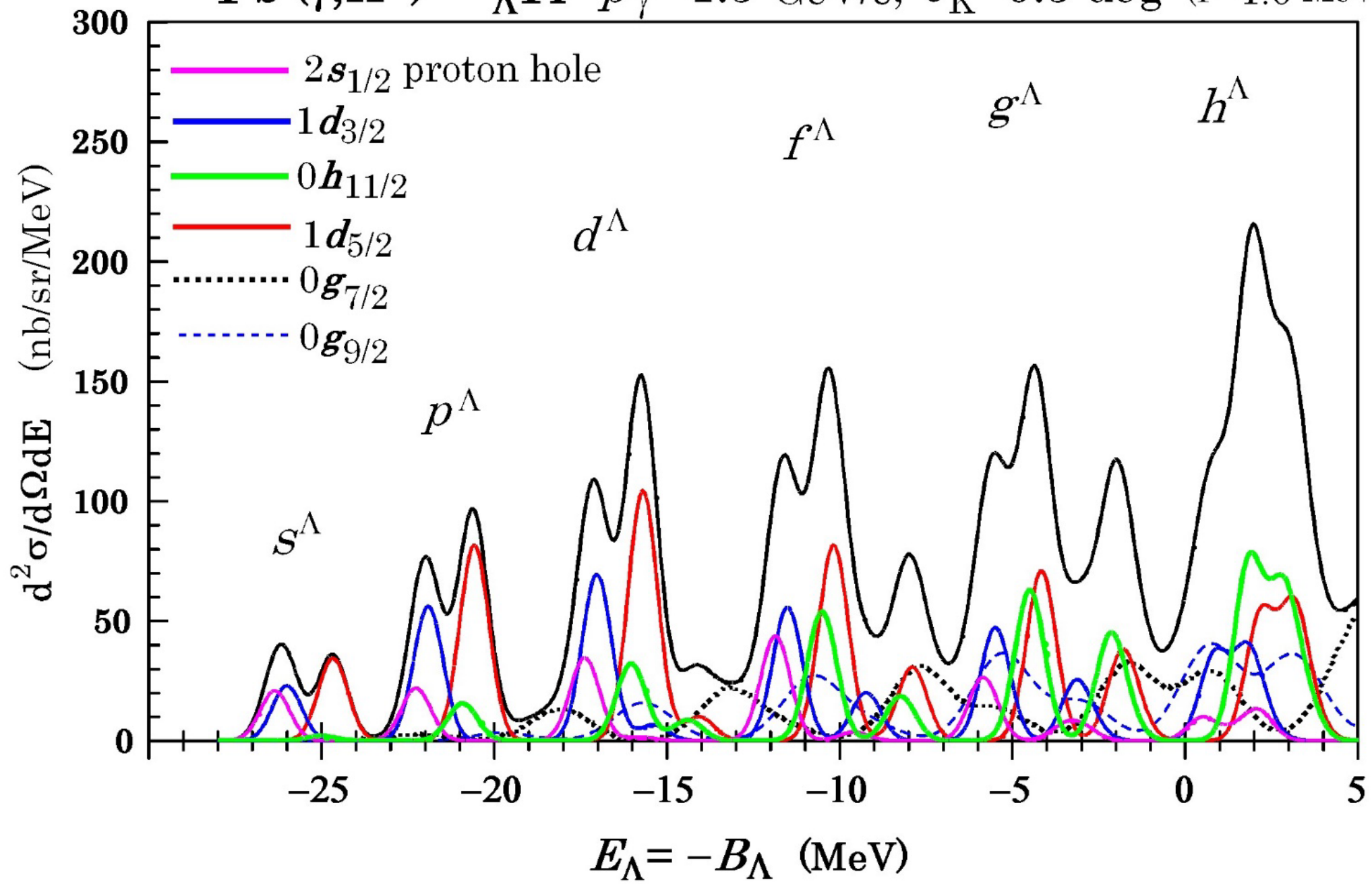


Table 2: DWIA Cross Sections (nb/sr) for $^{208}\text{Pb}(\gamma, K^+)_{\Lambda}^{208}\text{Tl}$ at $p_{\gamma} = 1.5\text{GeV}/c$ and $\theta_K = 0.5$ deg.

	s^{Λ}	p^{Λ}	d^{Λ}		f^{Λ}	
Core state $(E_x) \times (nlj)^{\Lambda}$	$0s_{1/2}^{\Lambda}$	$0p_{3/2}^{\Lambda} + 0p_{1/2}^{\Lambda}$	$0d_{5/2}^{\Lambda} + 0d_{3/2}^{\Lambda}$	$1s_{1/2}^{\Lambda}$	$0f_{7/2}^{\Lambda} + 0f_{5/2}^{\Lambda}$	$1p_{3/2}^{\Lambda} + 1p_{1/2}^{\Lambda}$
$1/2^+(E_x=0.0): \pi(2s_{1/2}^{-1})$	22.2	15.6+ 7.8	22.1+14.8	1.6	26.6+20.0	2.6+1.3
$3/2^+(0.351): \pi(1d_{3/2}^{-1})$	24.4	46.4+13.4	52.4+21.6	7.1	41.0+18.4	15.8+5.6
Left peak $d\sigma/d\Omega$	46.6	83.2	110.9	9.7	106.0	25.3
$(E_{\Lambda} = -B_{\Lambda}$ in MeV)	(-26.16)	(-21.97)	(-17.11)		(-11.64)	
$11/2^-(1.348): \pi(0h_{11/2}^{-1})$	2.1	10.6+6.1	18.7+15.6	9.3	29.9+27.4	12.8+7.0
$5/2^+(1.682): \pi(1d_{5/2}^{-1})$	36.7	51.7+35.2	58.1+52.4	10.6	42.9+44.0	20.3+12.4
Right peak $d\sigma/d\Omega$	39.8	103.5	144.8	19.9	149.3	52.5
$(E_{\Lambda} = -B_{\Lambda}$ in MeV)	(-24.70)	(-20.60)	(-15.74)		(-10.32)	(-7.96)

Table 2 corresponds to set $E_{gs}({}_{\Lambda}^{208}\text{Tl}) = -26.35$ MeV (Exp). Figure shows a series of doublet peaks indicated respectively by s^{Λ} , p^{Λ} , d^{Λ} , f^{Λ} , etc. As known from the energy differences between low-lying energy levels of ^{207}Tl , the proton-hole states are classified into two nearly degenerate groups in view of the ‘critical’ value $\Delta E_x \simeq 0.35$ MeV. The left member of each doublet is attributed to the structure $[\text{core}(1/2^+, 3/2^+) \times (nlj)^{\Lambda}]$, while the right member to $[\text{core}(11/2^-, 5/2^+) \times (nlj)^{\Lambda}]$. In the present calculation the elementary amplitude from the Saclay-Lyon model A is employed, but it should be noted that SLA leads to considerable overestimate at very forward angle $\theta_K^{\text{Lab}} \lesssim 5$ deg when compared with other theoretical models and/or experimental behaviors ($p_{\gamma} \simeq 1.3$ GeV/c). One may refer to JPS Conf. Proc. 17, 011003 (2017) to see how the theoretical cross sections change depending on p_{γ} and θ_K .